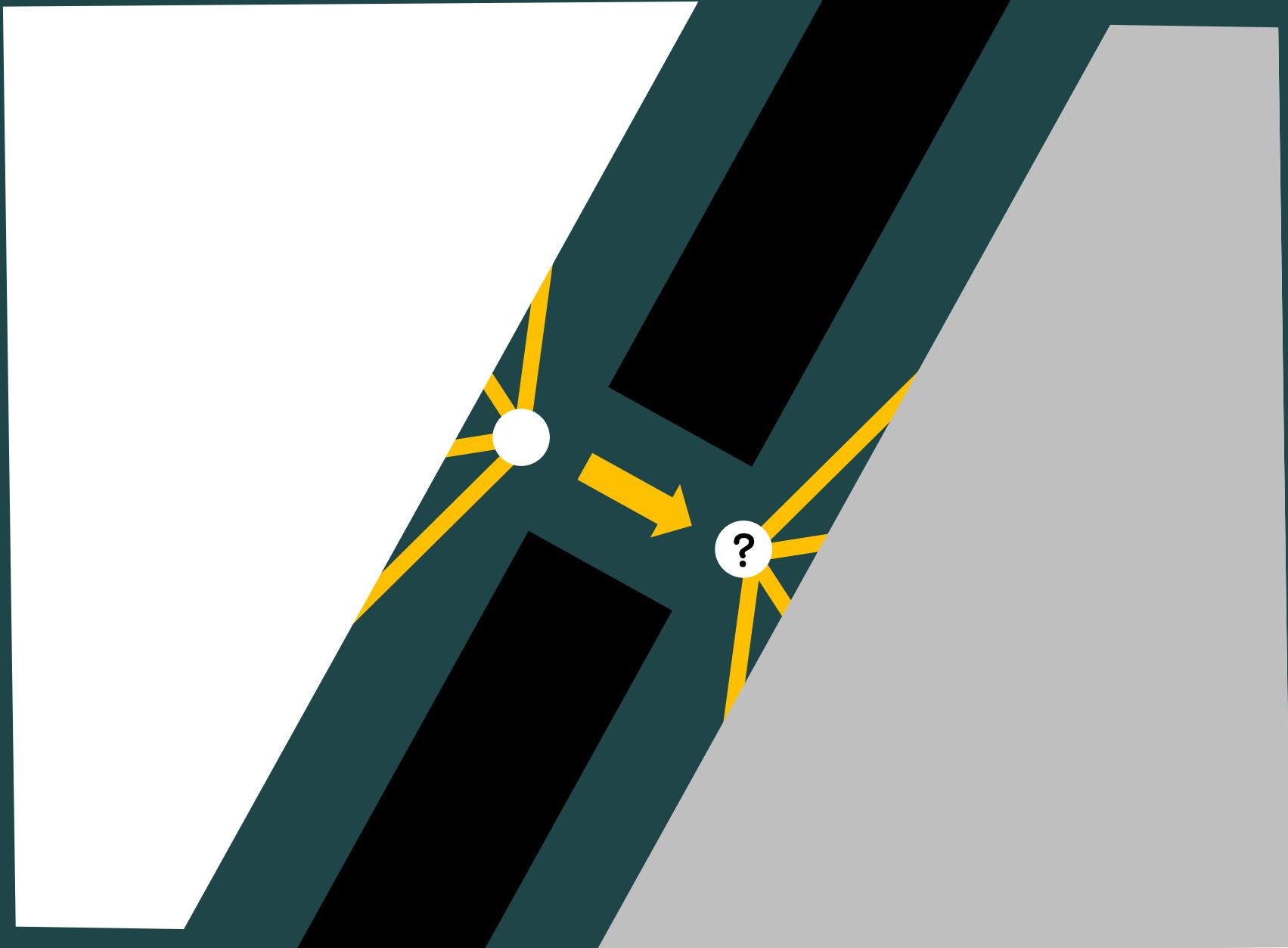


local tests of global entanglement and a counterexample to the generalized area law







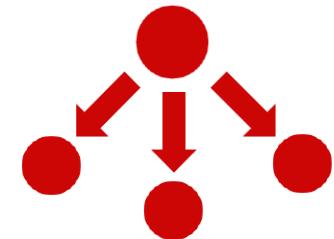
!



1

q. expanders

maximally entangled states



2

entanglement

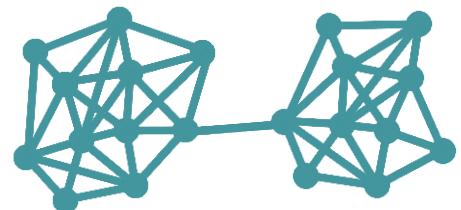
testing and communication

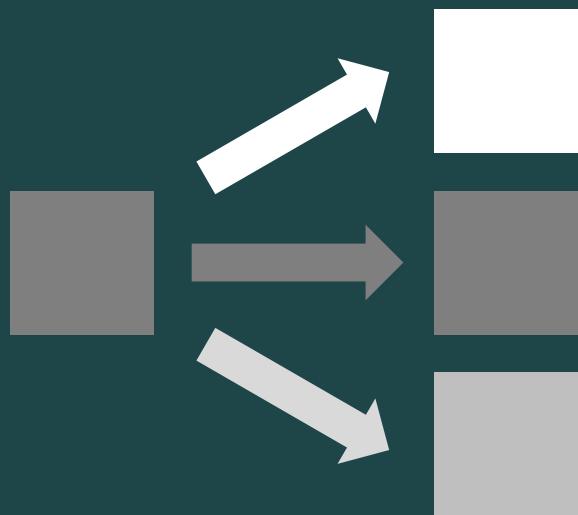


3

area law

gaps, connections, correlations

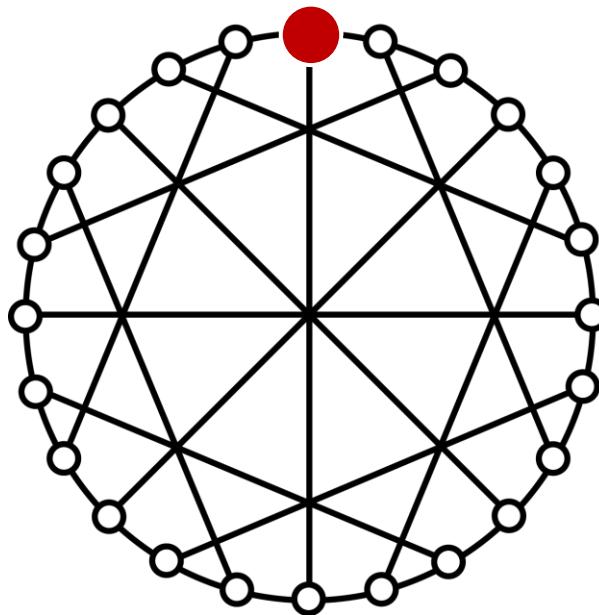




Quantum
Expanders

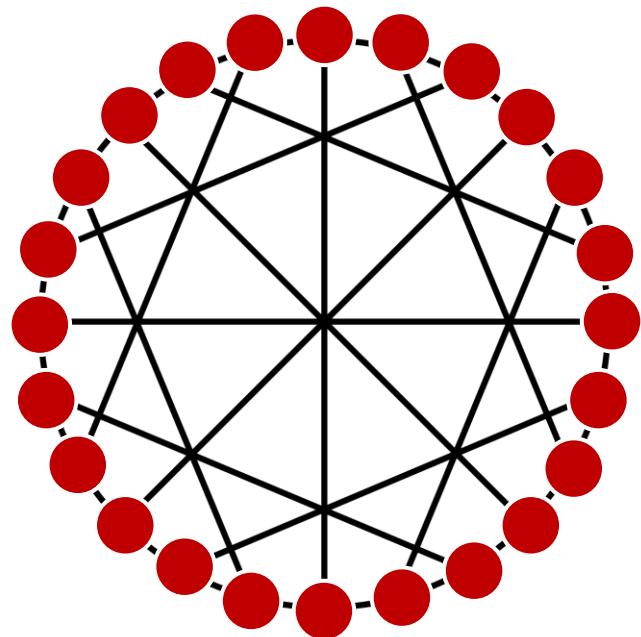
1 Classical expanders

- walk on these graphs? mix fast!



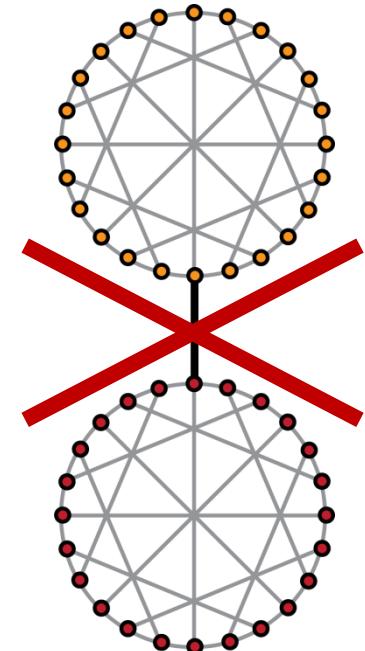
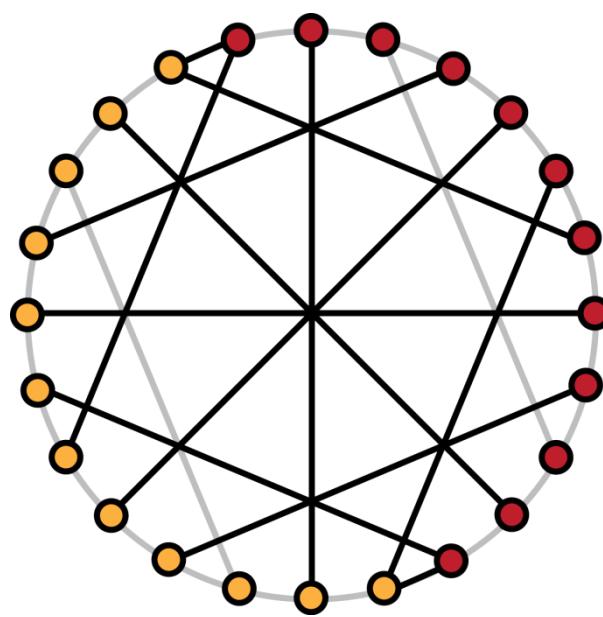
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1 Classical expanders

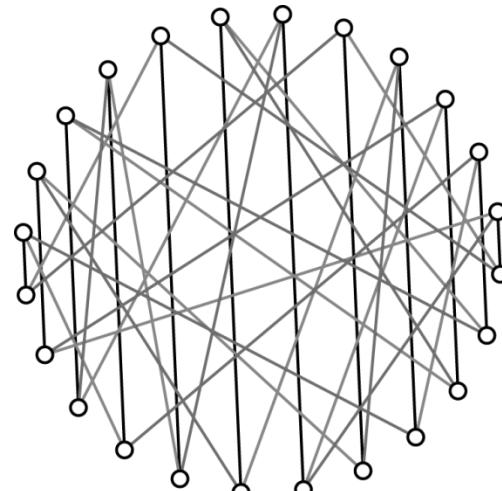
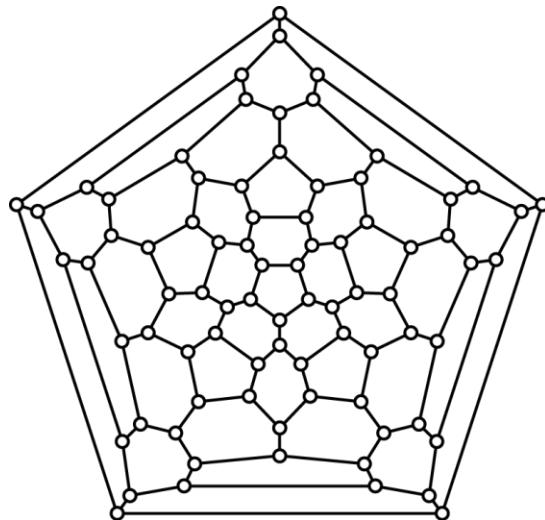
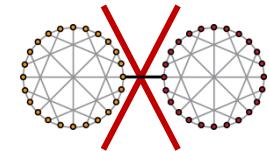
- walk on these graphs? mix fast!
divide in two? cut a lot (fraction) of edges!



1 Classical expanders

- walk on these graphs? mix fast!
divide in two? cut a lot (fraction) of edges!

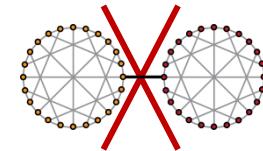
examples: Ramanujan, Cayley



1 Classical expanders

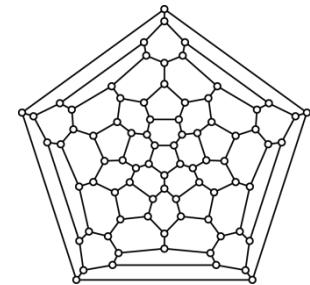
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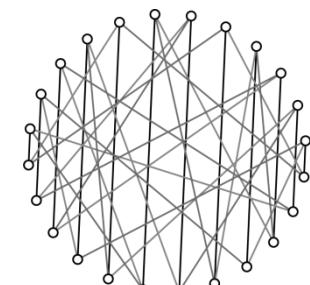


- explicit, constant-degree
approximations to the full graph

normalized adjacency matrix
second largest eigenvalue $1-\lambda$



- a review [Hoory Linial Wigderson]
a talk [Harrow quantum expanders youtube]

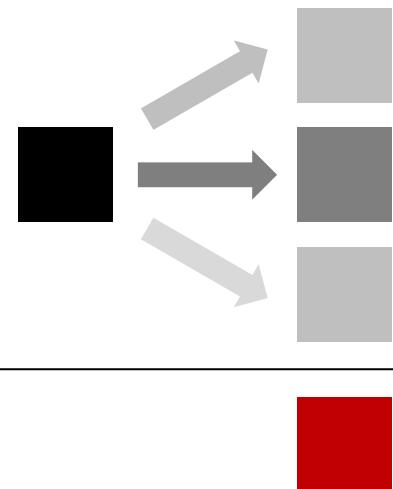


1 Mixing up something quantum

- applying random unitaries

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

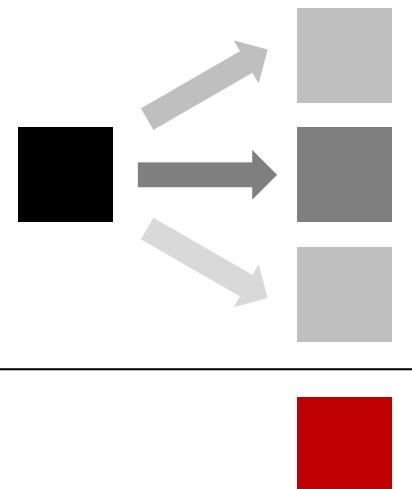
- classical expanders:
explicit, constant-degree
approximations to the full graph
fast-mixing



1 Quantum expanders

- applying random unitaries from a small set a discrete approx. to Haar

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$



- classical expanders:
explicit, constant-degree
approximations to the full graph
fast-mixing

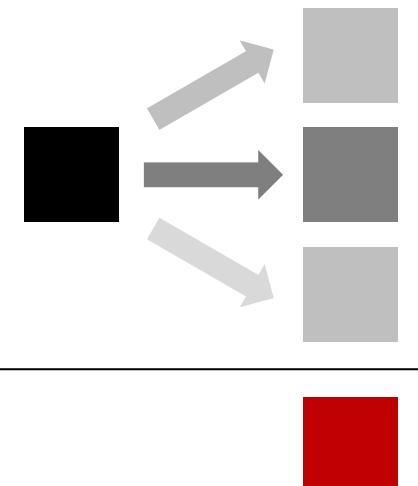
1 Quantum expanders

- applying random unitaries from a small set a discrete approx. to Haar
- transform $N \times N$ matrices \mathbb{I} stays, everything else changes

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

$$\|\mathcal{E}(X)\|_2 \leq \lambda \|X\|_2$$

small 2nd largest sing. value



- QE constructions for fixed $k (\dots, 8, 4, 3)$

$$\lambda \approx k^{-c}$$

[Ben-Arroyo+ 07, Hastings '07, Gross & Eisert '08, Hast. & Harr. '09, Gross '15]

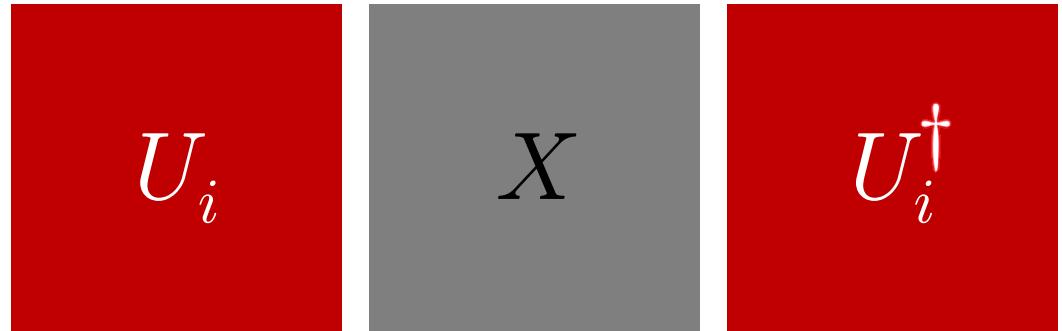
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



$$U_i \ X \ U_i^\dagger$$

$$U_i X = X U_i$$

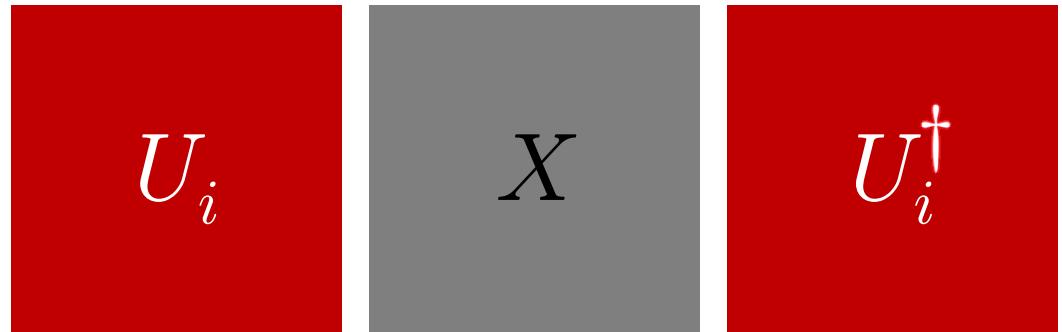
1 Quantum expanders

- transform $N \times N$ matrices

a matrix that
doesn't change?

$$X = \mathbb{I}$$

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$



$$X \quad | \quad U_i \quad | \quad U_i^\dagger$$

$$U_i X = X U_i$$

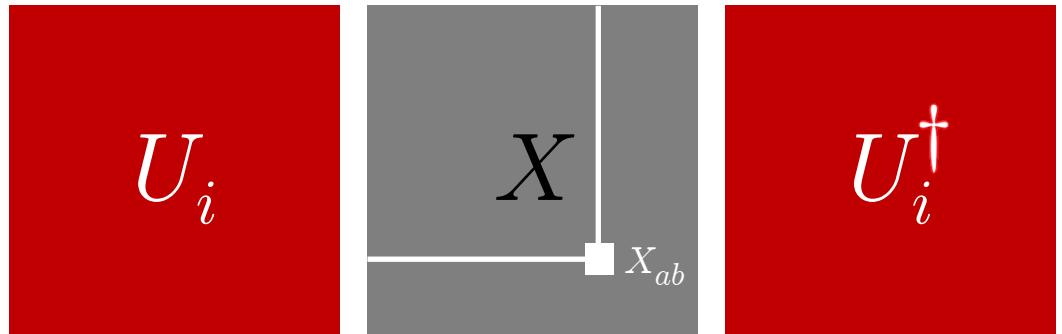
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



- interpreting matrices
as 2-register states

$$X = \sum_{a,b} X_{ab} |a\rangle\langle b|$$

density matrix

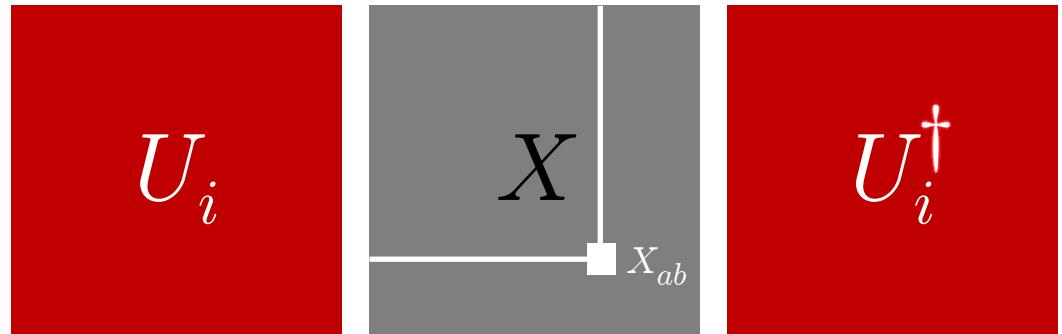
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



- interpreting matrices
as 2-register states

$$|X\rangle = \sum_{\substack{\text{pure} \\ \text{state}}} X_{ab} |a\rangle |b\rangle$$

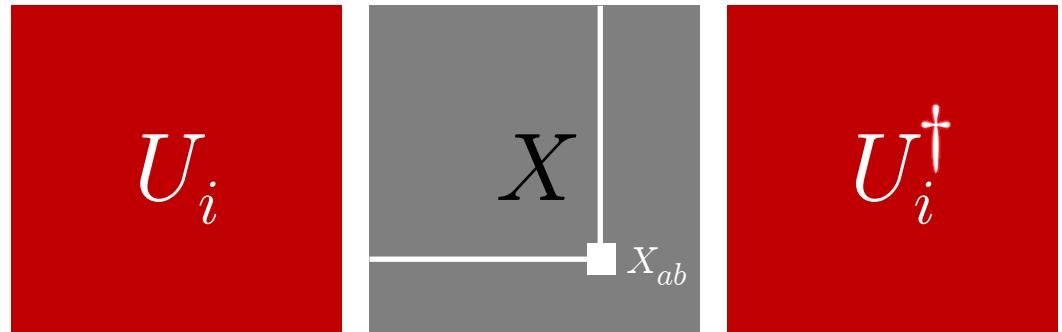
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



- interpreting matrices
as 2-register states

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle|b\rangle$$

applying an expander distributively

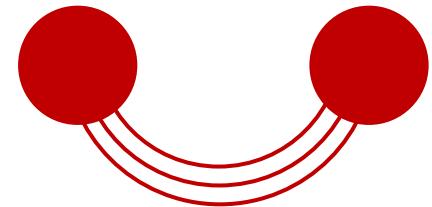
a stationary state?

max. entangled!

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$

1 Quantum expanders & 2 registers

- transform $N \otimes N$ states close to the depolarizing channel



$$\|\tilde{\mathcal{E}} - |\phi_N\rangle\langle\phi_N|\| = \lambda$$

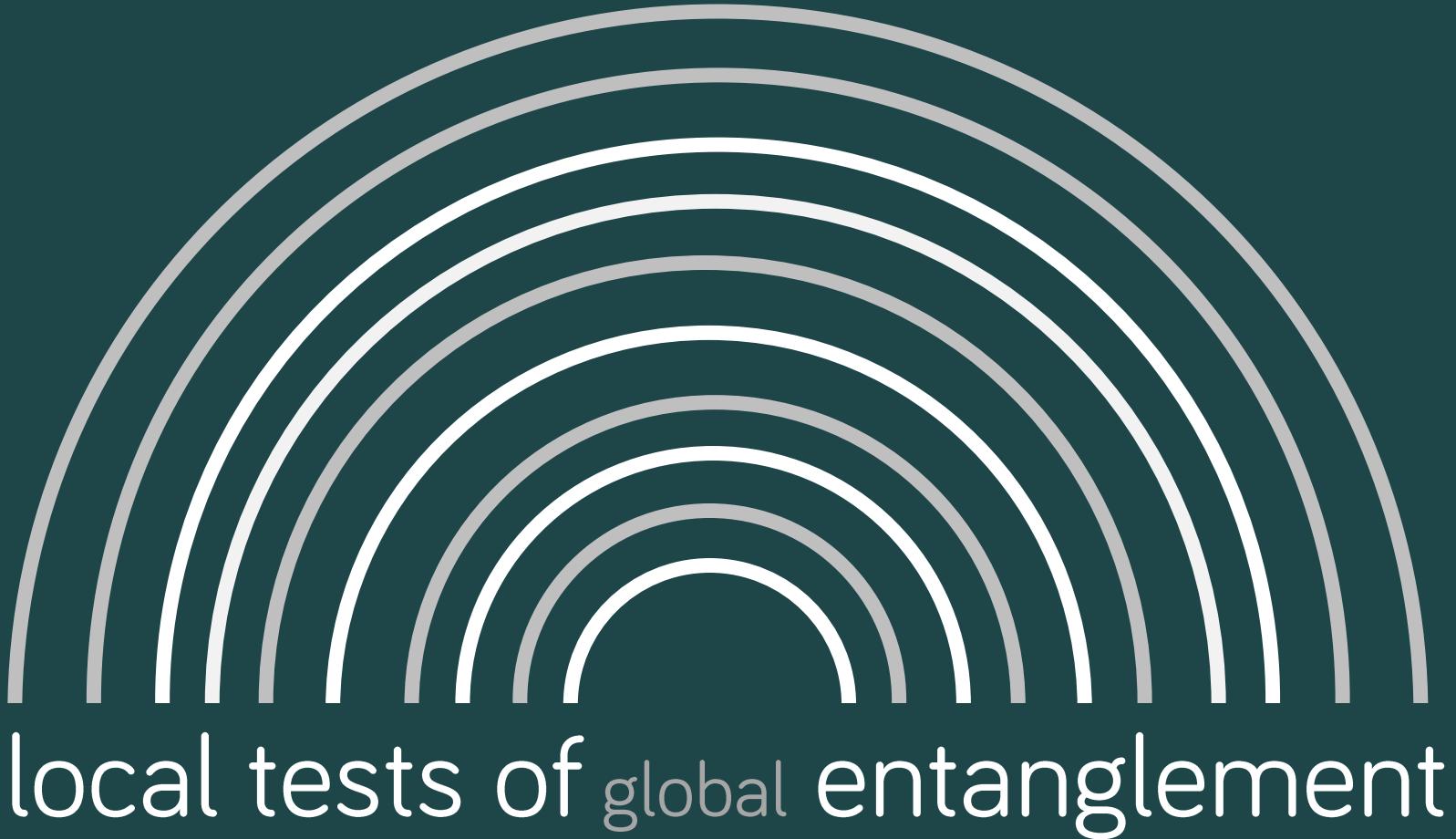
- interpreting matrices as 2-register states

a stationary state?

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle|b\rangle$$

applying an expander distributively

max. entangled! $\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$



local tests of global entanglement

2 EPR testing

- how costly is it to certify that we share a maximally entangled state?

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$



$O(\log\log(N) + \log(1/\lambda))$ qubits. [BDSW '96, BCGST '02]

2 EPR testing

- how costly is it to certify that we share a maximally entangled state?

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$



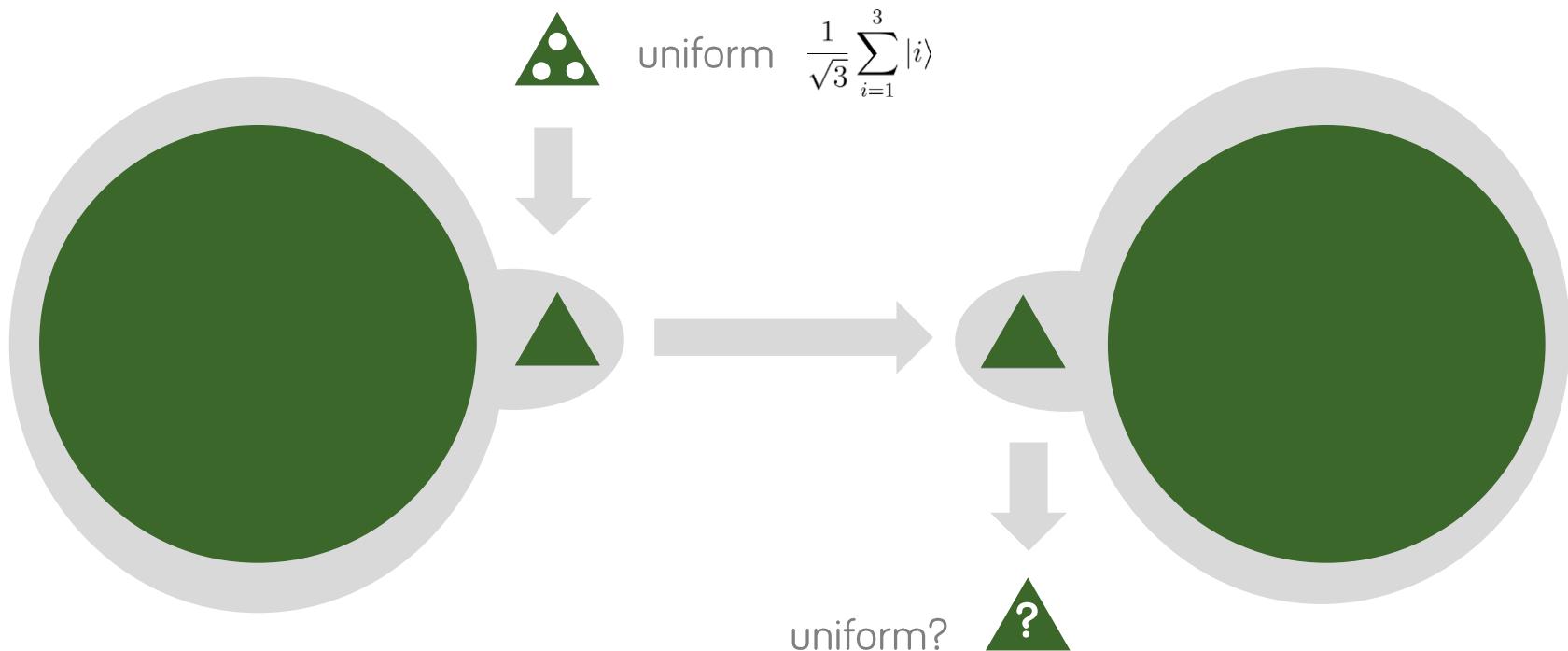
$O(\log(1/\lambda))$ qubits for error λ .

2 EPR testing

- how costly is it to certify that we share a maximally entangled state?
- apply a quantum expander distributively

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$

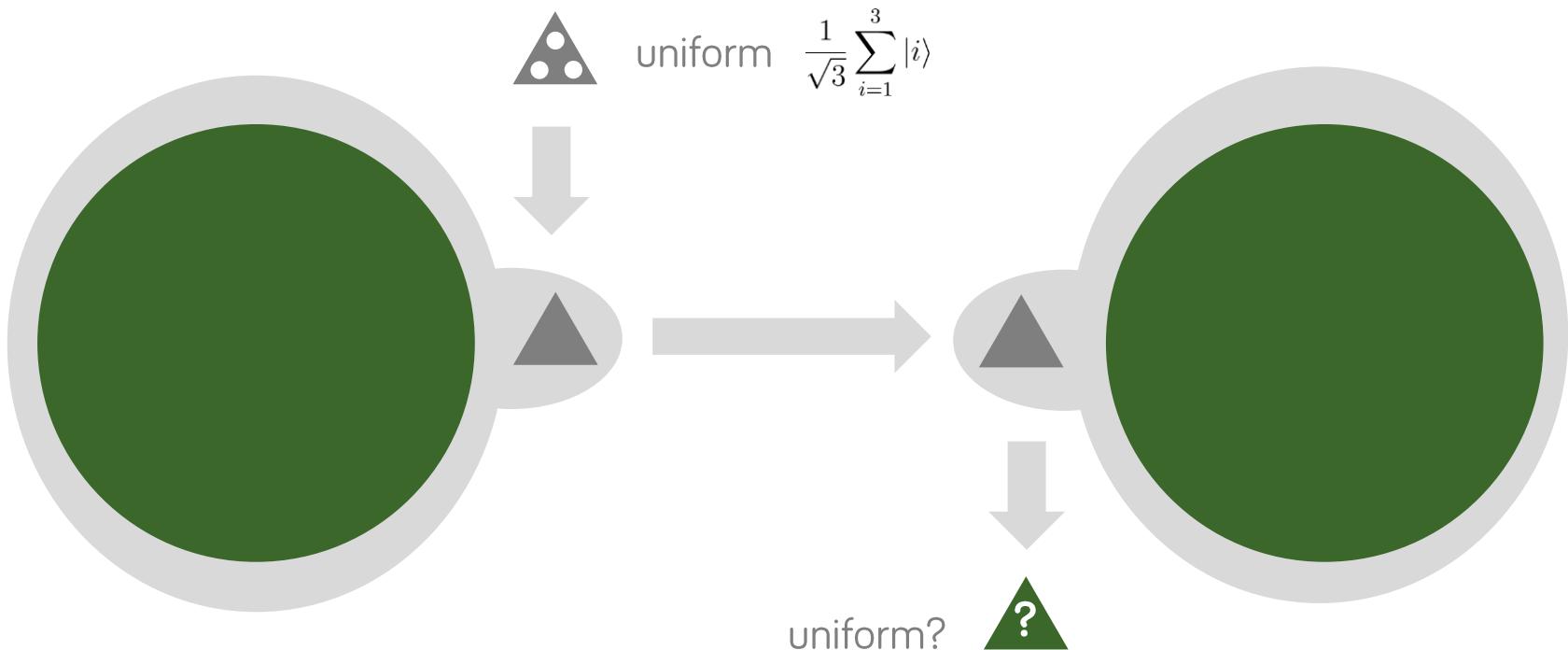
$$U_i \otimes U_i^*$$



2 EPR testing

- when does the qutrit remain uniform?

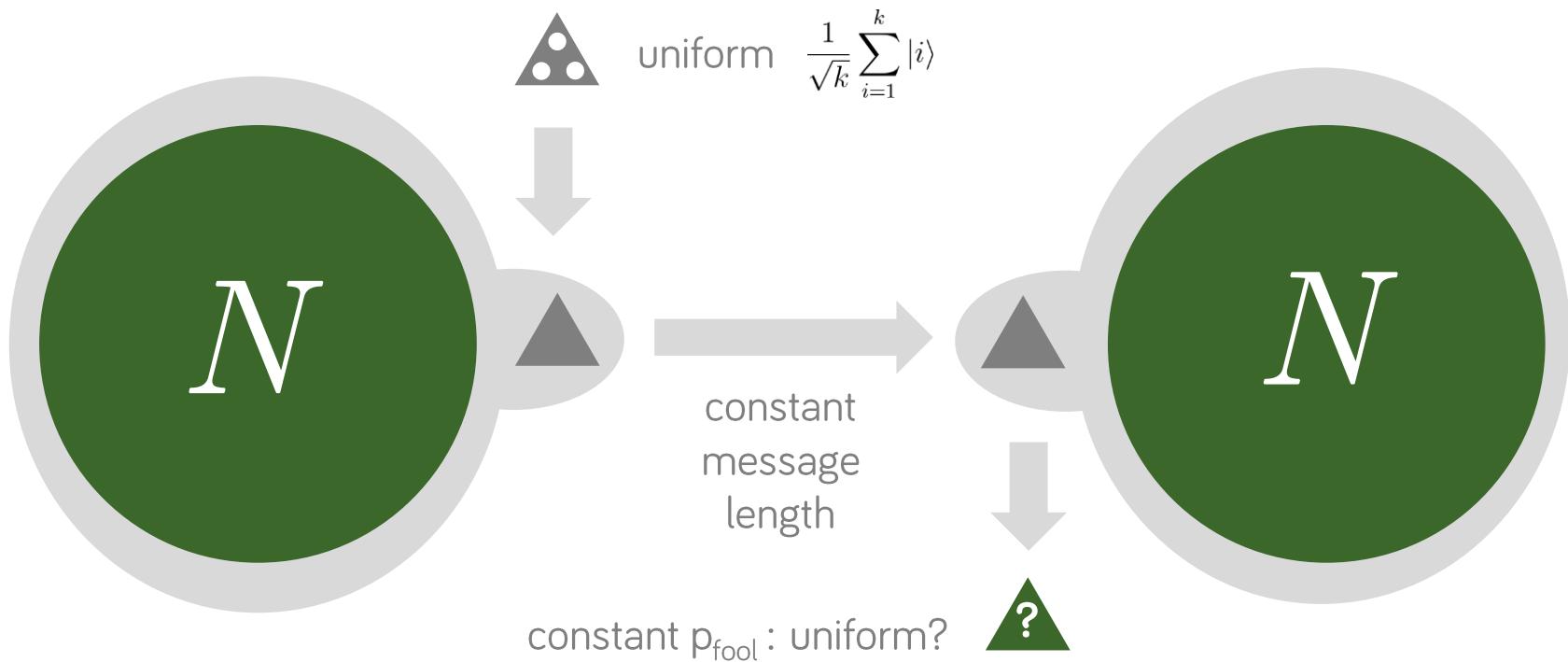
$$\frac{1}{\sqrt{3}} \sum_{i=1}^3 |i\rangle (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle |b\rangle$$



2 EPR testing

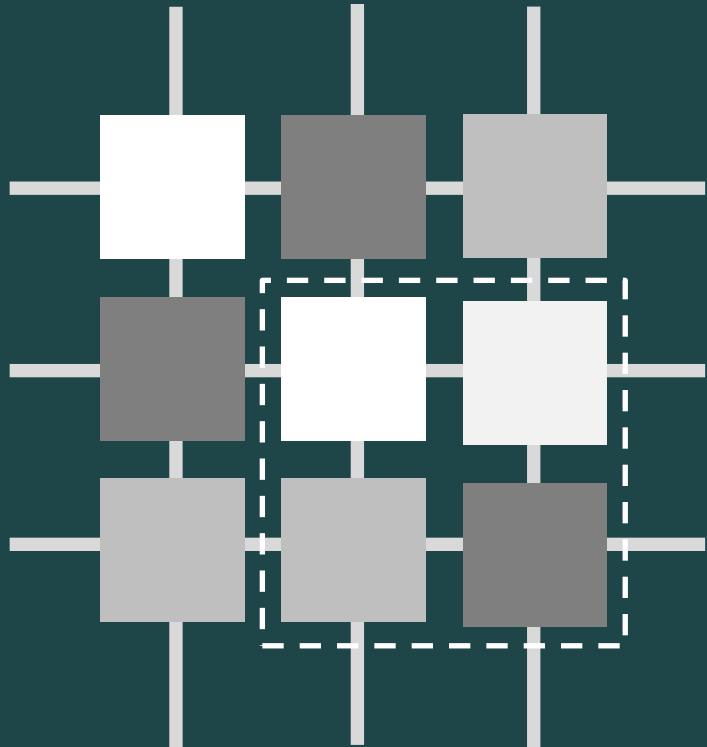
- when does the qutrit remain uniform?
for max. entangled X
- quantum expander property ... soundness

$$\frac{1}{\sqrt{3}} \sum_{i=1}^3 |i\rangle (U_i \otimes U_i^*) \frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$$



little communication
local tests
constant error

global correlations



a counterexample to the
generalized
area law

few connections

local interactions

constant gap

local correlations

few connections

local interactions

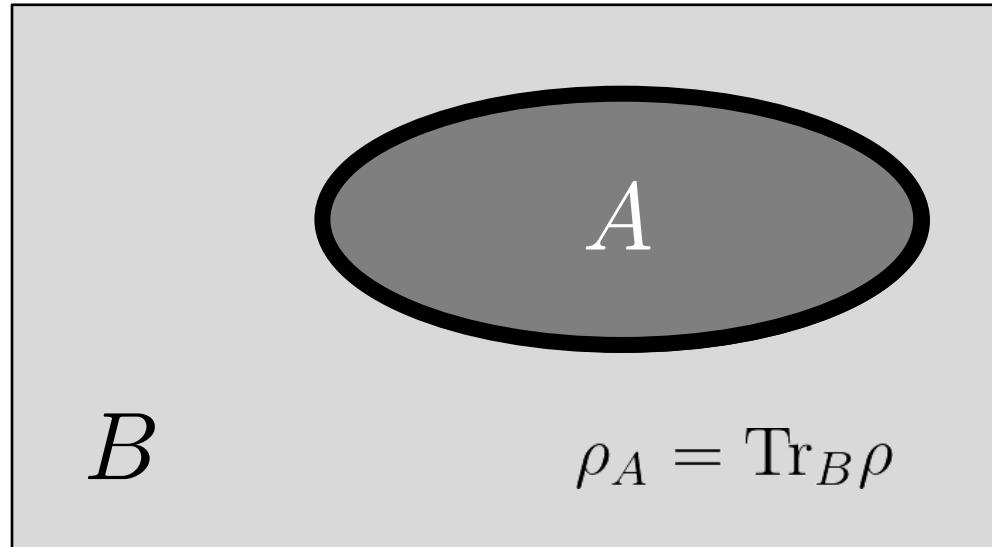
constant gap

global correlations

3 Ground states of gapped quantum spin systems

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \frac{\cancel{\text{volume}}}{\text{surface area}}$$



area law \rightarrow a simple ground state?

3 Gapped 1D Hamiltonians

- Nothing closer than Δ to the ground state.

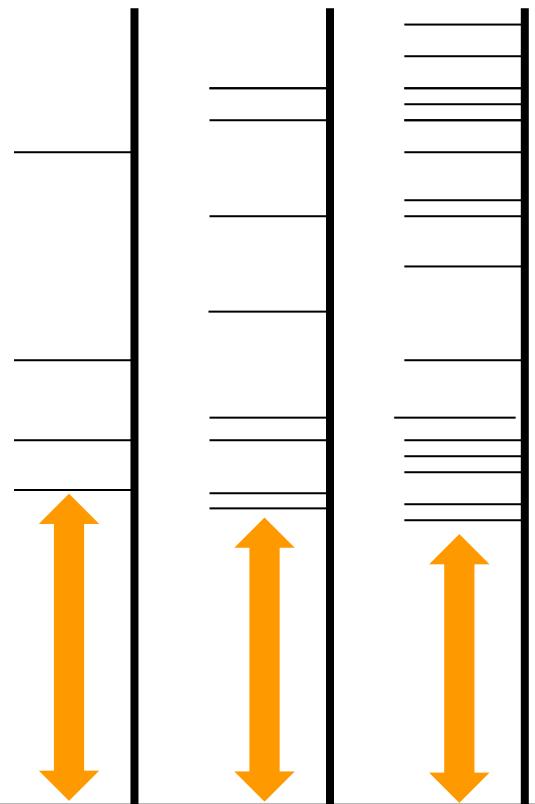
$$N \rightarrow \infty$$

the AKLT (spin-1) chain

$$\sum_{j=1}^{N-1} \vec{S}_j \cdot \vec{S}_{j+1} + \frac{1}{3} \left(\vec{S}_j \cdot \vec{S}_{j+1} \right)^2$$

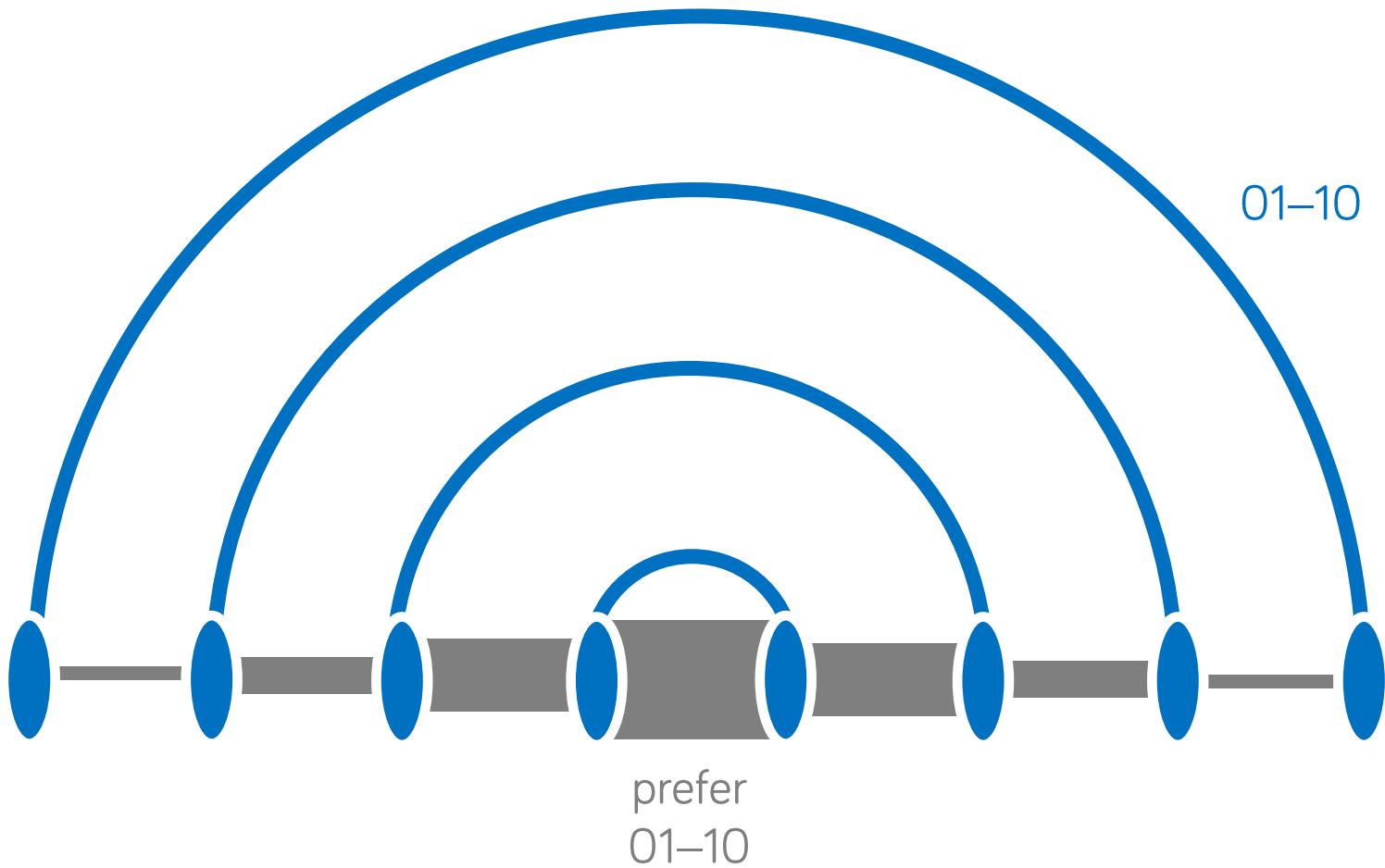
a biased walk in 1D

$$\sum_{j=1}^{N-1} (|j\rangle - B|j+1\rangle) (\langle j| - B\langle j+1|)$$



exp. falloff of correlations, MPS/DMRG work well

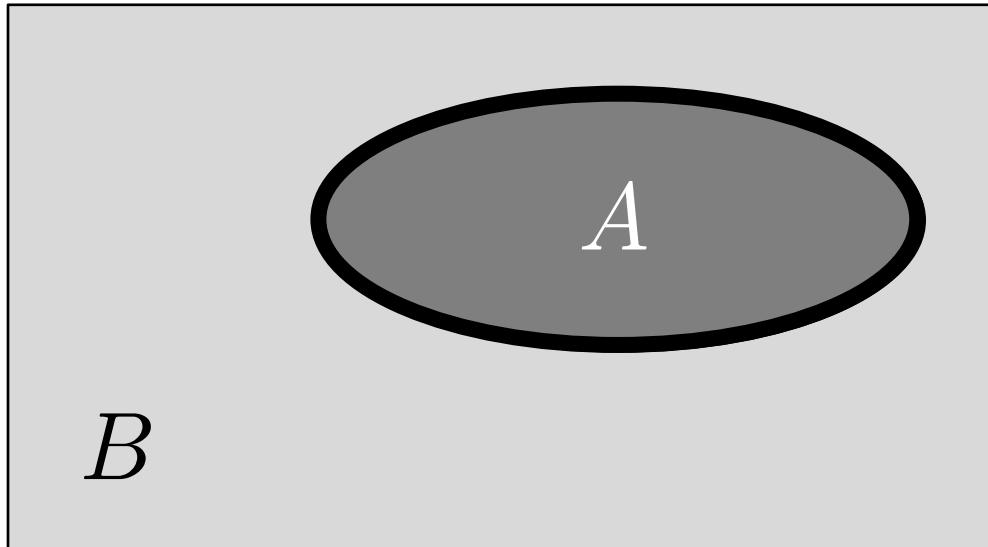
Without a gap, the entropy can be large. [Verstraete, Latorre+]



3 Ground states of gapped H's & the area law

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \frac{\cancel{\text{volume}}}{\text{surface area}}$$

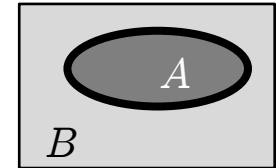


a gapped system  an area law

3 Ground states of gapped H's & the area law

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \frac{\cancel{\text{volume}}}{\text{surface area}}$$



1D ... theorems [Hastings 07, Arad+ 13]

algorithms [White 92, Vidal 03, Landau+ 13]

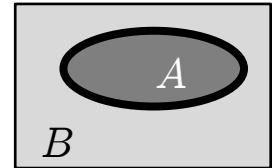
2D ... we're close

a gapped system  an area law

3 Ground states of gapped H's & the area law

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \frac{\cancel{\text{volume}}}{\text{surface area}}$$

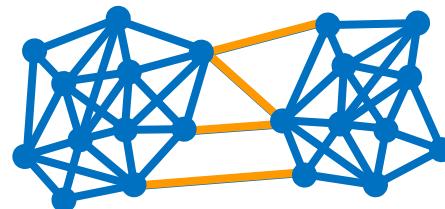


1D ... theorems [Hastings 07, Arad+ 13]

algorithms [White 92, Vidal 03, Landau+ 13]

2D ... we're close

- generalized area conjecture
entropy \sim cut size



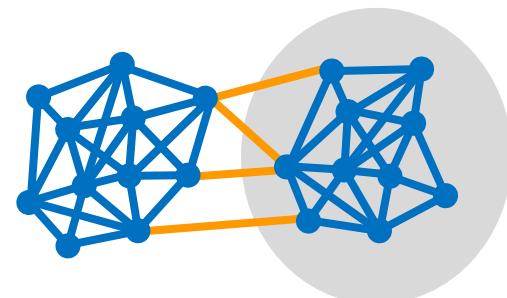
NO

a gap
a few links
 $O(1)$ terms



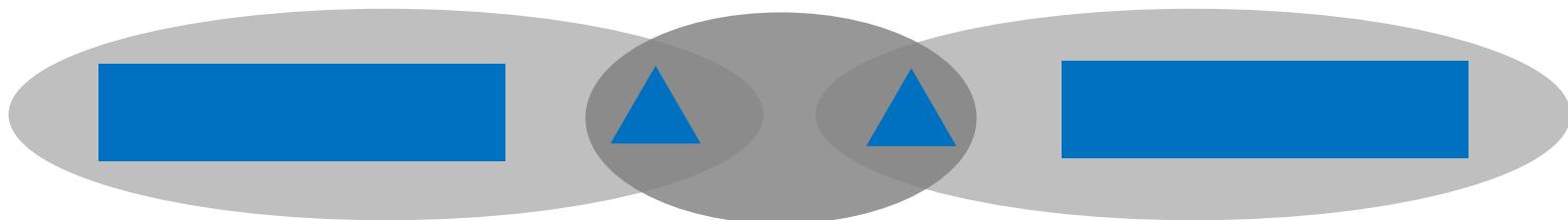
not much
entanglement
(a “simple” ground state)

- generalized area conjecture
entropy \sim cut size



3 Generalized area conjecture: the counterexample

- an $N \times 3 \times 3 \times N$ dimensional system
- a gapped, frustration-free Hamiltonian
- an $O(1)$ interaction between the qutrits



3 Generalized area conjecture: the counterexample

- an $N \times 3 \times 3 \times N$ dimensional system
- a gapped, frustration-free Hamiltonian
an $O(1)$ interaction between the qutrits
- a unique, very entangled ground state
 $O(N)$ entanglement entropy across the cut



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L with ground states

$$\frac{1}{\sqrt{3}}(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$$

- as a vector

x

$Ax \otimes |j\rangle \otimes |y\rangle$

Bx

- as a matrix

X_1

AX_1

BX_1

X_2

AX_2

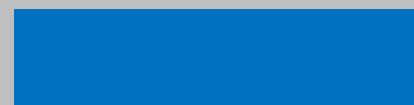
BX_2

X_3

AX_3

BX_3

P_L



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_R
with ground states

$$\frac{1}{\sqrt{3}} (|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle)$$

- as a vector

$$|i\rangle \otimes |x\rangle \otimes \begin{array}{c} y \\ yA \\ yB \end{array}$$

as a matrix

$$\begin{matrix} Y_1 & Y_1A & Y_1B \\ Y_2 & Y_2A & Y_2B \\ Y_3 & Y_3A & Y_3B \end{matrix}$$



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L
- a projector P_R

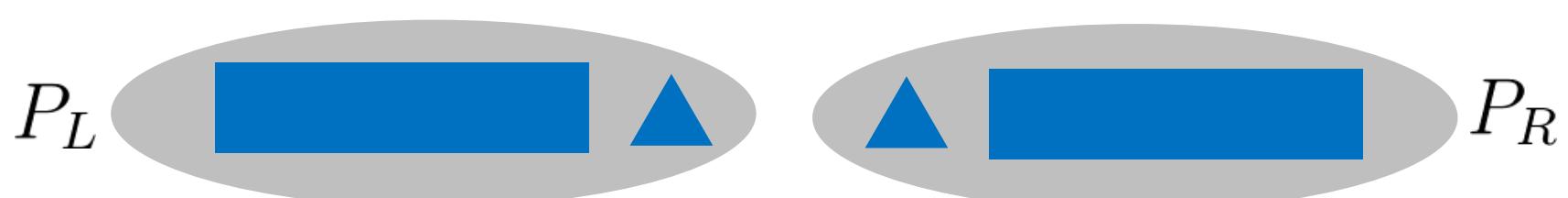
$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

X_1	X_2	X_3
AX_1	AX_2	AX_3
BX_1	BX_2	BX_3

X	XA	XB
AX	AXA	AXB
BX	BXA	BXB

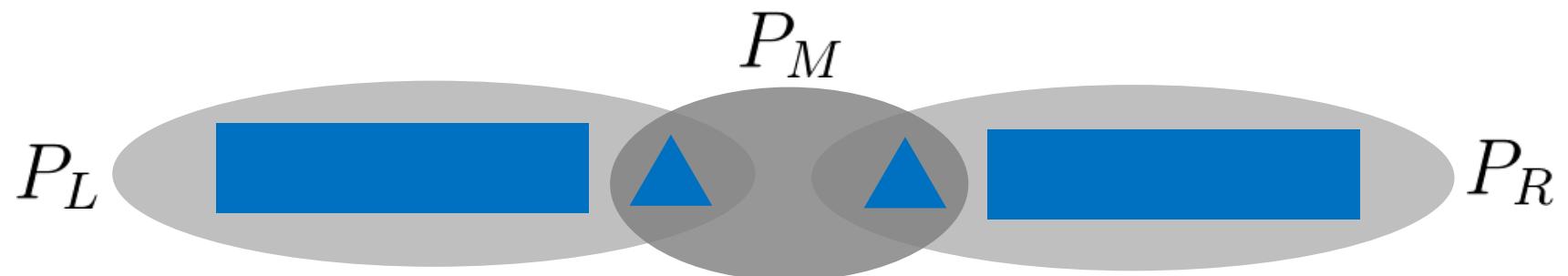
Y_1	Y_1A	Y_1B
Y_2	Y_2A	Y_2B
Y_3	Y_3A	Y_3B



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L $(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$
- a projector P_R $(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$
- a projector P_M

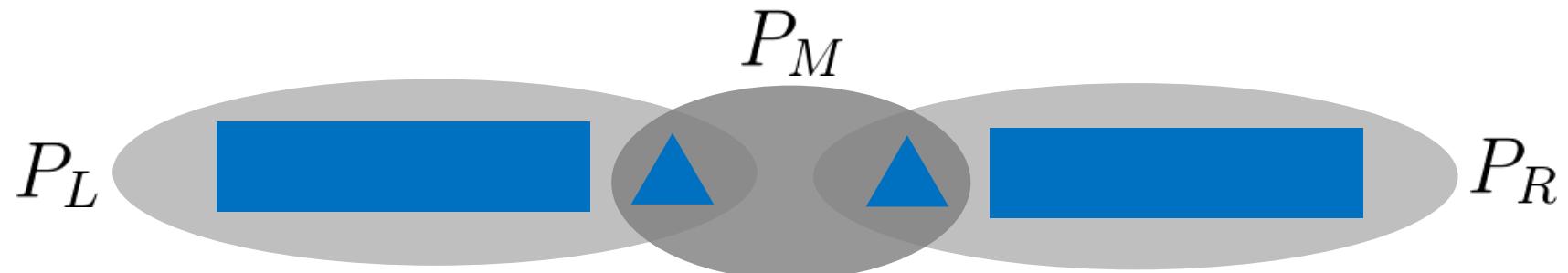
X	XA	XB
AX	AXA	AXB
BX	BXA	BXB



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

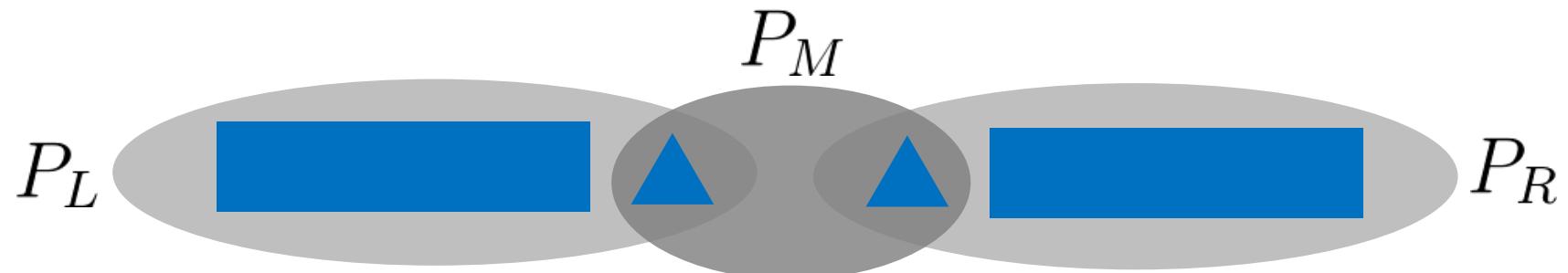
- a projector P_L $(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$
- a projector P_R $(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$
- a projector P_M enforce symmetry: 12 & 21
13 & 31
- who commutes with A and B ?
only the identity, as $[I, A, B]$ is a q. expander

X	XA	XB
AX	AXA	AXB
BX	BXA	BXB



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L $(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$
 - a projector P_R $(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$
 - a projector P_M enforce symmetry: 12 & 21
13 & 31
 - who commutes with A and B ?
only the identity, as $[I, A, B]$ is a q. expander
- | | | |
|-----|------|------|
| I | A | B |
| A | AA | AB |
| B | BA | BB |
- $\frac{1}{3\sqrt{N}}$

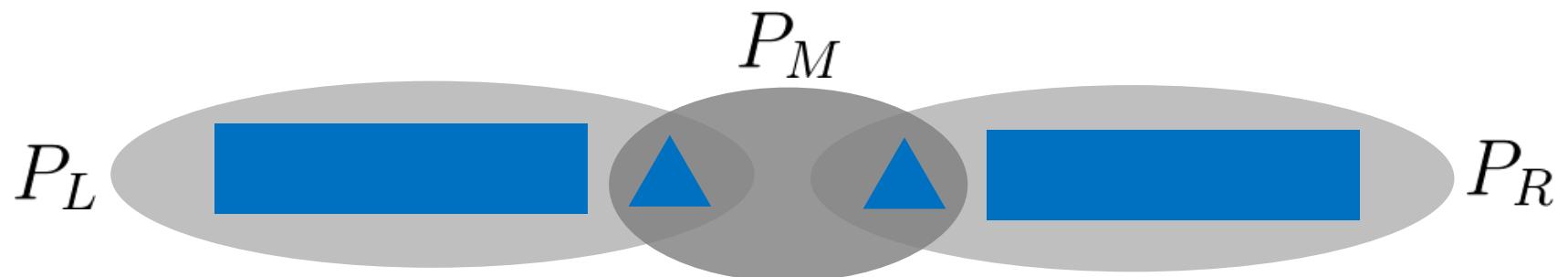


3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

a unique, very entangled ground state

a gapped Hamiltonian, $O(1)$ terms, frust. free

$$\begin{matrix} & I & A & B \\ A & & AA & AB \\ B & & BA & BB \end{matrix} \quad \frac{1}{3\sqrt{N}}$$



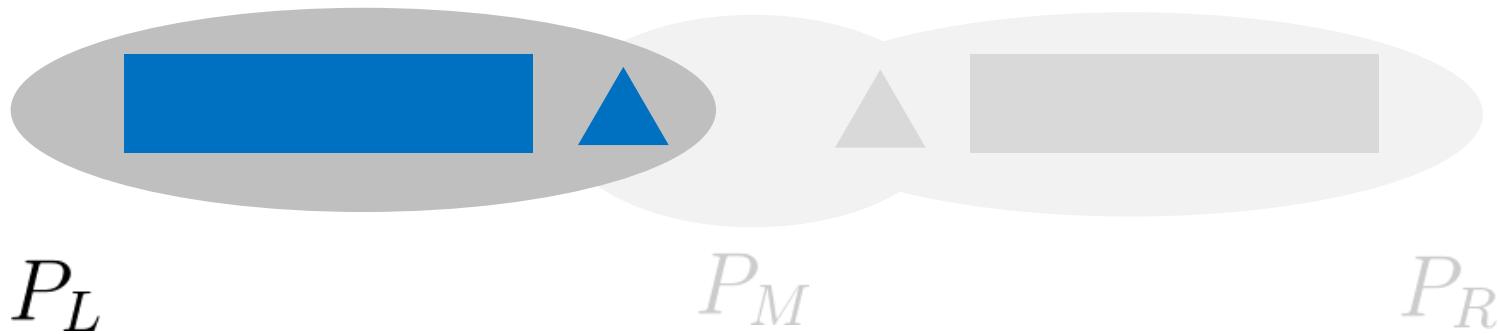
3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits



prepare $\frac{1}{\sqrt{3}} (|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$

from $\frac{1}{\sqrt{3}} (|1\rangle + |2\rangle + |3\rangle) |x\rangle$



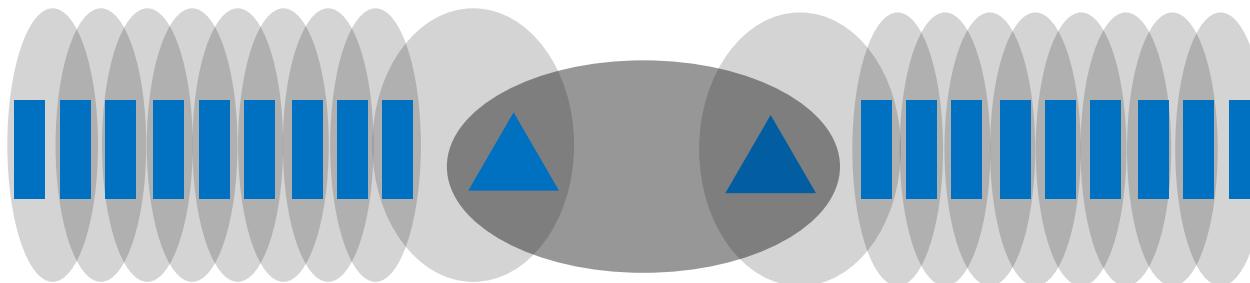
3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits



prepare $\frac{1}{\sqrt{3}} (|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$

from $\frac{1}{\sqrt{3}} (|1\rangle + |2\rangle + |3\rangle) |x\rangle$



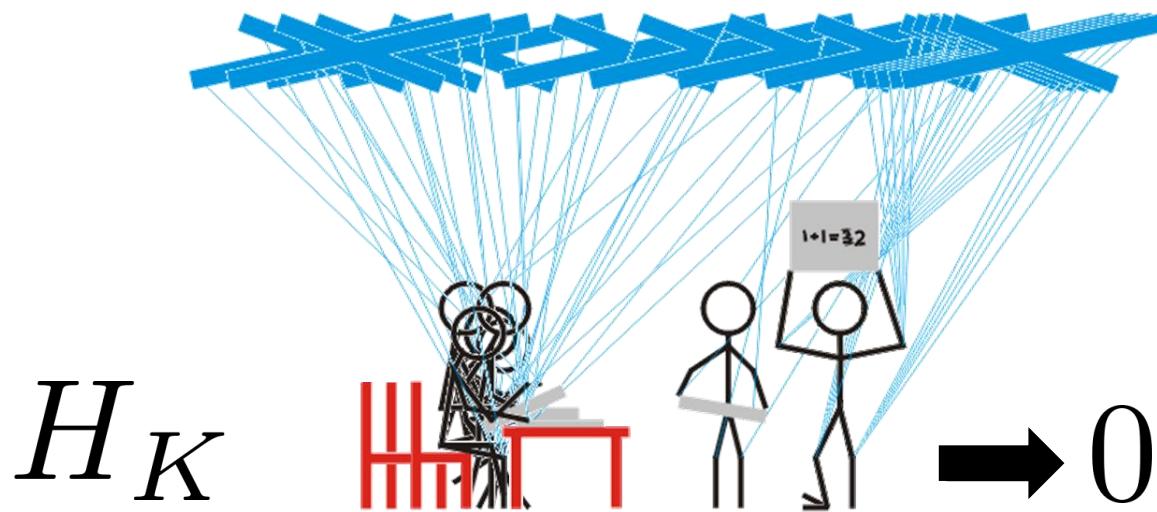
P_L

P_M

P_R

3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits



$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

$\overbrace{\quad\quad\quad\quad\quad\quad\quad\quad\quad}^{U_t \cdots U_1 |\varphi_0\rangle}$

3 The history state: a ground state of a qudit chain

2-local

c-o-n-d-i-t-i-o-n-s

d=13 [AGIK'08]

d=8 [HNN'13]

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$



3 The history state: a ground state of a qudit chain

2-local

c-o-n-d-i-t-i-o-n-s

$$| \cdots \rangle |0\rangle \otimes |0\rangle$$

$$|\varphi_t\rangle \otimes |t\rangle$$
$$|\varphi_{t+1}\rangle \otimes |t+1\rangle$$

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

clock encoding

state progression

initialization

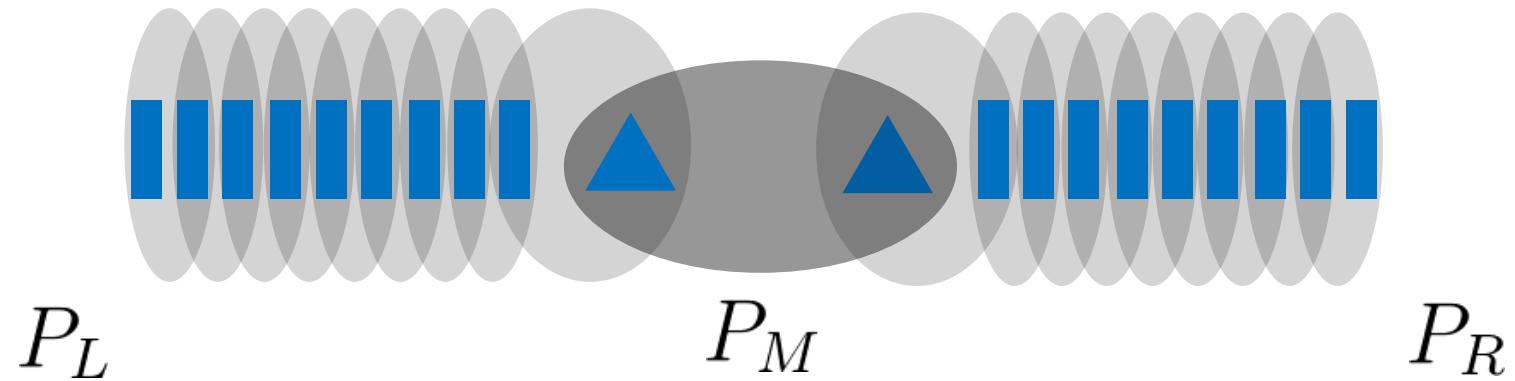
idling



most of the state has the result

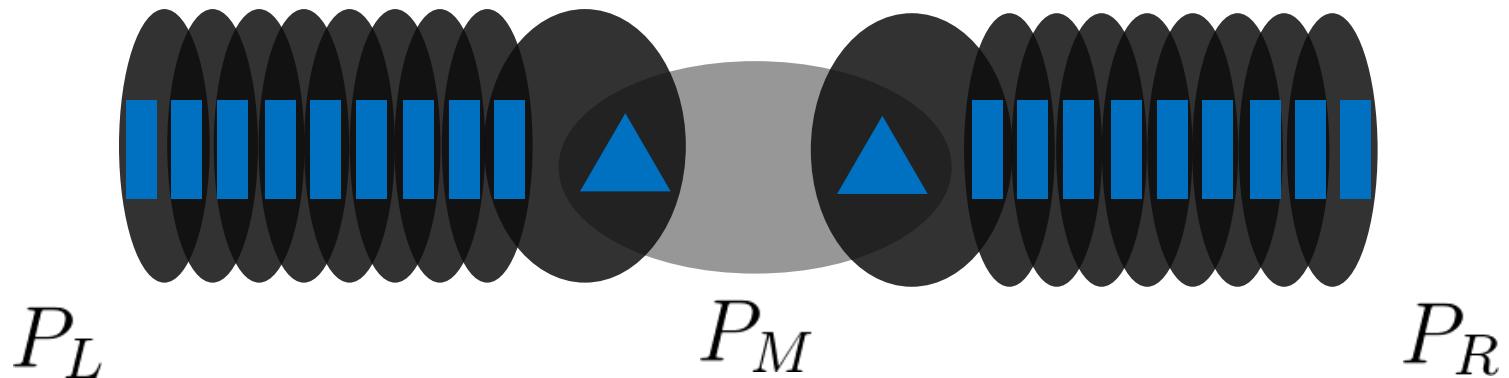
3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits
an approx. groundstate, a small $1/\text{poly}(n)$ gap



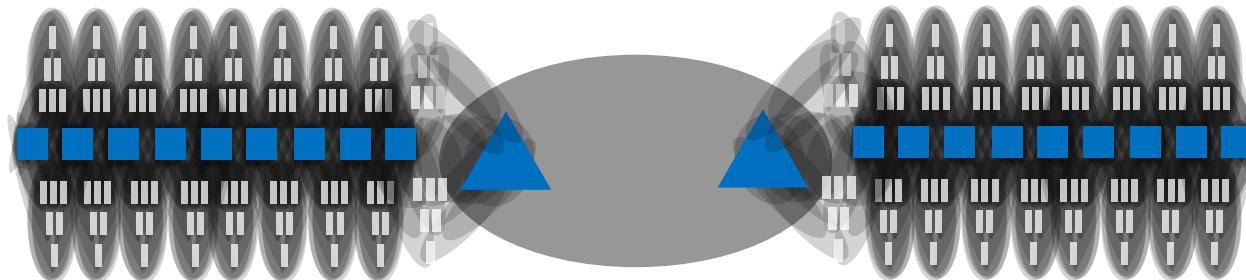
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- rescale P_L, P_R (not the middle!)
a constant gap, huge couplings

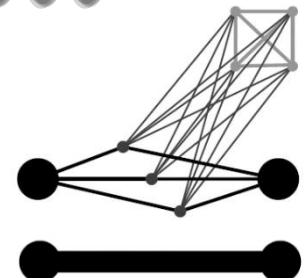


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- decompose using gadgets [Cao, N.]
~~huge couplings~~ many spins, high degree

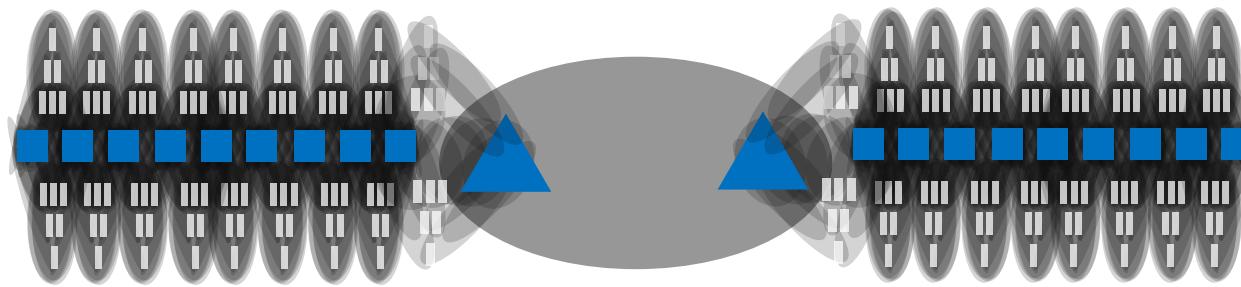


3 A local Hamiltonian, $(N \times M) \times 3 \times 3 \times (N \times M)$

- a unique and very entangled ground state

$$\approx |w\rangle \otimes$$

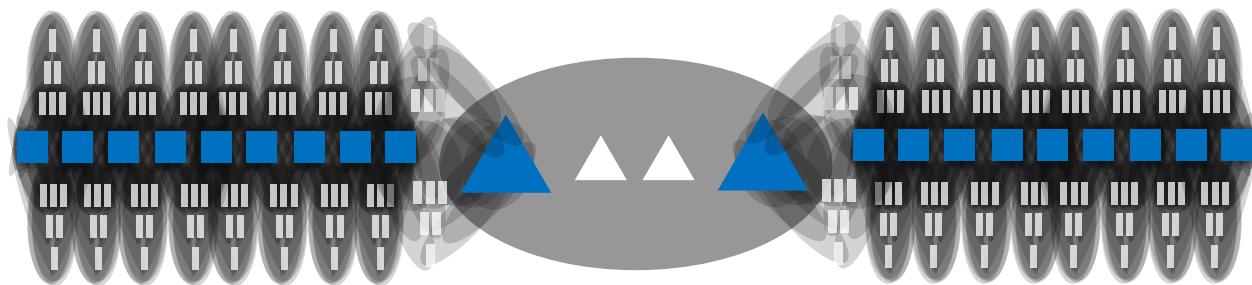
I	A	B
A	AA	AB
B	BA	BB



- a constant gap $O(1)$ or smaller norm terms, frustration

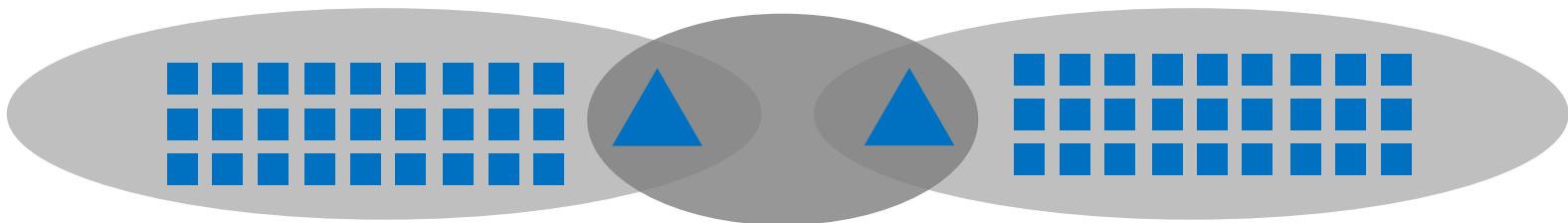
3 What next?

- a longer middle chain?



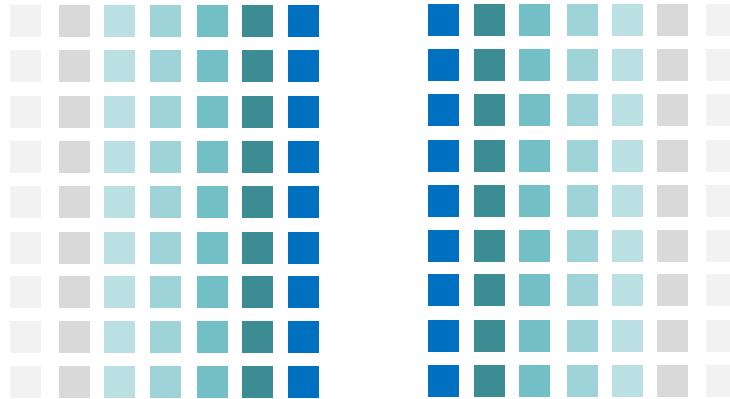
3 What next?

- a longer middle chain?
- a nice lattice on the sides?



3 What next?

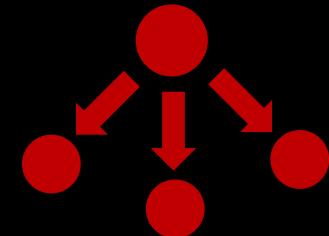
- a longer middle chain?
- a nice lattice on the sides?
- a new area conjecture: count the cut & things nearby?



1

q. expanders

maximally entangled states



2

entanglement

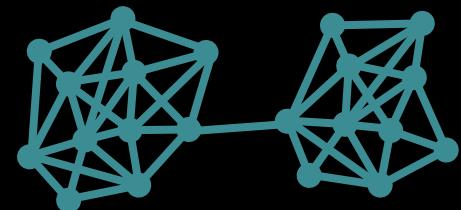
testing and communication



3

area law

gaps, connections, correlations



local tests of global entanglement and a counterexample to the generalized area law

