

An abstract geometric composition featuring a teal background. Two large yellow shapes, one on the left and one on the right, are separated by a central teal channel. Two black diagonal bars cross this channel. A small white circle with the number '1' is positioned in the center of the teal channel.

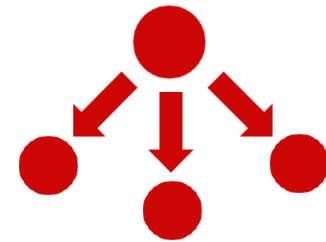
1



1

q. expanders

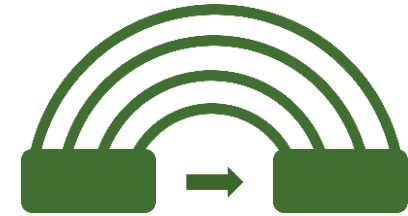
maximally entangled states



2

entanglement

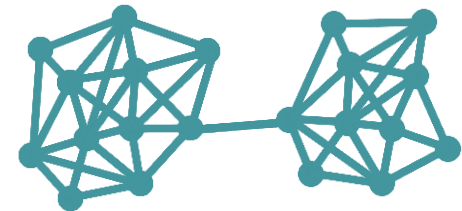
testing and communication

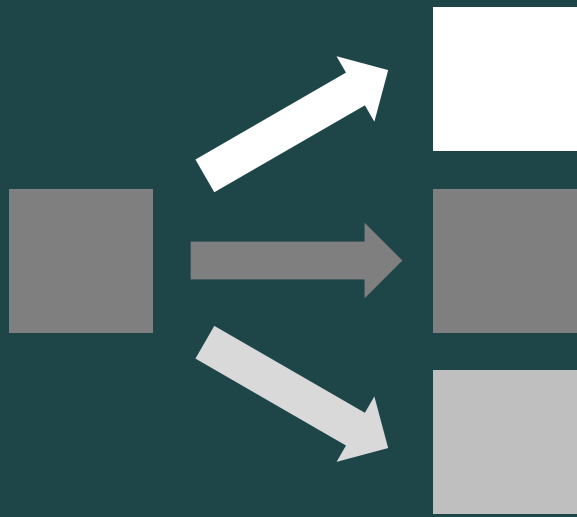


3

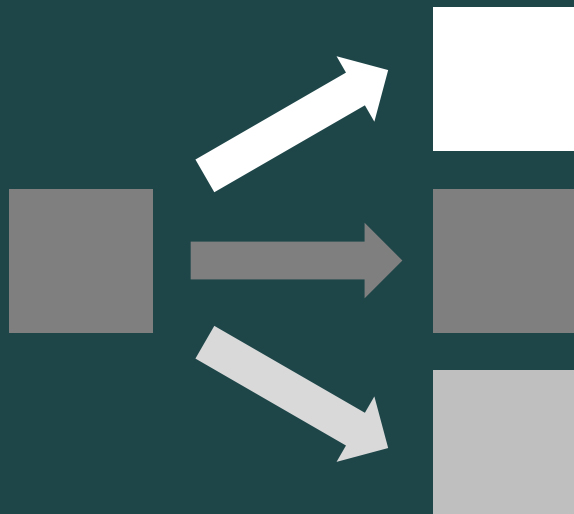
area law

gaps, connections, correlations





Quantum
Expanders



Expanders Everywhere!

1-5 December 2014

Neuchâtel

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Ana Khukhro, *Université de Neuchâtel*
Alain Valette, *Université de Neuchâtel*

<https://sites.google.com/site/expanderseverywhere/>

(part of the Swiss Doctoral Program in Mathematics supported by CUSO)

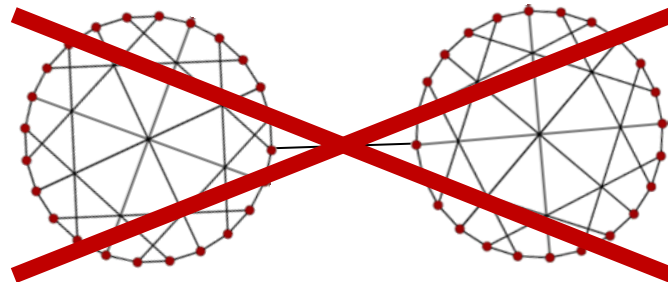
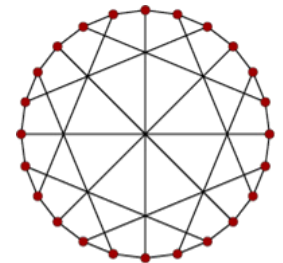
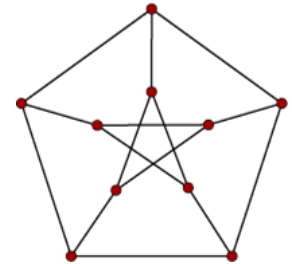
1 Classical expanders

- Aram Harrow's talk, QHC workshop at the [youtube Harrow quantum expanders]



- graphs that “mix” well
divide in two? cut a lot (fraction) of edges!

examples: Cayley graphs

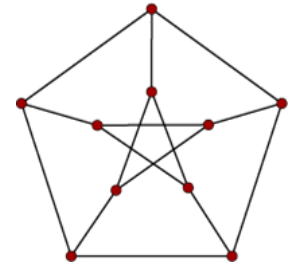


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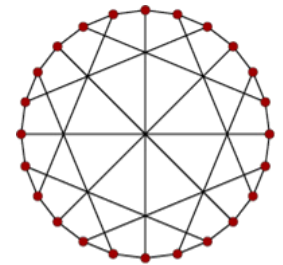


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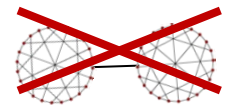


examples: Cayley graphs

- normalized adjacency matrix
second largest eigenvalue $1-\lambda$



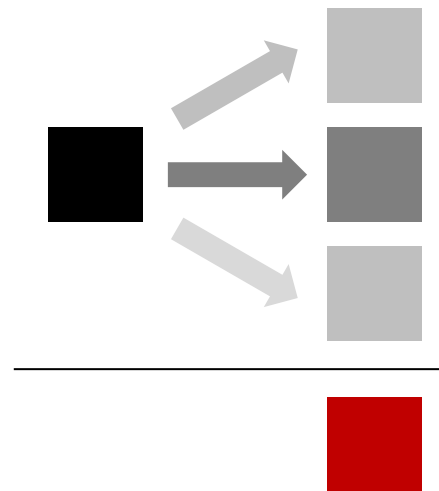
- classical expanders: constant-degree
approximations to the full graph



1 Mixing up something quantum

- applying random unitaries

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

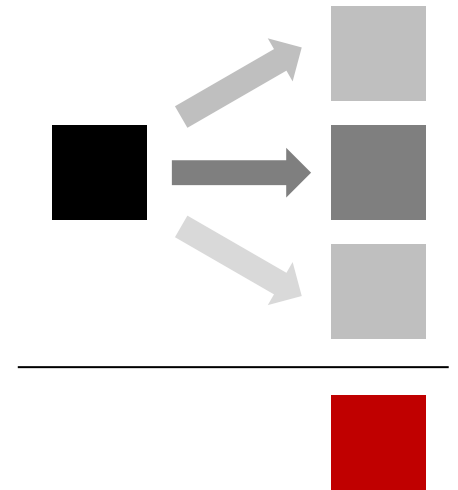


- classical expanders: constant-degree approximations to the full graph

1 Quantum expanders

- applying random unitaries from a set a discrete approximation to the Haar measure

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

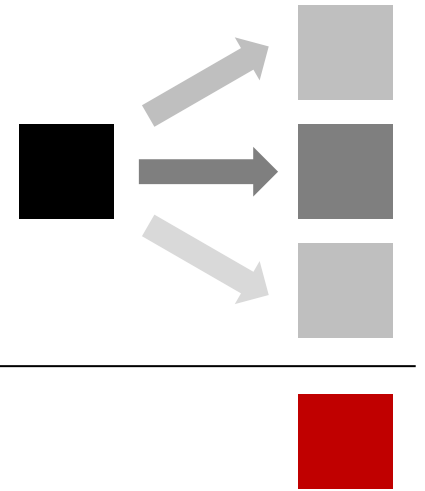


- classical expanders: constant-degree approximations to the full graph

1 Quantum expanders

- applying random unitaries from a set
a discrete approximation to the Haar measure
- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$



- a small second largest singular value λ
not far from the depolarizing channel $\|\hat{\mathcal{E}} - |\phi_D\rangle\langle\phi_D|\| = \lambda$
- q. expander constructions, also for fixed k ($=8$. $=3$?)
[Ben-Aroya+ '07, Hastings '07, Gross & Eisert '08, Hastings & Harrow '09]

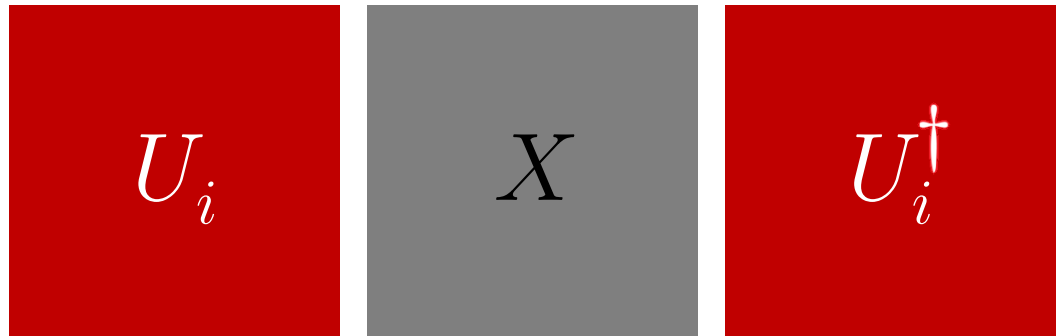
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

- a matrix that doesn't change?

$$X = \mathbb{I}$$



$$U_i X = X U_i$$



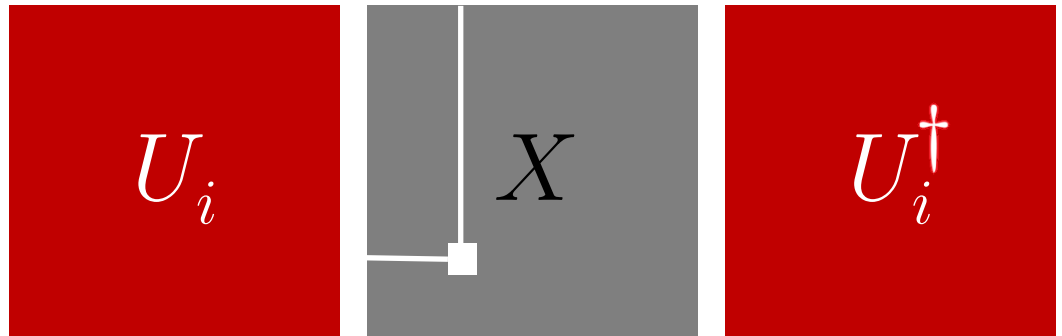
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- interpreting matrices as 2-register states

$$\sum_{a,b} X_{ab} |a\rangle\langle b|$$

density matrix

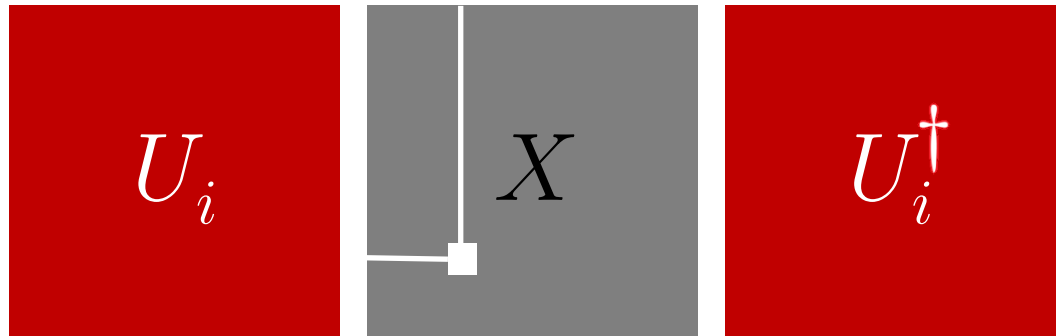
1 Quantum expanders

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- interpreting matrices as 2-register states

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle |b\rangle$$

pure state
↓

- distributively applying an expander stationary?

$$\sum_{a,b} X_{ab} |a\rangle |b\rangle$$

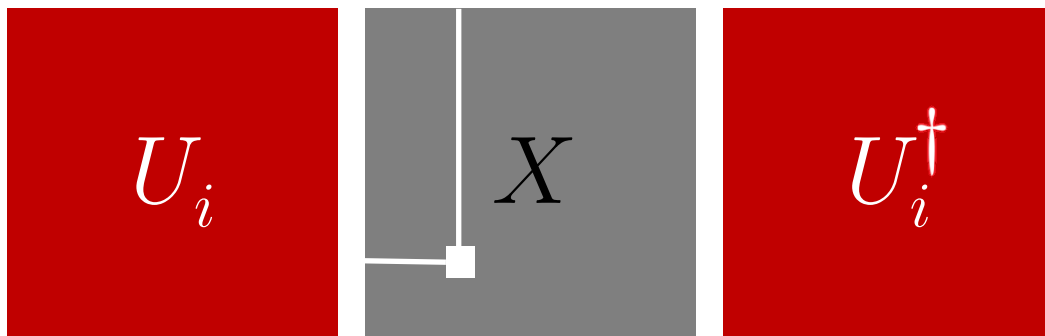
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- interpreting matrices as 2-register states

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle |b\rangle$$

- distributively applying an expander

stationary? $X=I$... max. entangled! $|\Phi_N\rangle = \frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$



local tests of global entanglement

2 EPR testing

- how costly is it to certify that we share a maximally entangled state?

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$

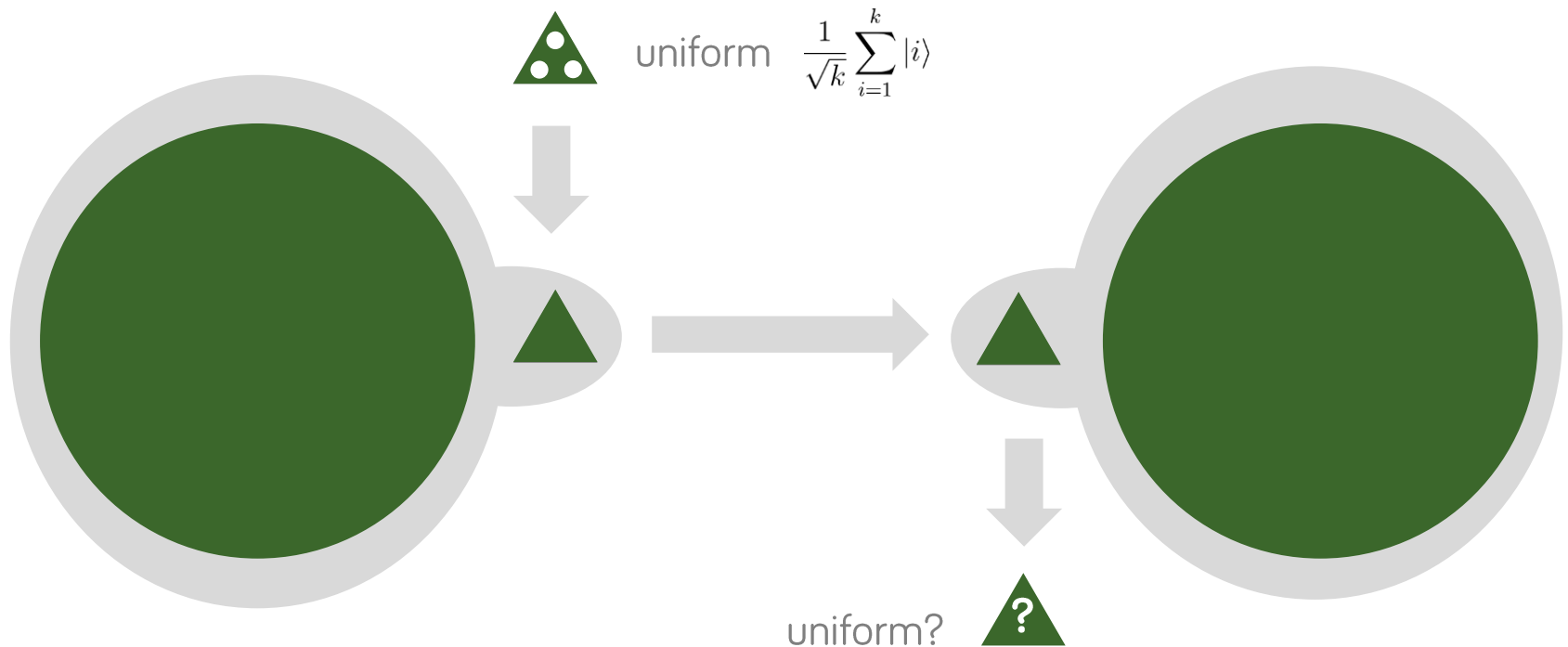


2 EPR testing

- how costly is it to certify that we share a maximally entangled state?
- apply a quantum expander distributively

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$

$$U_i \otimes U_i^*$$

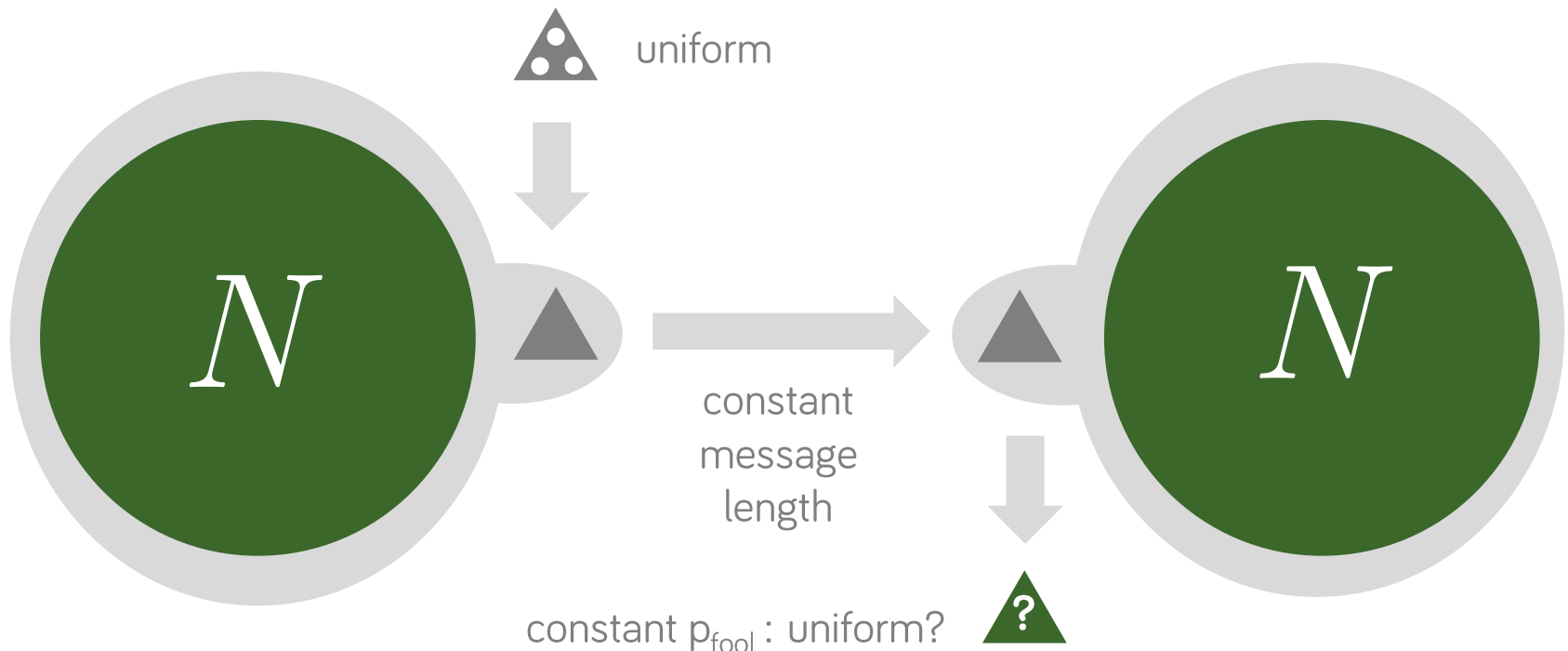


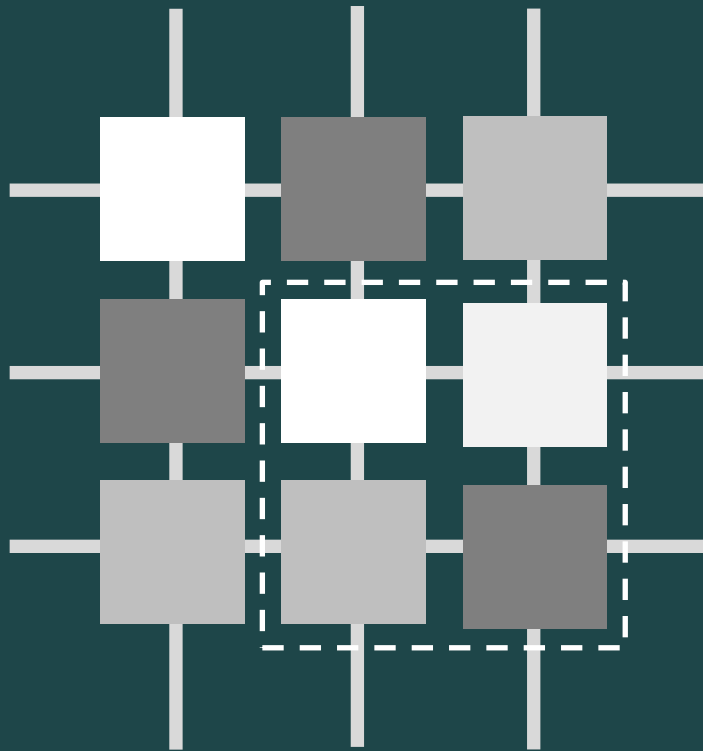
2 EPR testing

- action on states
- does the qutrit remain uniform?
when X commutes with all U_i
- quantum expander property ... soundness

$$\frac{1}{\sqrt{k}} \sum_{i=1}^k |i\rangle (U_i \otimes U_i^*) |X\rangle$$

$$\sum_{a,b} X_{ab} |a\rangle |b\rangle$$



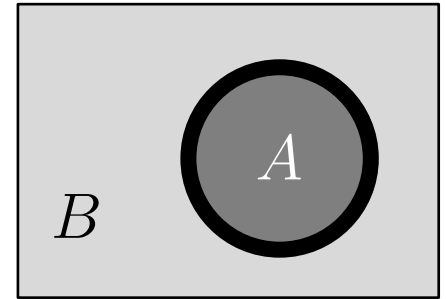


a counterexample to the
generalized
area law

3 Area law: ground states of (gapped) q. spin systems

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \begin{array}{l} \text{volume} \\ \text{area} \end{array}$$



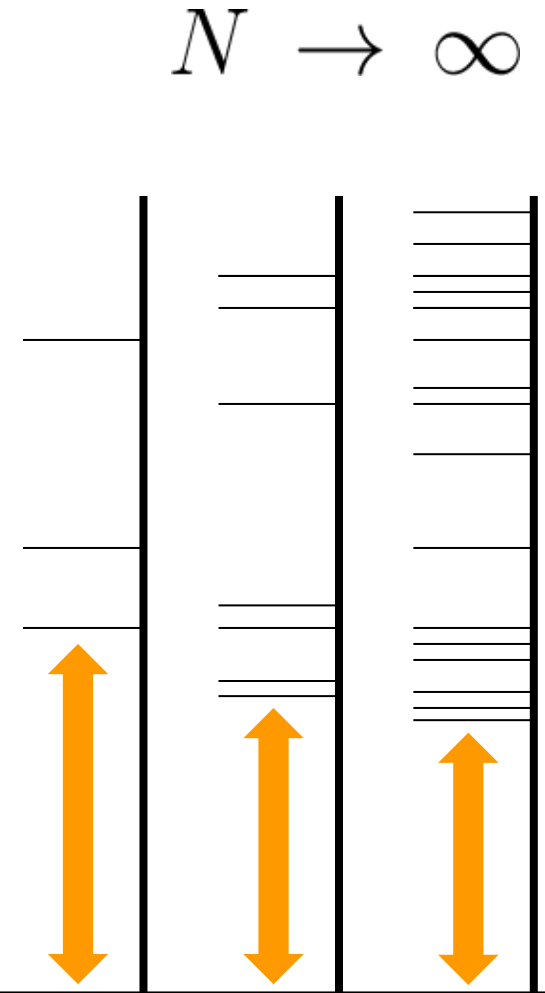
$$\rho_A = \text{Tr}_B \rho$$

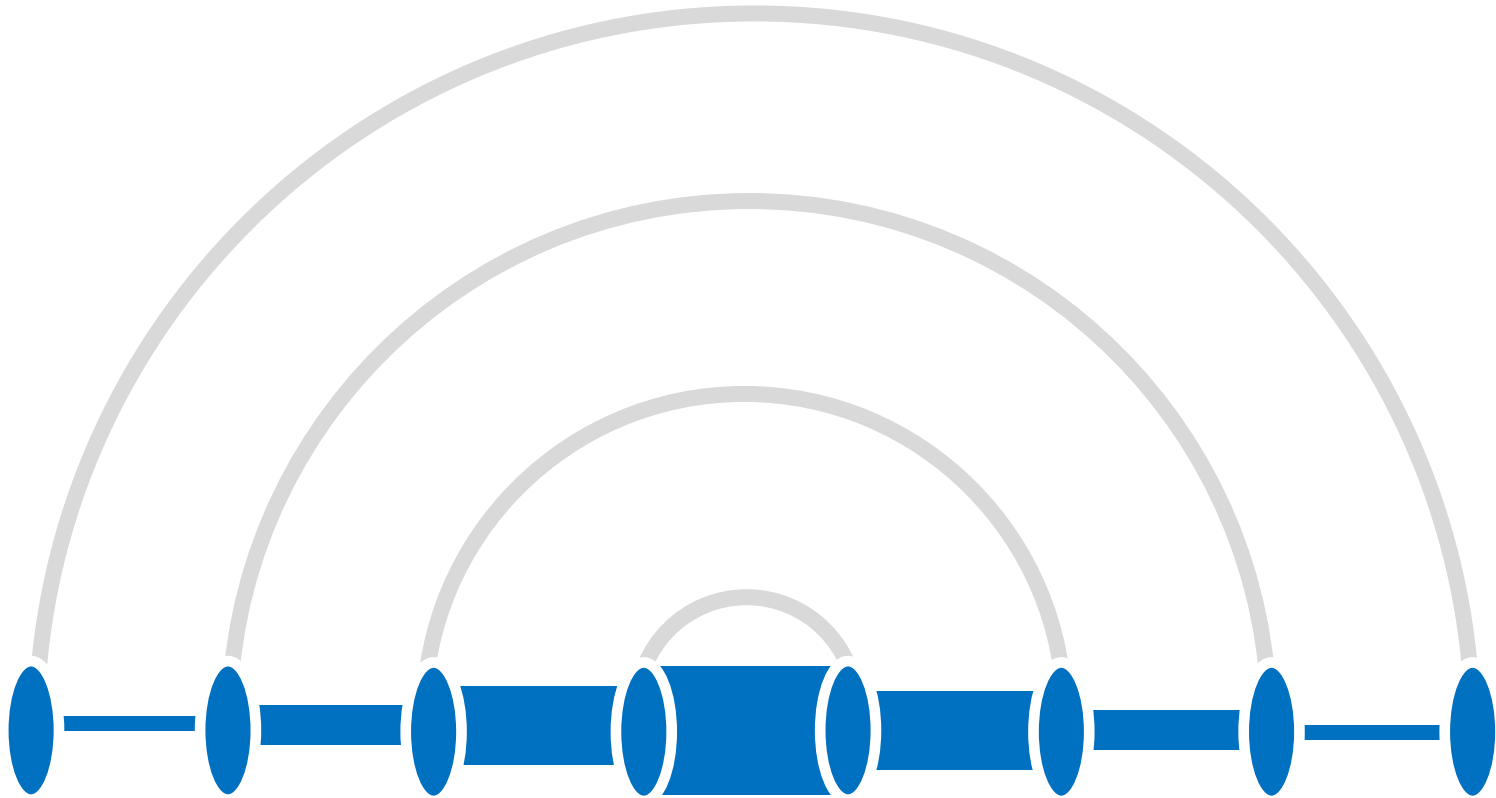
a gapped system ... a simple ground state?

3 Gapped Hamiltonians

- Are there any states close to the ground state when we take the thermodynamic limit?
- local, $O(1)$ norm terms

~~an inverse-poly gap? $\Delta = \frac{c}{N} \rightarrow 0$~~





- without a gap, the entropy can be large

[Verstraete, Latorre+]

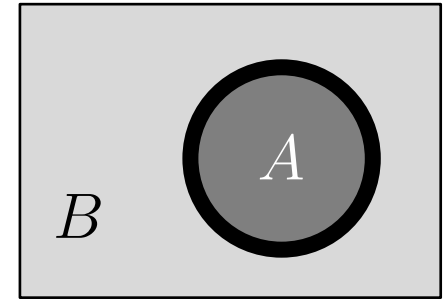
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area

Schmidt coeff's fall off



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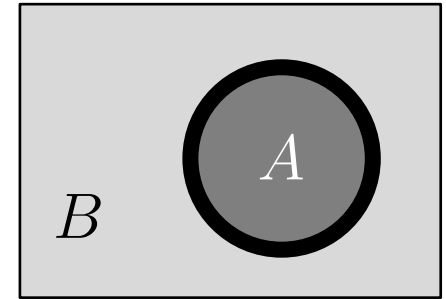
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area

Schmidt coeff's fall off



$$\rho_A = \text{Tr}_B \rho$$

- 1D ... algorithms [White 92, Vidal 03, Landau+ 13]

theorems [Hastings 07, Arad+ 13]

- 2D ... we're close

small gap? large local dimension?

a gapped system ... a simple ground state?

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- entanglement entropy

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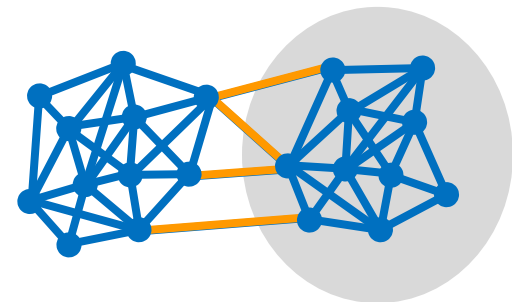
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- 2D ... we're close

small gap? large local dimension?

- generalized area conjecture
entropy \sim cut size



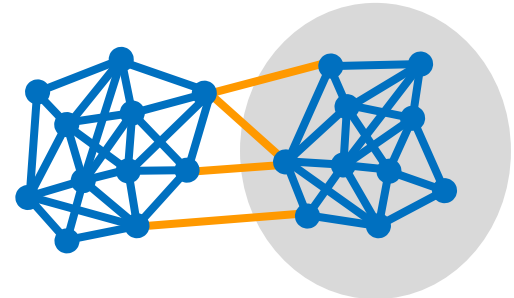
Not true.

a gap
a few links
 $O(1)$ terms



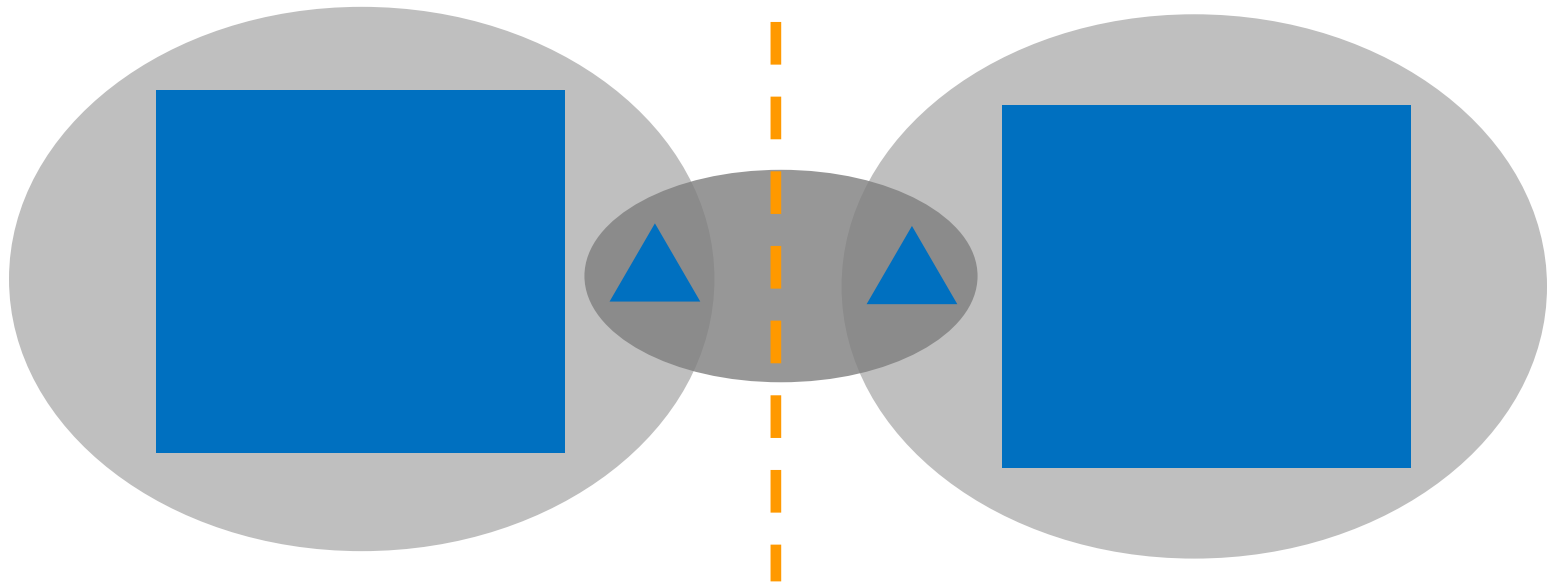
not much
entanglement
(a “simple” ground state)

- generalized area conjecture
entropy \sim cut size



3 Our counterexample to the generalized area conjecture

- an $N \times 3 \times 3 \times N$ dimensional system
- a frustration-free, gapped, Hamiltonian
- a single $O(1)$ interaction of two 3×3 subsystems
- a unique, very entangled ground state with $O(N)$ entanglement entropy across the cut

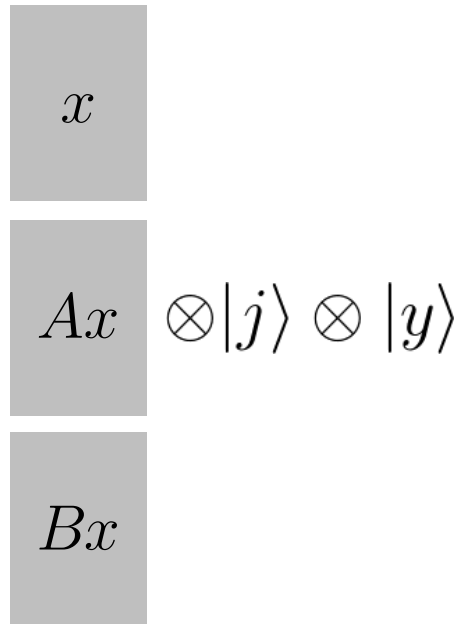


3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

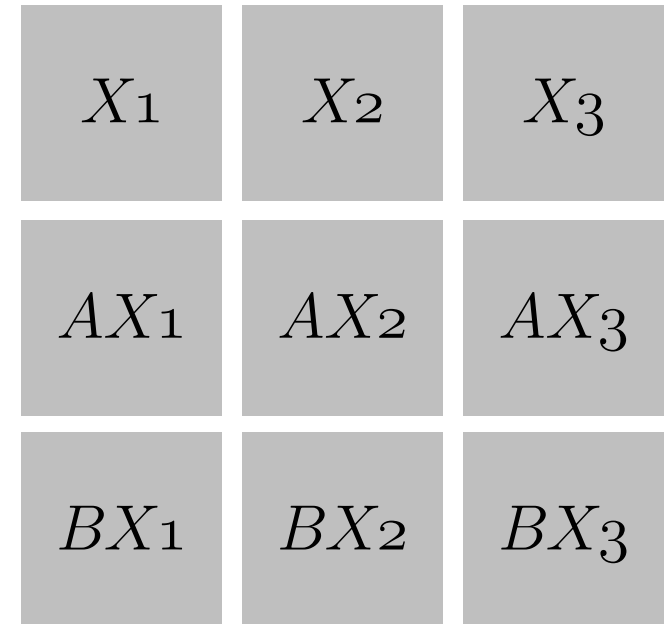
- a projector P_L with ground states

$$\frac{1}{\sqrt{3}} (|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$$

- as a vector



- as a matrix



3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

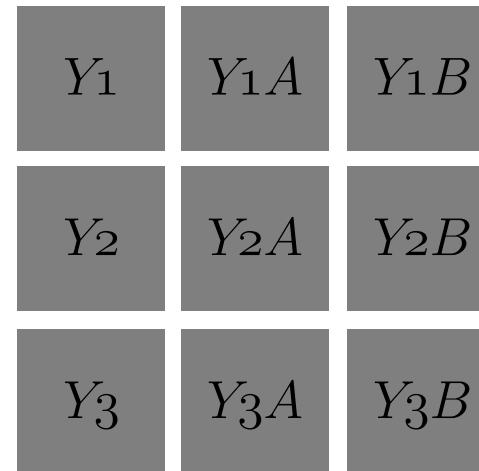
- a projector P_R
with ground states

$$\frac{1}{\sqrt{3}} (|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle)$$

- as a vector $|i\rangle \otimes |x\rangle \otimes$



as a matrix

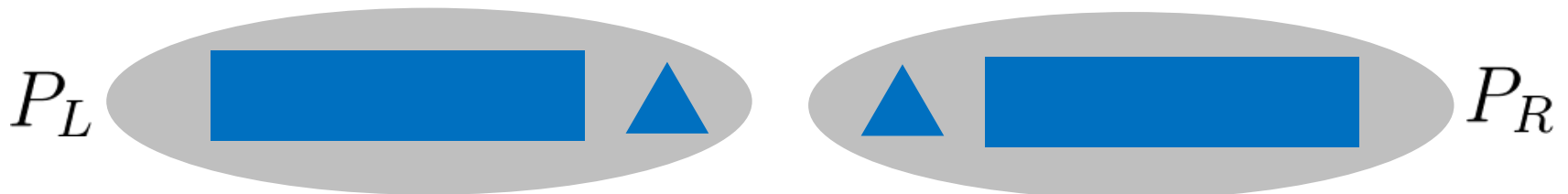
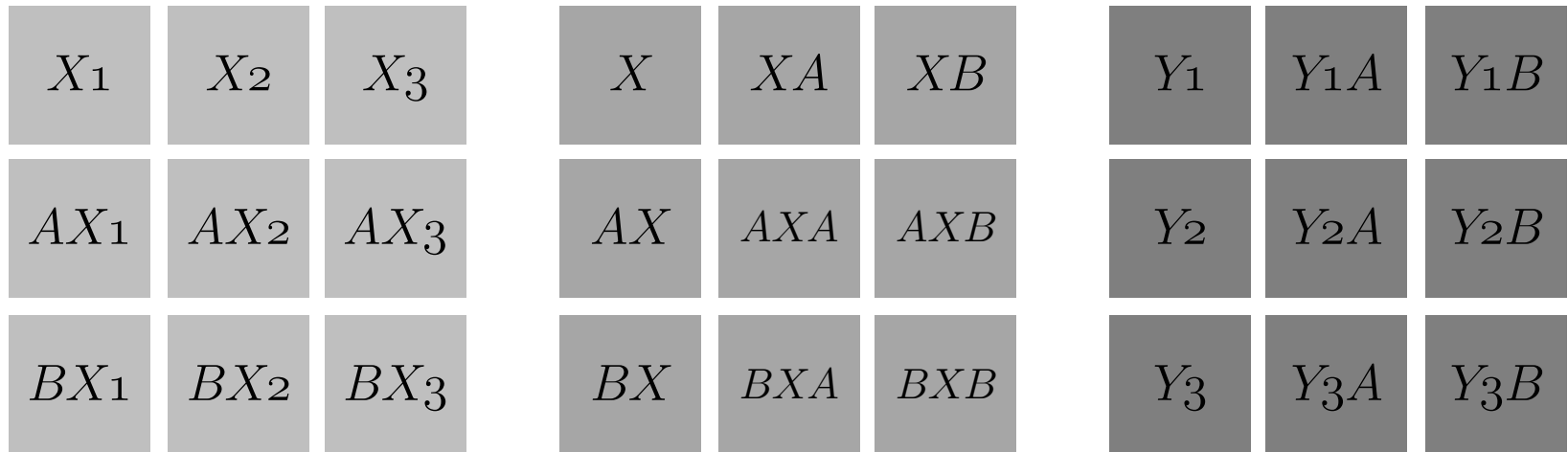


3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

- a projector P_L
- a projector P_R

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$



3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

■ a projector P_L

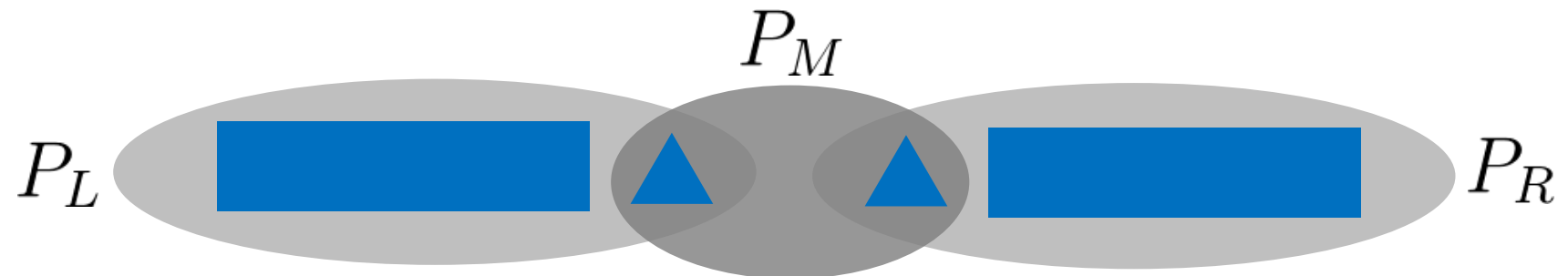
a projector P_R

a projector P_M

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

X	XA	XB
AX	AXA	AXB
BX	BXA	BXB



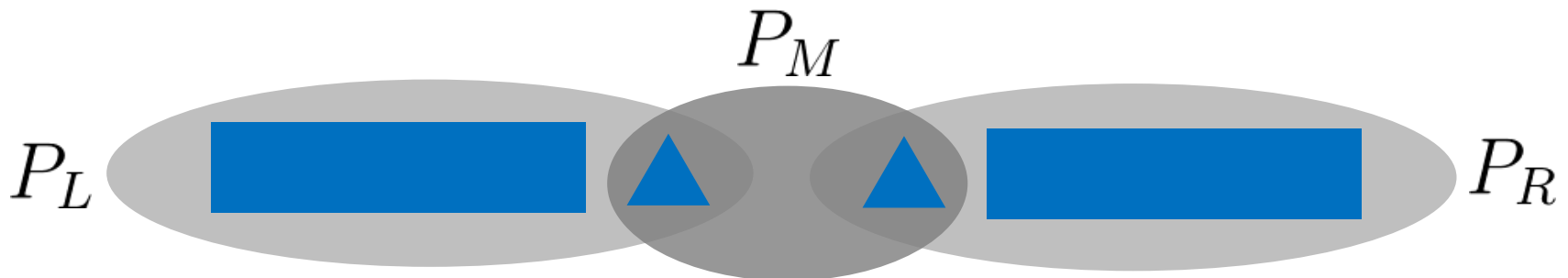
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- a projector P_R $(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$
- a projector P_M enforce symmetry

- who commutes with A and B ?
 only the identity,
 as $[I, A, B]$ are
 a q. expander

X	XA	XB
AX	AXA	AXB
BX	BXA	BXB

for 12 & 21
 for 13 & 31



3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

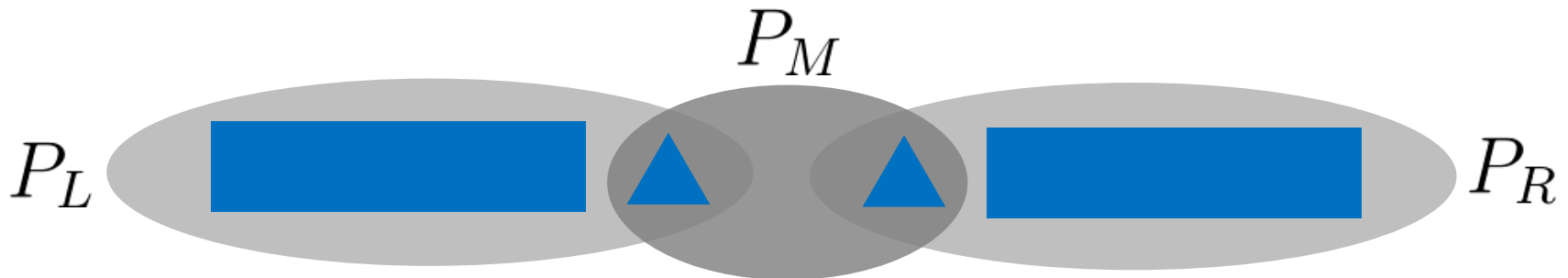
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 a projector P_M enforce symmetry

- who commutes with A and B ?
 only the identity,
 as $[I, A, B]$ are
 a q. expander

I	A	B
A	AA	AB
B	BA	BB

for 12 & 21
for 13 & 31

$$\frac{1}{3\sqrt{N}}$$



3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

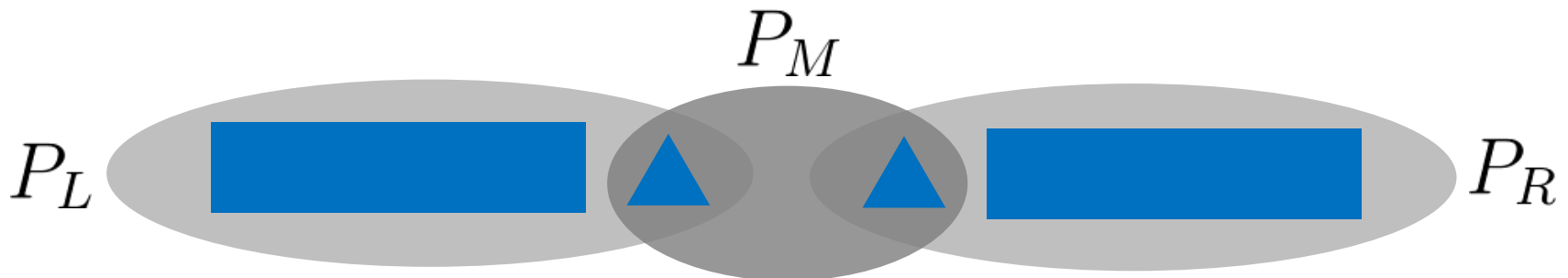
- ground state: unique
very entangled

- Hamiltonian: frustration free
gapped


I	A	B
A	AA	AB
B	BA	BB

$$\sum_{x=1}^N |x\rangle|x\rangle$$

$$\frac{1}{3\sqrt{N}}$$

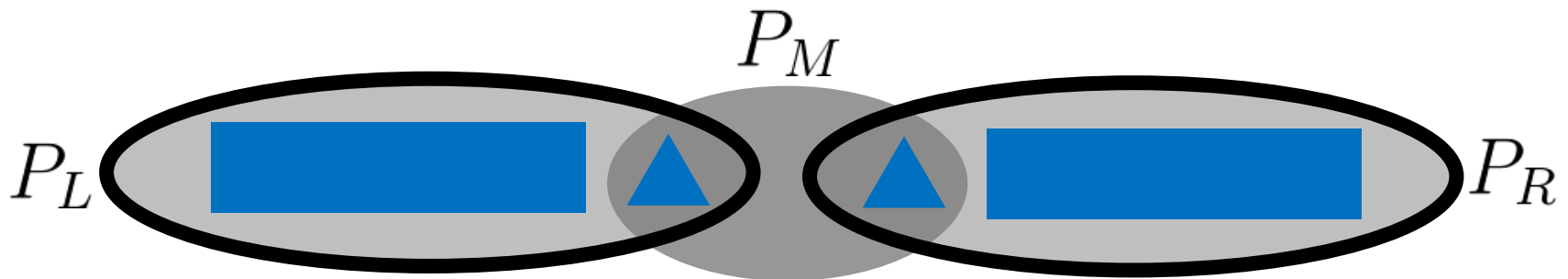


3 Making the counterexample local

- quantum expander $[I, A, B]$... quantum circuits ...
~~nonlocal projectors~~ ... Kitaev's LH & history states 
an approximate groundstate, a very small gap

- rescale P_L, P_R (not the middle!)
huge, nonphysical couplings

- use new “strengthening gadgets” [N., Cao] 
~~large interaction strength~~ ... extra particles, high degree



3 A local Hamiltonian, $(N+n) \times 3 \times 3 \times (N+n)$

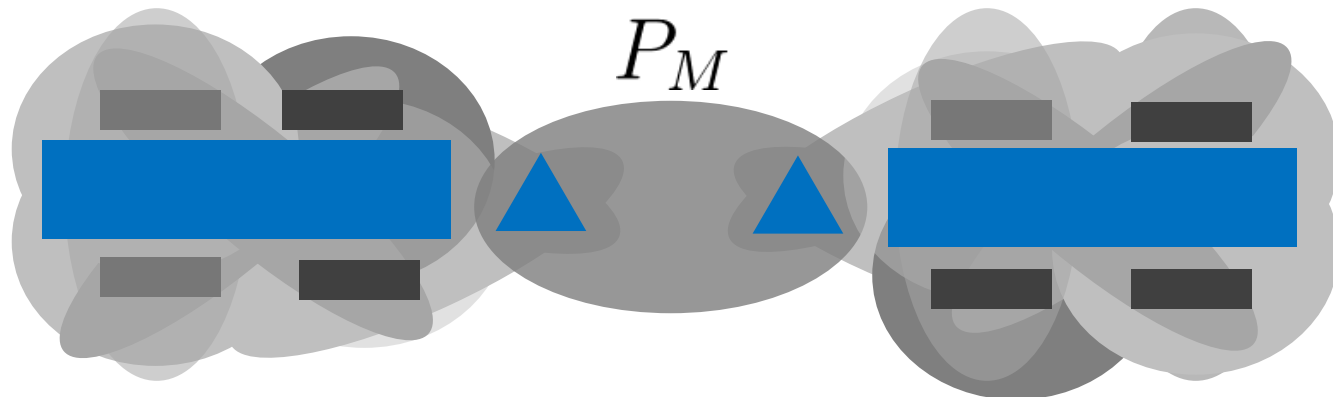
- frustrated, but still gapped

$O(1)$ norm terms

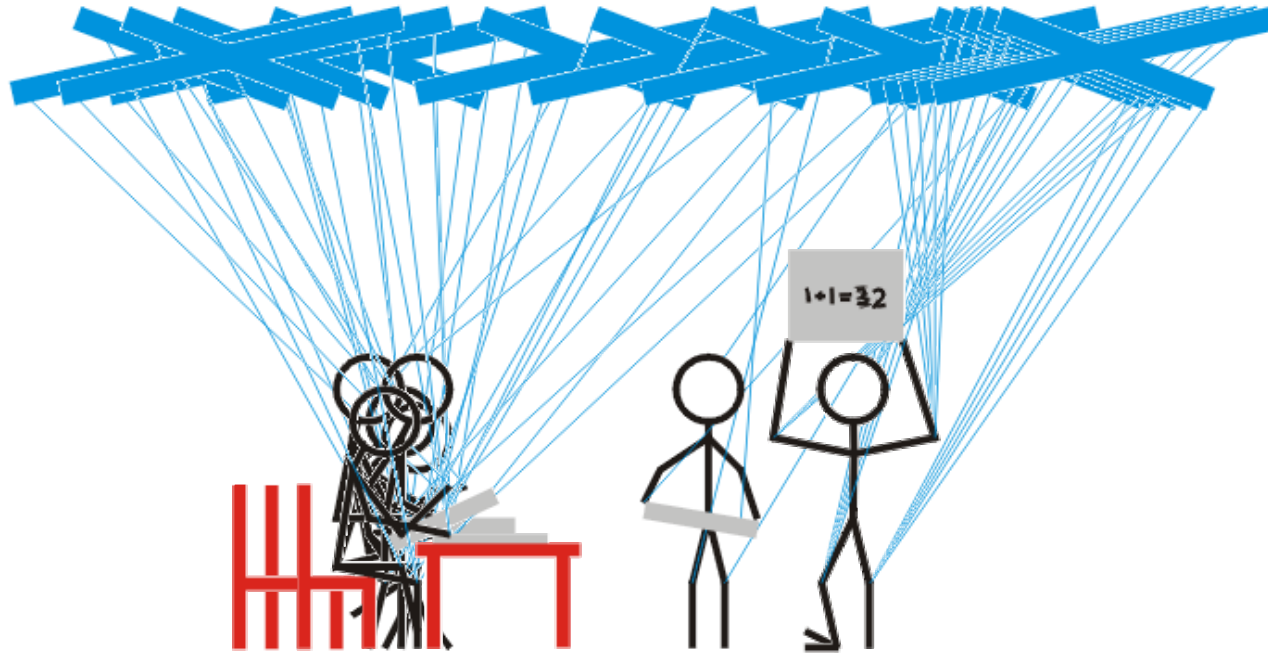
- a unique and still very entangled ground state

$$\approx |w\rangle \otimes$$

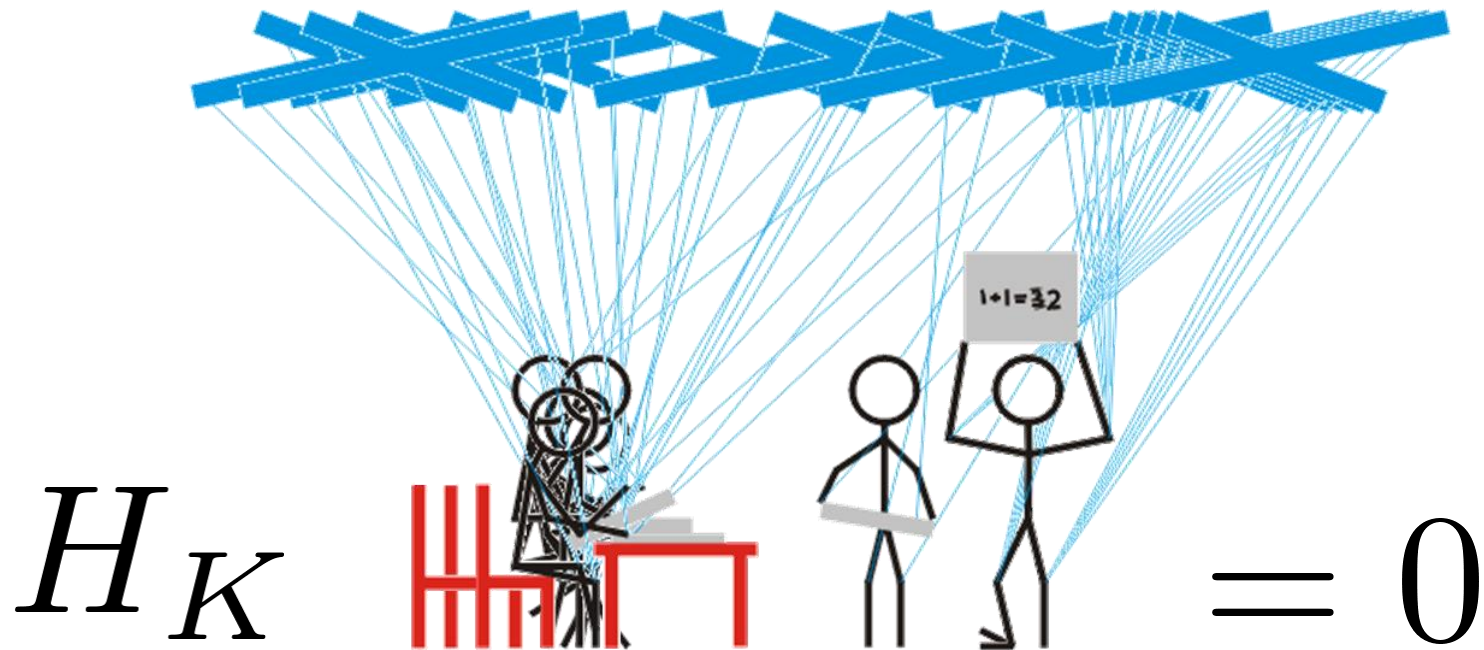
<i>I</i>	<i>A</i>	<i>B</i>
<i>A</i>	<i>AA</i>	<i>AB</i>
<i>B</i>	<i>BA</i>	<i>BB</i>



3 Implementing circuits locally: Feynman's computer



3 The history state: a ground state



$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

$\underbrace{\hspace{10em}}_{U_t \cdots U_1 |\varphi_0\rangle}$

3 The history state: a ground state

k-local

c-o-n-d-i-t-i-o-n-s

clock encoding
state progression
initialization

$$|\cdots\rangle|u\rangle \otimes |0\rangle$$

$$|\varphi_t\rangle \otimes |t\rangle$$

$$|\varphi_{t+1}\rangle \otimes |t+1\rangle$$

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

idling



most of the state has the result

3 A local Hamiltonian, $(N+n) \times 3 \times 3 \times (N+n)$

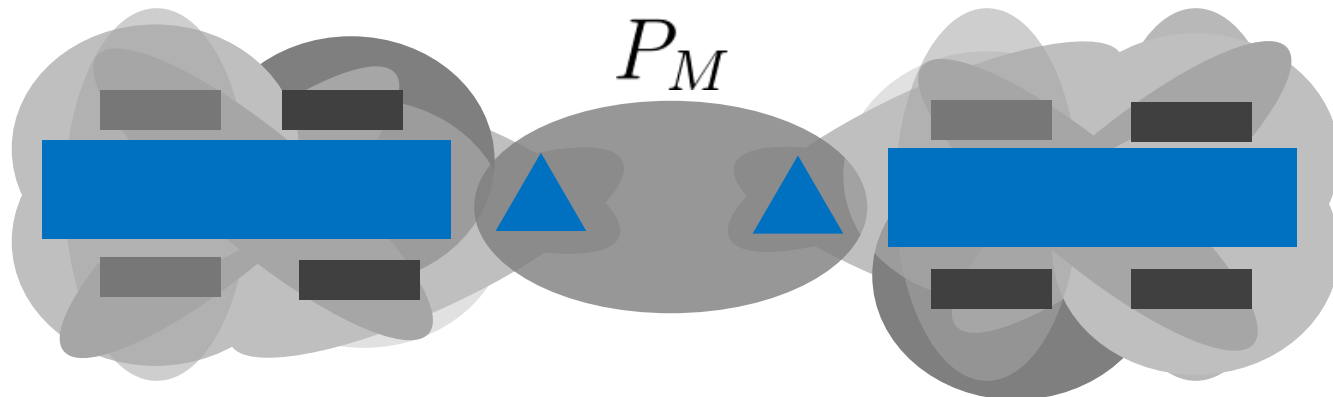
- frustrated, but still gapped

$O(1)$ norm terms

- a unique and still very entangled ground state

$$\approx |w\rangle \otimes$$

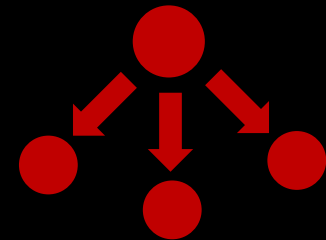
<i>I</i>	<i>A</i>	<i>B</i>
<i>A</i>	<i>AA</i>	<i>AB</i>
<i>B</i>	<i>BA</i>	<i>BB</i>



1

q. expanders

maximally entangled states



2

entanglement

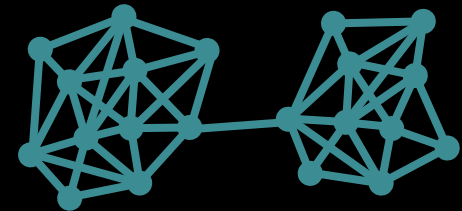
testing and communication



3

area law

gaps, connections, correlations



local tests of global entanglement and a counterexample to the generalized area law



