



# Local tests of global entanglement and a counterexample to the generalized area law.

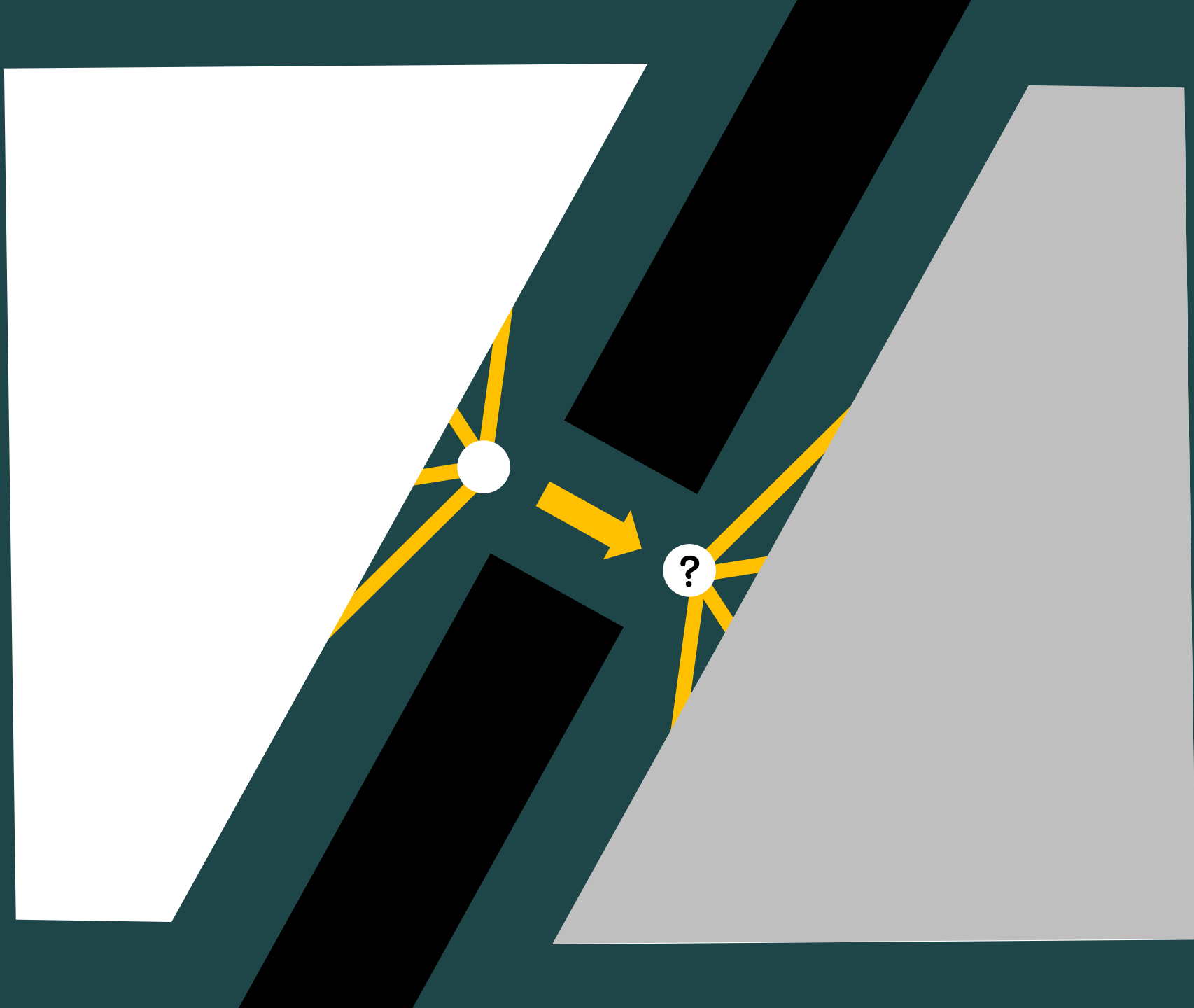


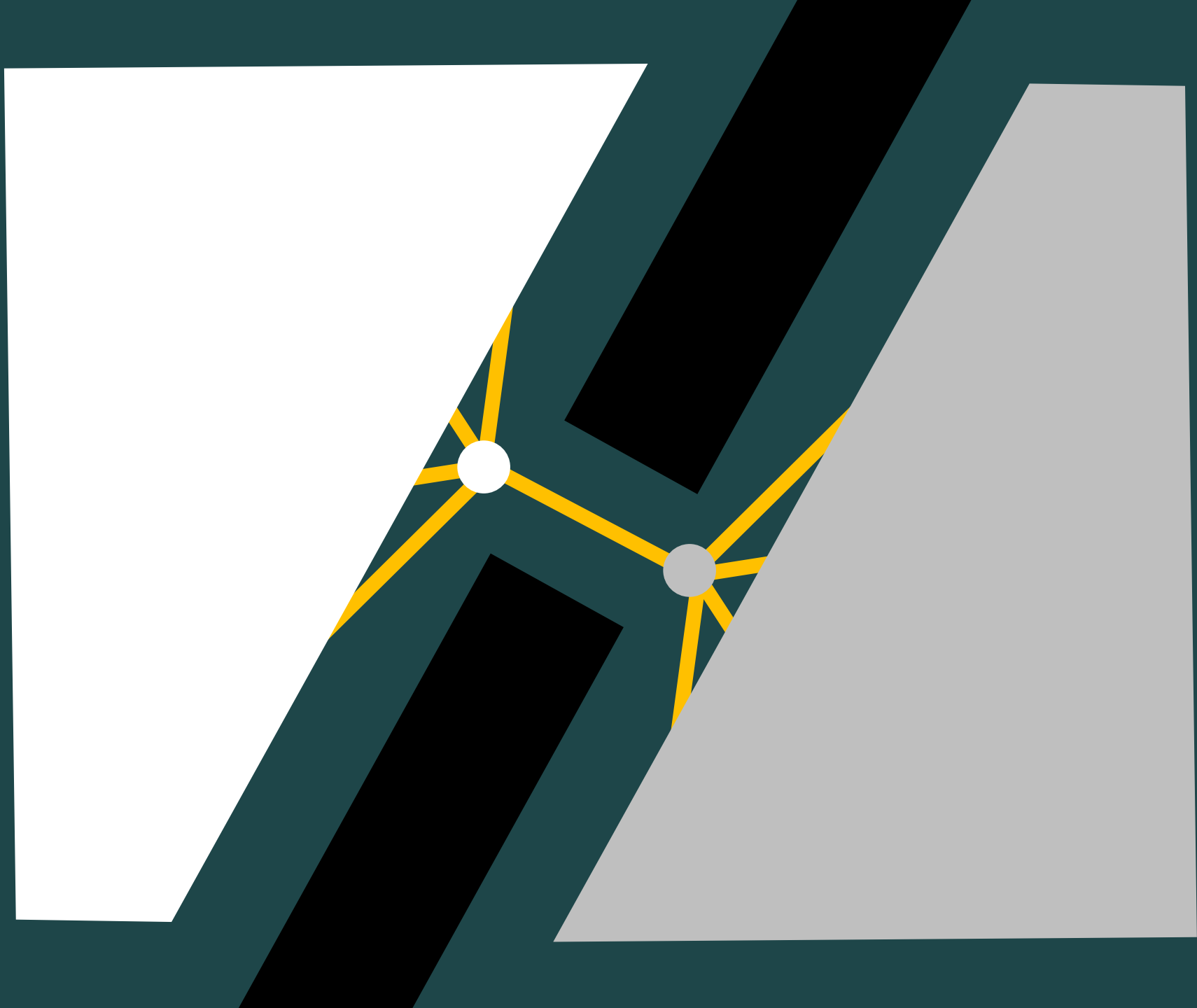


# Local tests of global entanglement and a counterexample to the generalized area law.



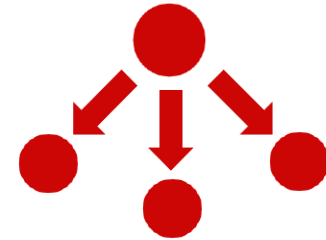






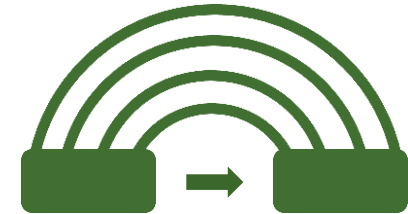
# 1 q. expanders

maximally entangled states



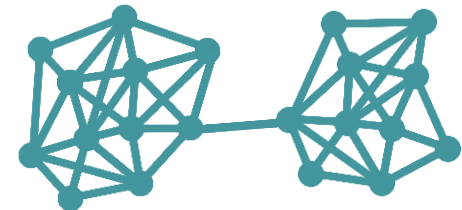
# 2 entanglement

testing and communication

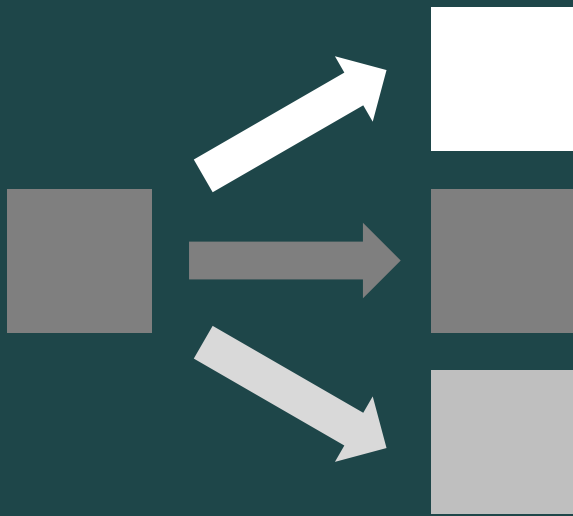


# 3 area law

gaps, connections, correlations





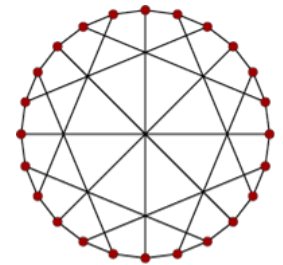
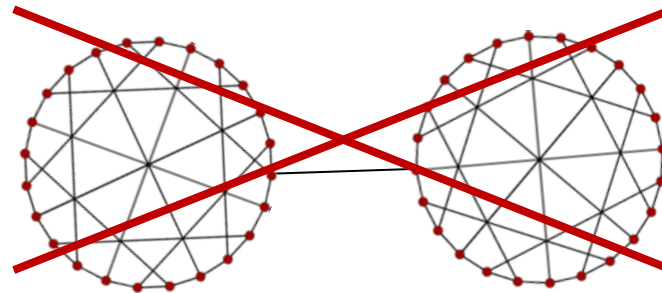
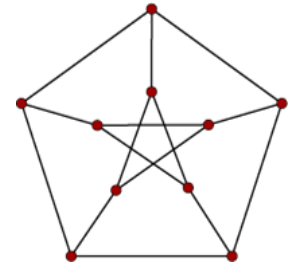


Quantum  
Expanders



# 1 Classical expanders

- Aram Harrow's talk, QHC workshop at the [youtube Harrow quantum expanders]
- graphs that mix well  
divide in two? cut a lot (fraction) of edges!

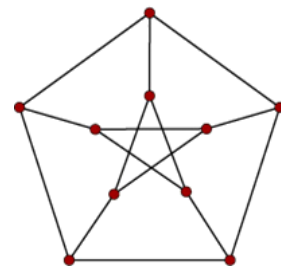


# 1 Classical expanders

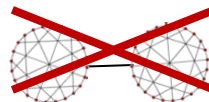
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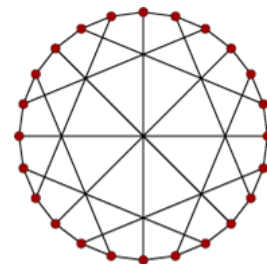
- graphs that mix well  
divide in two? cut a lot (fraction) of edges!



examples: Cayley graphs



- normalized adjacency matrix  
second largest eigenvalue  $1-\lambda$



- a motivation for quantum expanders  
 $d$ -regular graphs, random permutations

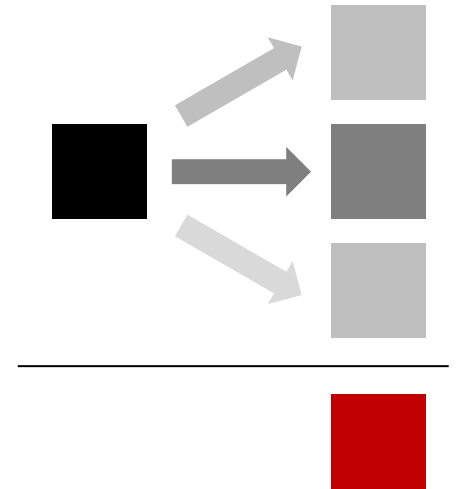
$$A = \frac{1}{k} \sum_{i=1}^k \Pi_i$$



# 1 Mixing up something quantum

- applying random unitaries

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$



- a motivation for quantum expanders  
*d*-regular graphs, random permutations

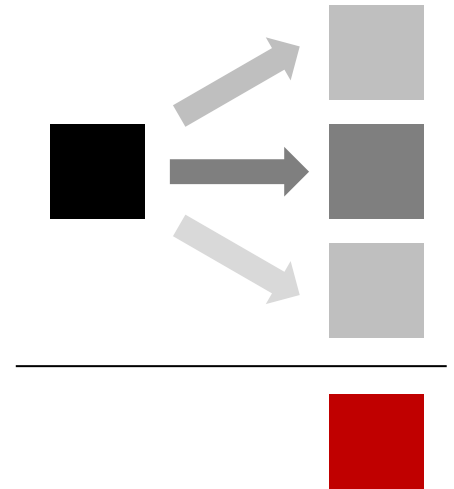
$$A = \frac{1}{k} \sum_{i=1}^k \Pi_k$$

# 1 Quantum expanders

- applying random unitaries from a set  
a discrete approximation to the Haar measure

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

- transform  $N \times N$  matrices



- a small second largest singular value  $\lambda$   
not far from the depolarizing channel

$$\|\mathcal{E} - D\|$$

- q. expander constructions, also for fixed  $k$  (8, 5, 3?)

[Ben-Aroya+ '07, Hastings '07, Gross & Eisert '08, Hastings & Harrow '09]



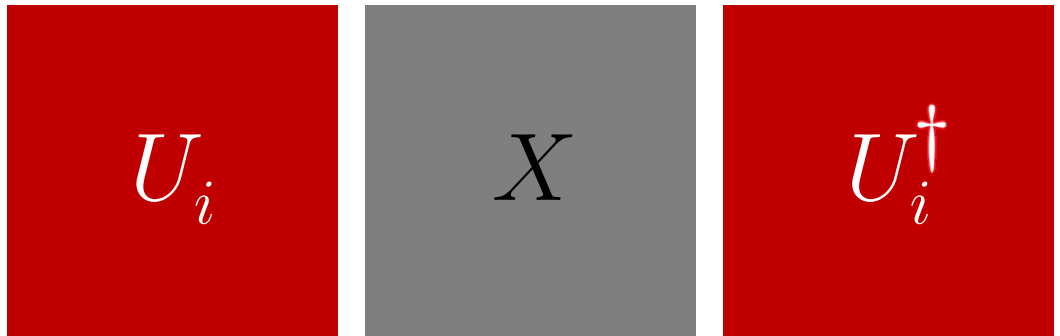
# 1 Quantum expanders

- transform  $N \times N$  matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

- a matrix that doesn't change?

$$X = \mathbb{I}$$



$$U_i X = X U_i$$



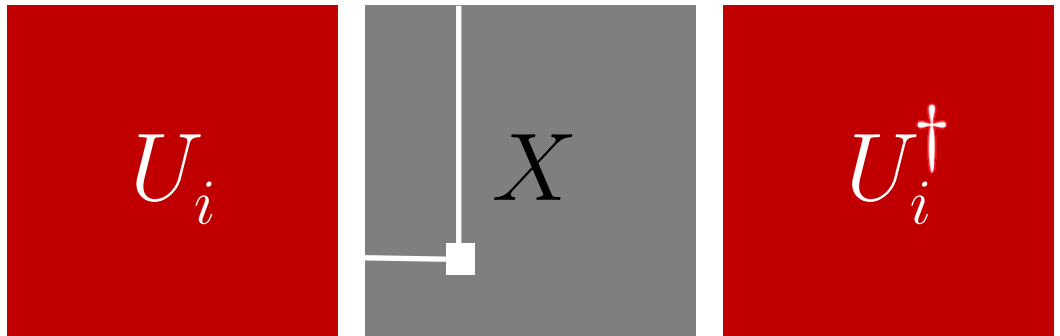
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- interpreting matrices as 2-register states

$$\sum_{a,b} X_{ab} |a\rangle \langle b|$$

density matrix



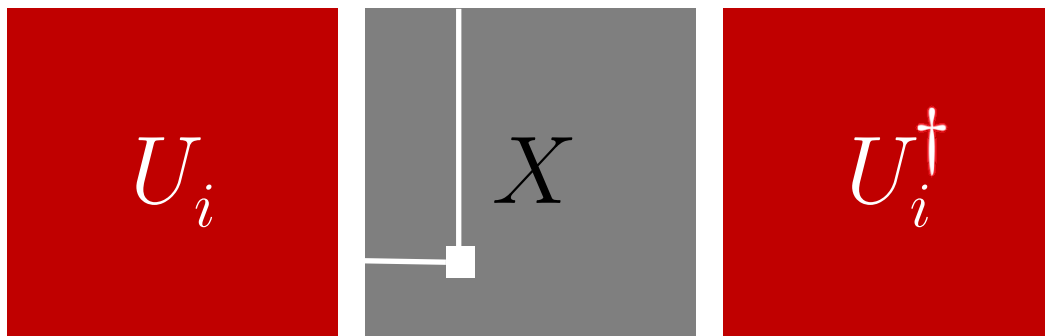
# 1 Quantum expanders

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- interpreting matrices as 2-register states

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle |b\rangle$$

- distributively applying an expander stationary?

pure state  
↓

$$\sum_{a,b} X_{ab} |a\rangle |b\rangle$$

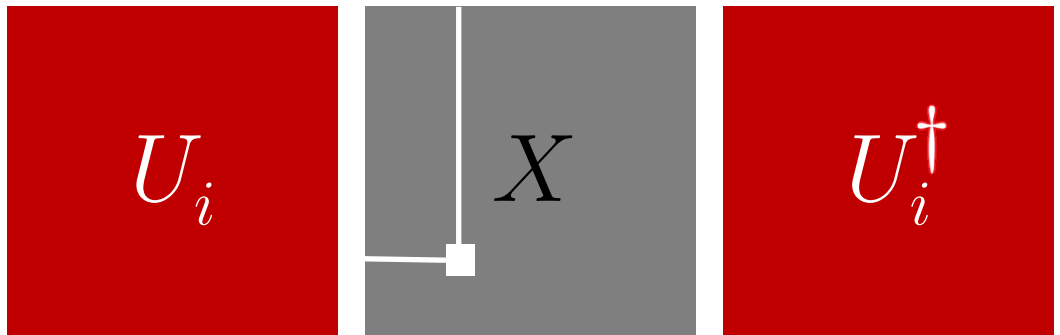
# 1 Quantum expanders

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- interpreting matrices as 2-register states

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle |b\rangle$$

- distributively applying an expander

stationary? max. entangled!

$$|\Phi_N\rangle = \frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$$



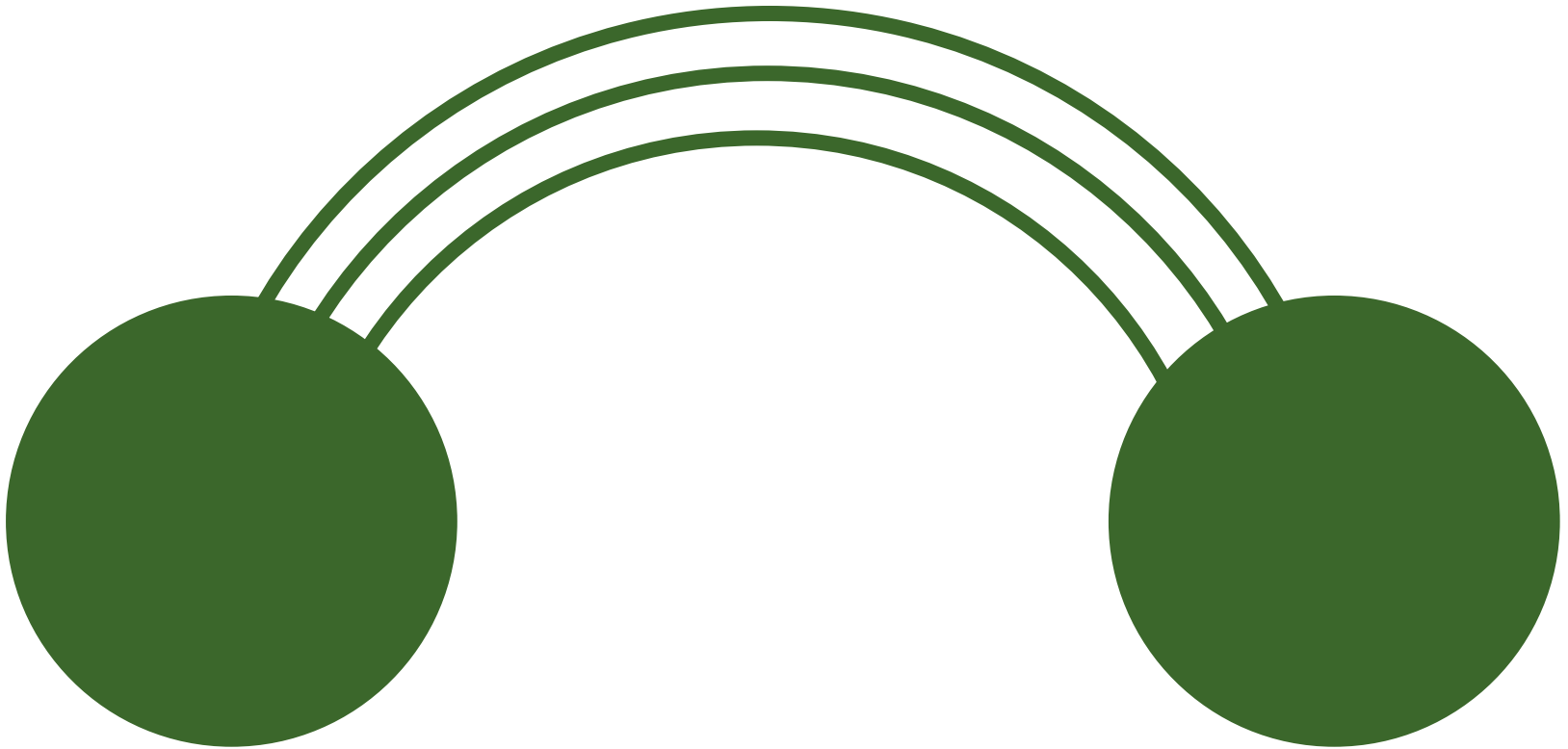
local tests of global entanglement



## 2 EPR testing

- how costly is it to certify that we share a maximally entangled state?

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$$

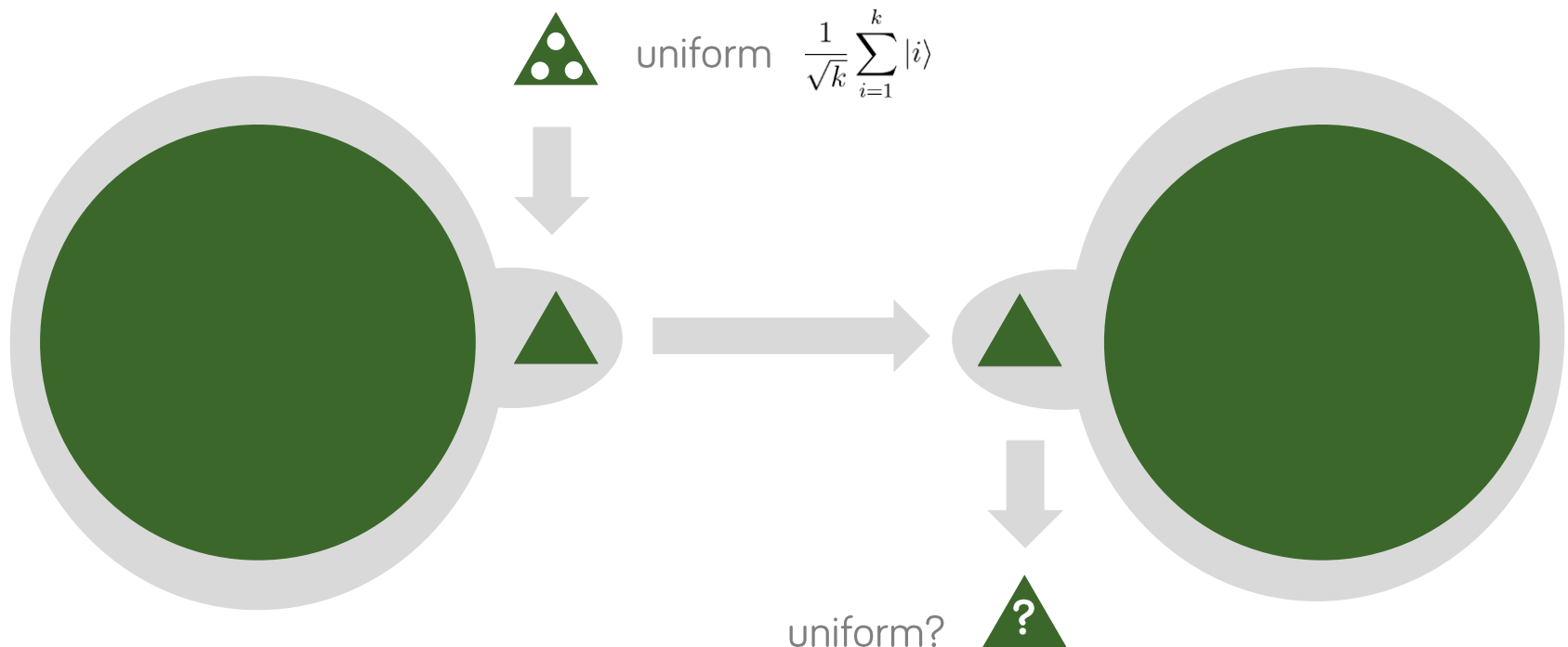


## 2 EPR testing

- how costly is it to certify that we share a maximally entangled state?
- apply a quantum expander distributively

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$$

$$U_i \otimes U_i^*$$



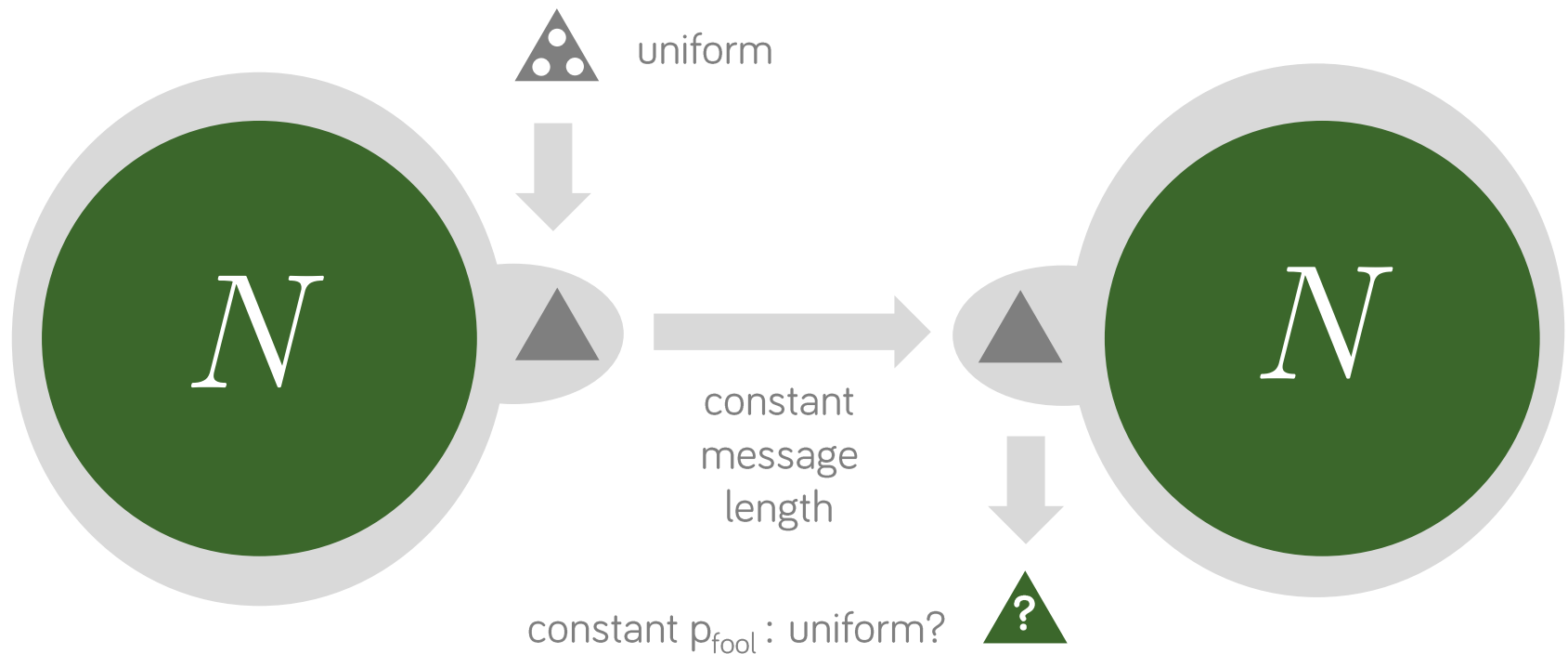
## 2 EPR testing

- action on states
- does the qutrit remain uniform?  
does the matrix  $X$  commute with  $U_i$ ?
- quantum expander ... soundness

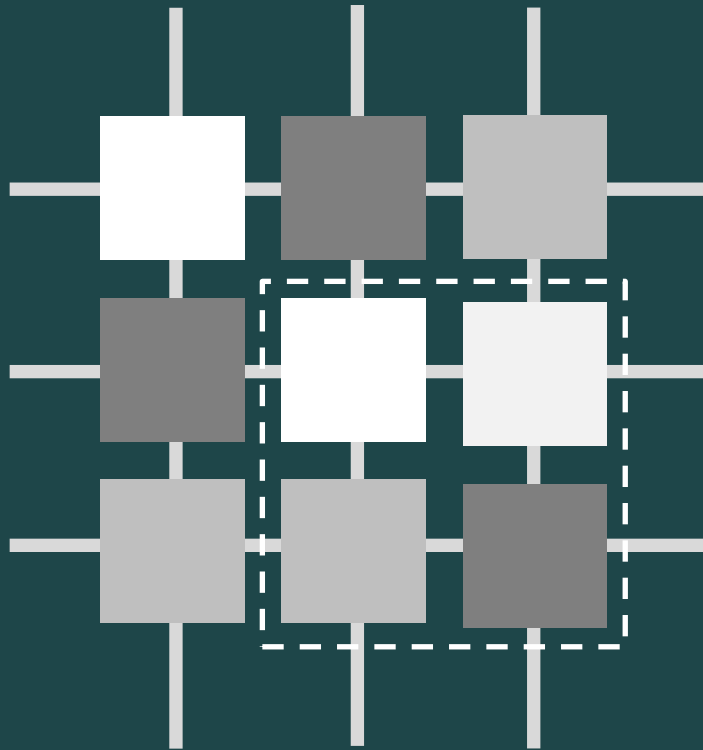
$$\frac{1}{\sqrt{k}} \sum_{i=1}^k |i\rangle (U_i \otimes U_i^*) |X\rangle$$

$$\sum_{a,b} X_{ab} |a\rangle |b\rangle$$

$$U_i \quad X \quad U_i^\dagger$$







a counterexample to the  
generalized  
area law

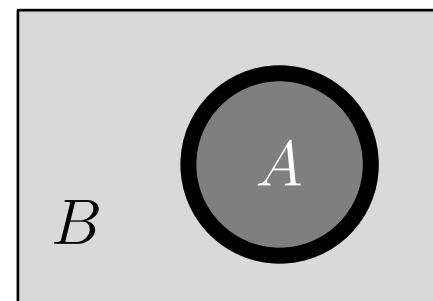
### 3 Area law: ground states of quantum spin systems

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \text{volume}$$

area

Schmidt coefficients



$$\rho_A = \text{Tr}_B \rho$$

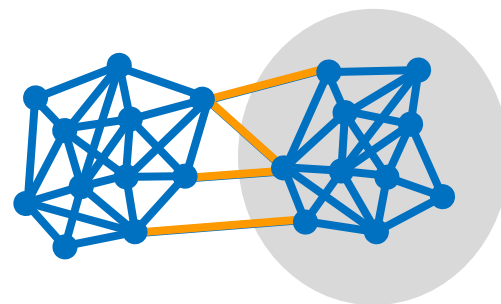
- 1D ... algorithms [White 92, Vidal 03, Landau+ 13]

theorems [Hastings 07, Arad+ 13]

- 2D ... we're close

small gap? large local dimension?

- generalized area conjecture  
entropy  $\sim$  cut size



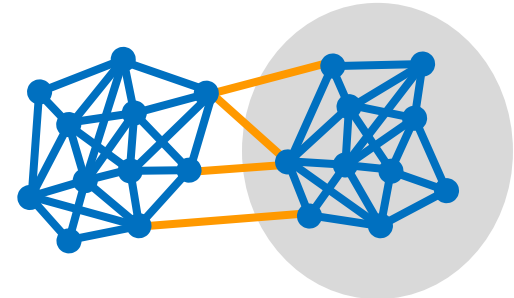
# Not true.

a gap  
a few links  
 $O(1)$  terms



not much  
entanglement  
(a “simple” ground state)

- generalized area conjecture  
entropy  $\sim$  cut size

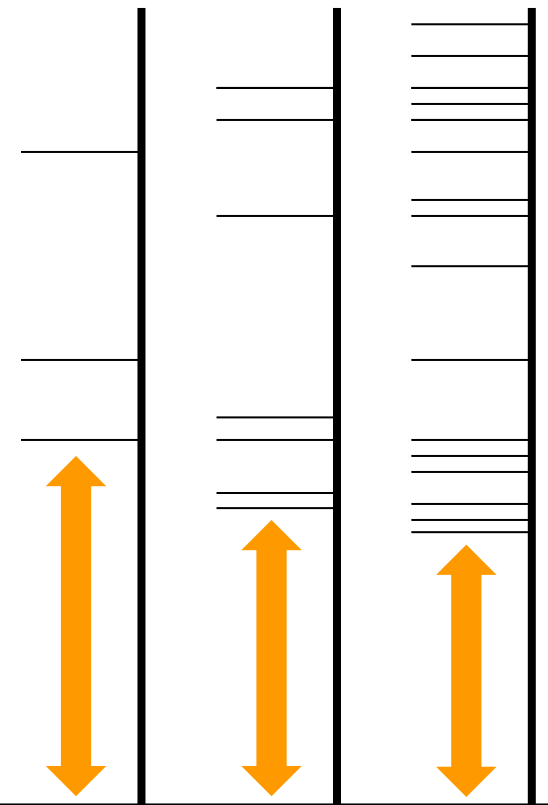


### 3 Gapped Hamiltonians

- Are there any states close to the ground state when we take the thermodynamic limit?
- local,  $O(1)$  norm terms

~~an inverse-poly gap?  $\Delta = \frac{c}{N} \rightarrow 0$~~

$$N \rightarrow \infty$$

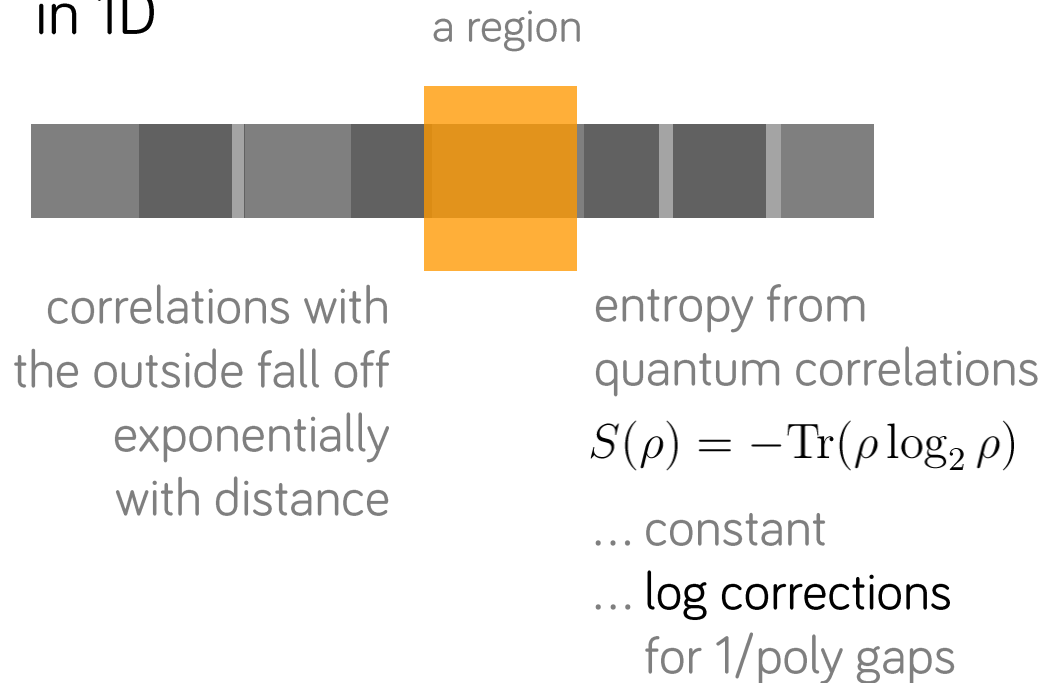




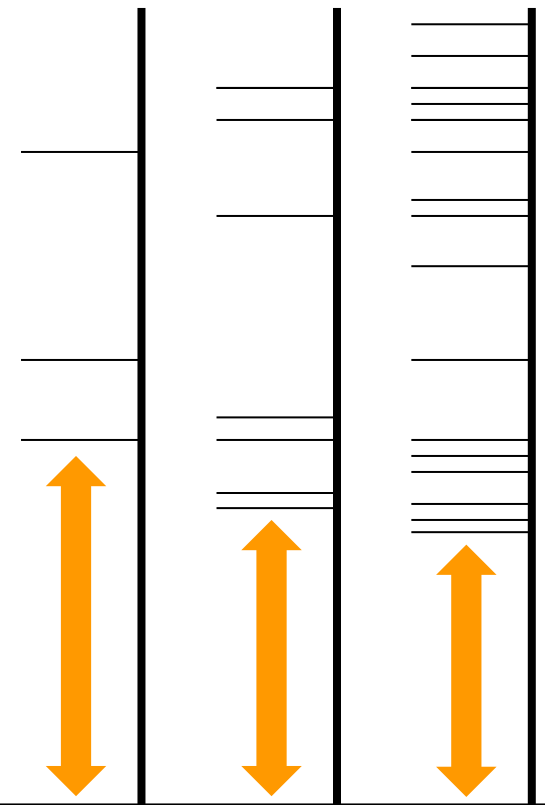
### 3 Gapped Hamiltonians

- Nothing closer than  $\Delta$  to the ground state.

- in 1D



$$N \rightarrow \infty$$



### 3 Gapped Hamiltonians

- Nothing closer than  $\Delta$  to the ground state.

- in 1D

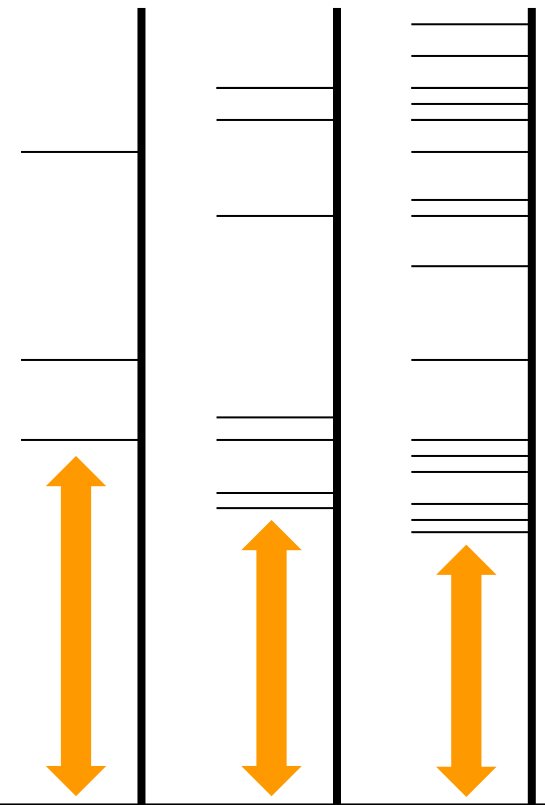
the AKLT (spin-1) chain

$$\sum_{j=1}^{N-1} \vec{S}_j \cdot \vec{S}_{j+1} + \frac{1}{3} \left( \vec{S}_j \cdot \vec{S}_{j+1} \right)^2$$

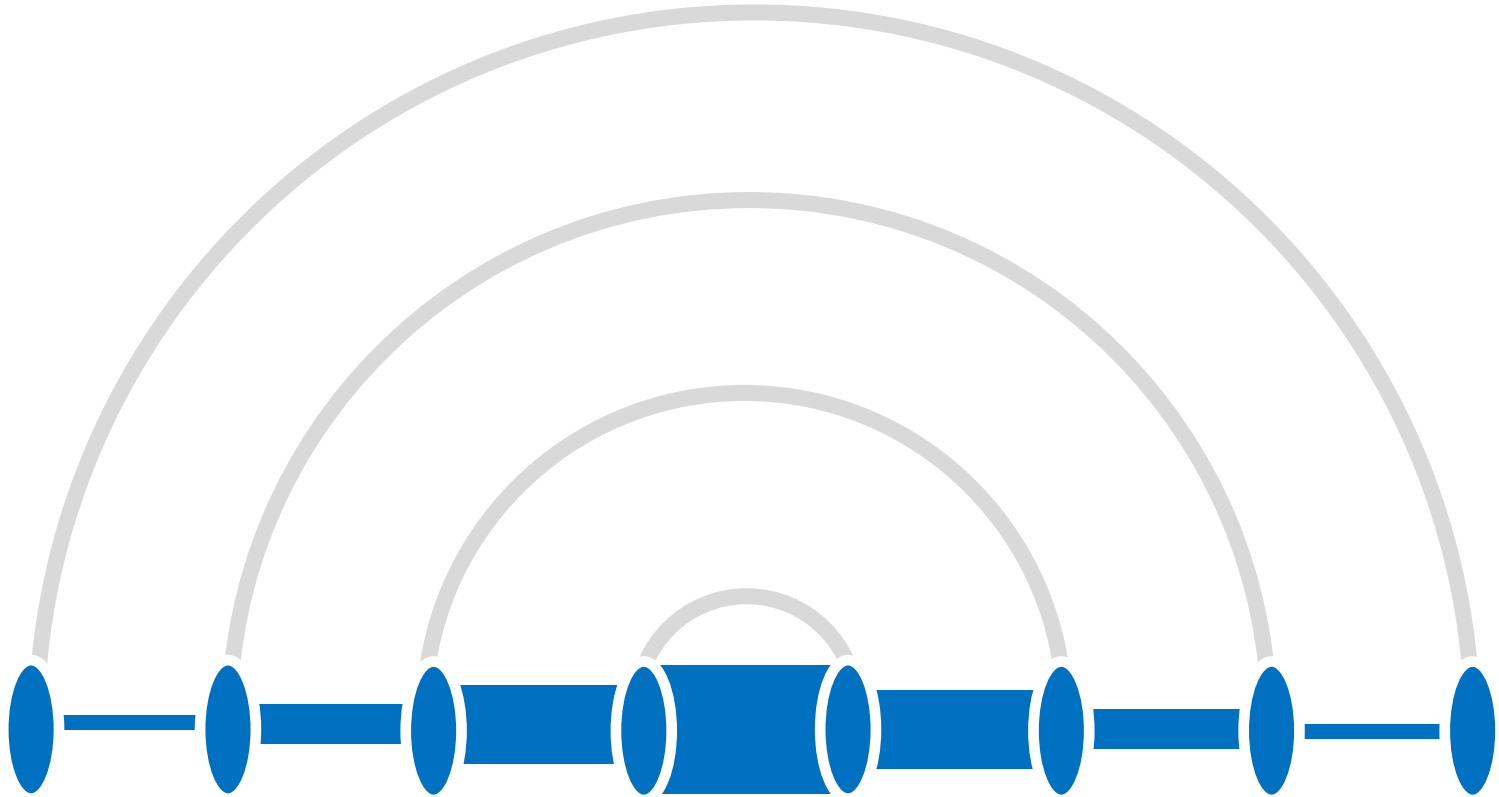
a biased walk in 1D

$$\sum_{j=1}^{N-1} (|j\rangle - B|j+1\rangle) (\langle j| - B\langle j+1|)$$

$$N \rightarrow \infty$$



- a large gap ... a simple (not too entangled) ground state?

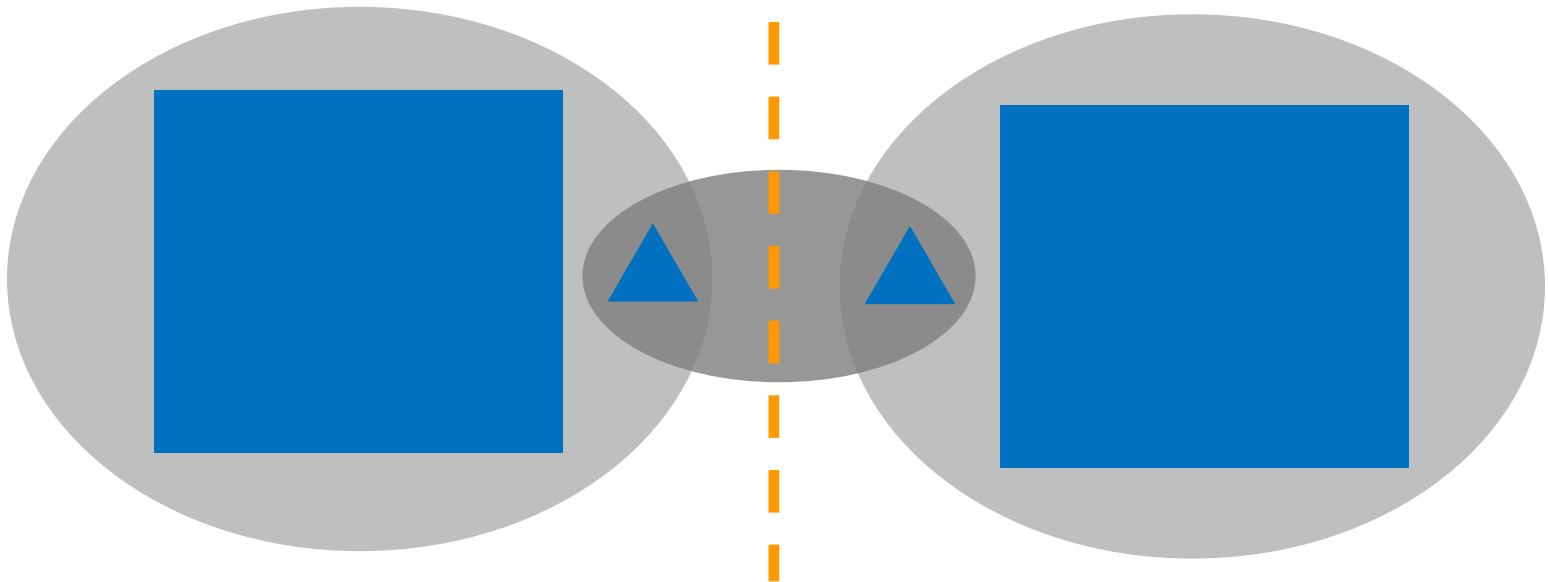


- without a gap, the entropy can be large

[Verstraete, Latorre+]

### 3 Our counterexample to the generalized area conjecture

- an  $N \times 3 \times 3 \times N$  dimensional system
- a frustration-free, gapped, Hamiltonian
- a single  $O(1)$  interaction of two  $3 \times 3$  subsystems
- a unique, very entangled ground state with  $O(N)$  entanglement entropy across the cut



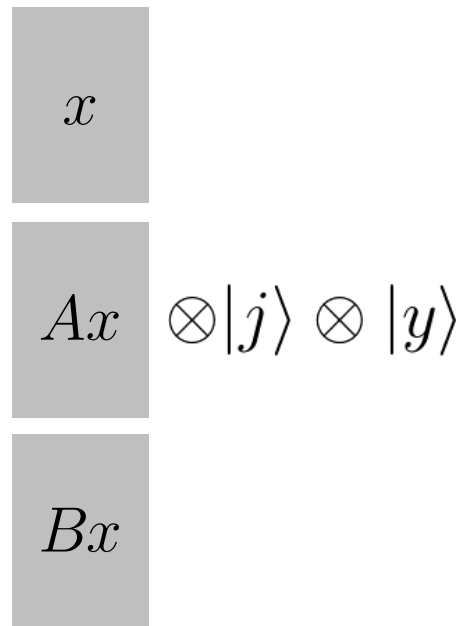


### 3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

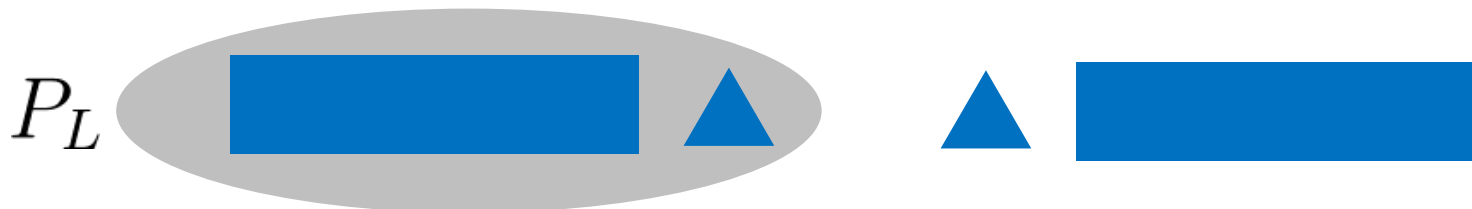
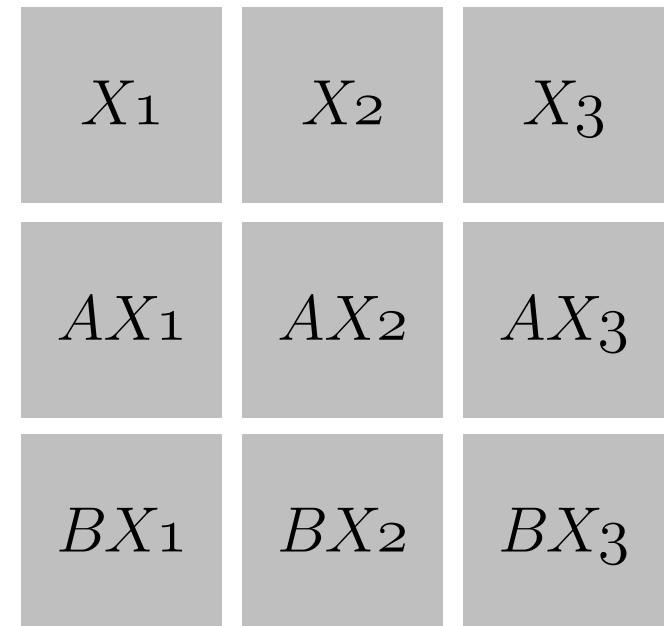
- a projector  $P_L$  with ground states

$$\frac{1}{\sqrt{3}} (|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$$

- as a vector



- as a matrix



### 3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

- a projector  $P_R$   
with ground states

$$\frac{1}{\sqrt{3}} (|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle)$$

- as a vector  $|i\rangle \otimes |x\rangle \otimes$ 

$y$	$yA$	$yB$
-----	------	------

as a matrix

$Y_1$	$Y_1A$	$Y_1B$
$Y_2$	$Y_2A$	$Y_2B$
$Y_3$	$Y_3A$	$Y_3B$



### 3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

- a projector  $P_L$
- a projector  $P_R$

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

$X_1$	$X_2$	$X_3$
$AX_1$	$AX_2$	$AX_3$
$BX_1$	$BX_2$	$BX_3$

$X$	$XA$	$XB$
$AX$	$AXA$	$AXB$
$BX$	$BXA$	$BXB$

$Y_1$	$Y_1A$	$Y_1B$
$Y_2$	$Y_2A$	$Y_2B$
$Y_3$	$Y_3A$	$Y_3B$



### 3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

■ a projector  $P_L$

a projector  $P_R$

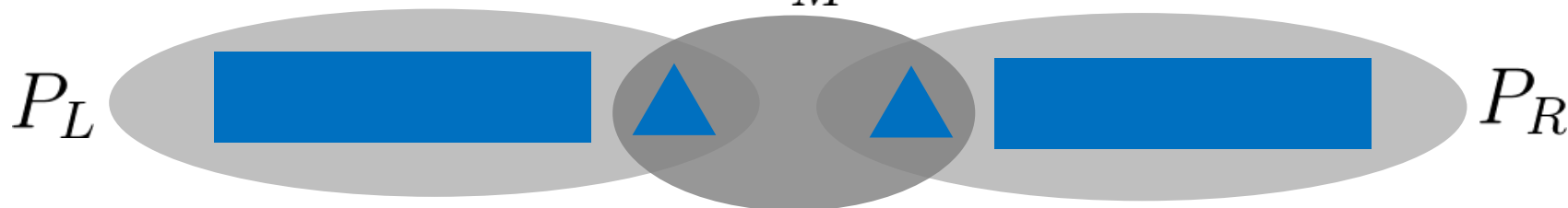
a projector  $P_M$

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

$X$	$XA$	$XB$
$AX$	$AXA$	$AXB$
$BX$	$BXA$	$BXB$

$P_M$



### 3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

- a projector  $P_L$
- a projector  $P_R$
- a projector  $P_M$

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

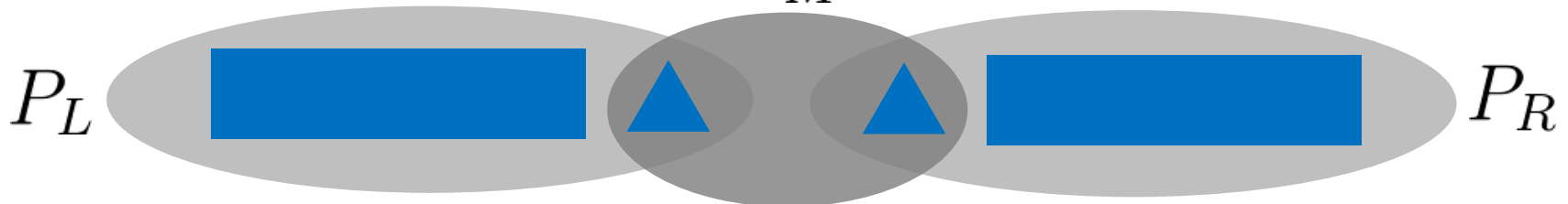
- who commutes with  $A$  and  $B$ ?
- only the identity,  
as  $[I, A, B]$  are  
a q. expander

enforce symmetry

for 12 & 21  
for 13 & 31

$X$	$XA$	$XB$
$AX$	$AXA$	$AXB$
$BX$	$BXA$	$BXB$

$P_M$





### 3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

- a projector  $P_L$
- a projector  $P_R$
- a projector  $P_M$

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

enforce symmetry

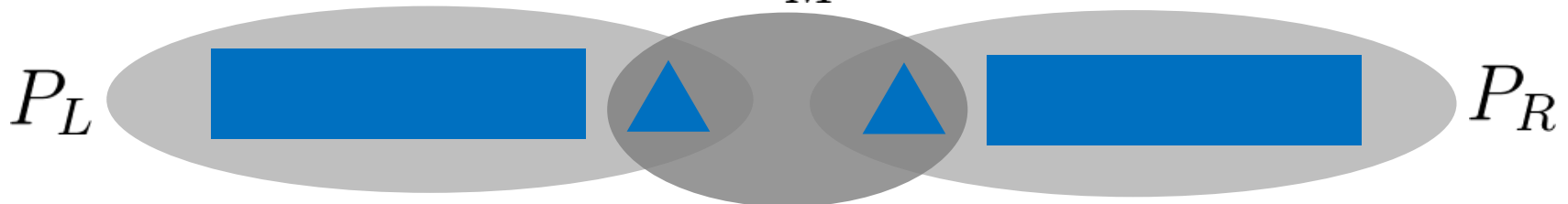
for 12 & 21  
for 13 & 31

- who commutes with  $A$  and  $B$ ?
- only the identity,  
as  $[I, A, B]$  are  
a q. expander

$I$	$A$	$B$
$A$	$AA$	$AB$
$B$	$BA$	$BB$

$$\frac{1}{3\sqrt{N}}$$

$P_M$



### 3 A nonlocal Hamiltonian, $N \times 3 \times 3 \times N$

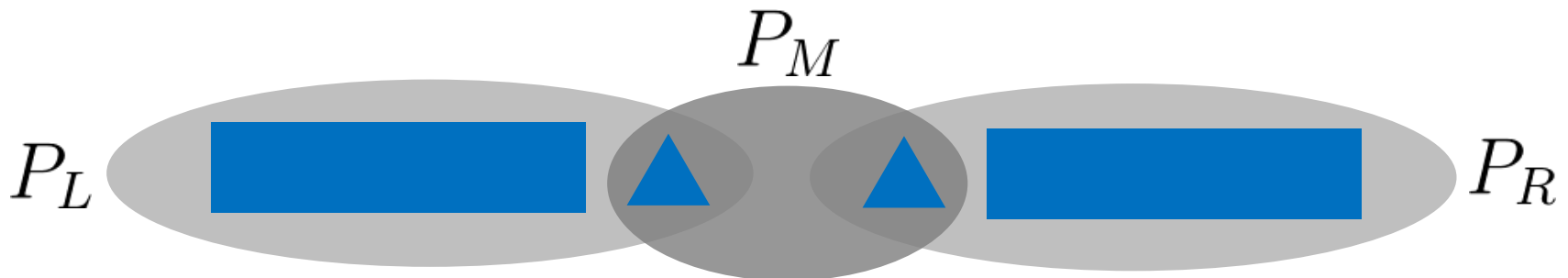
- ground state:                      unique  
    very entangled

- Hamiltonian:  
frustration free  
gapped



$I$	$A$	$B$
$A$	$AA$	$AB$
$B$	$BA$	$BB$

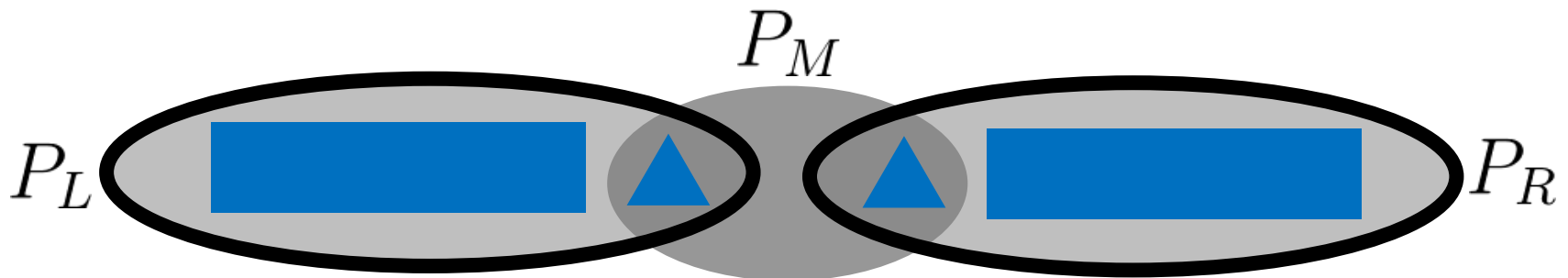
$$\sum_{x=1}^N |x\rangle |x\rangle$$

$$\frac{1}{3\sqrt{N}}$$



### 3 Making the counterexample local

- quantum expander  $[I, A, B]$  ... quantum circuits ...  
~~nonlocal projectors~~ ... Kitaev's LH & history states   
an approximate groundstate, a very small gap
- rescale  $P_L, P_R$  (not the middle!)  
huge, nonphysical couplings
- use new “strengthening gadgets” [N., Cao]   
~~large interaction strength~~ ... extra particles, high degree



### 3 A local Hamiltonian, $(N+n) \times 3 \times 3 \times (N+n)$

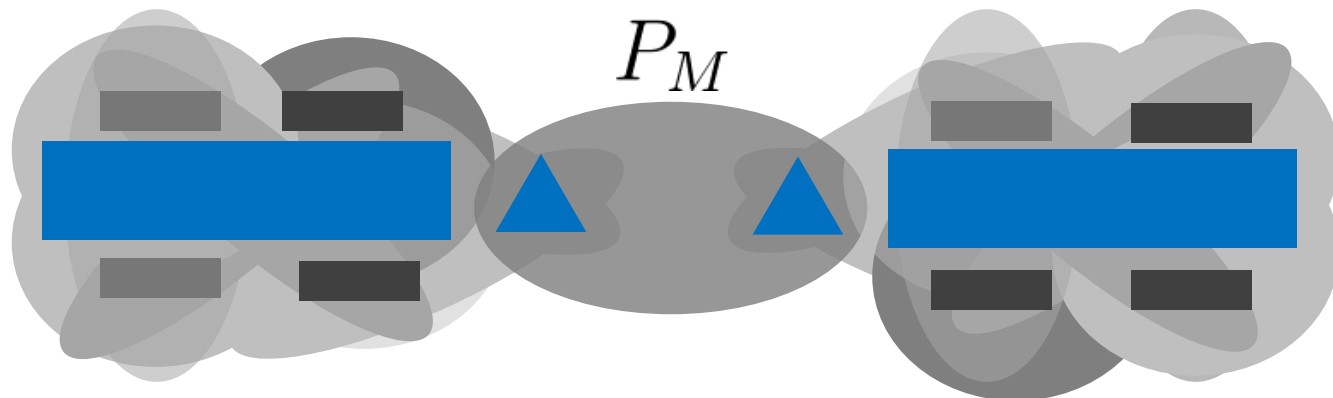
- frustrated, but still gapped

$O(1)$  norm terms

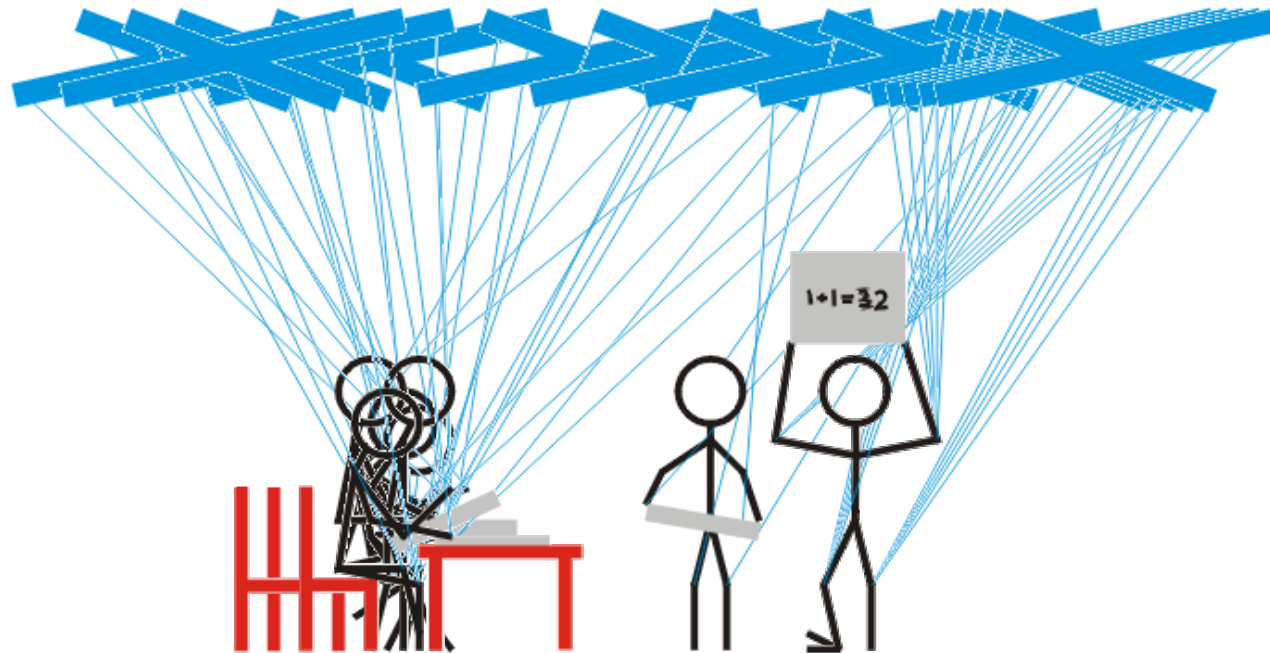
- a unique and still very entangled ground state

$$\approx |w\rangle \otimes$$

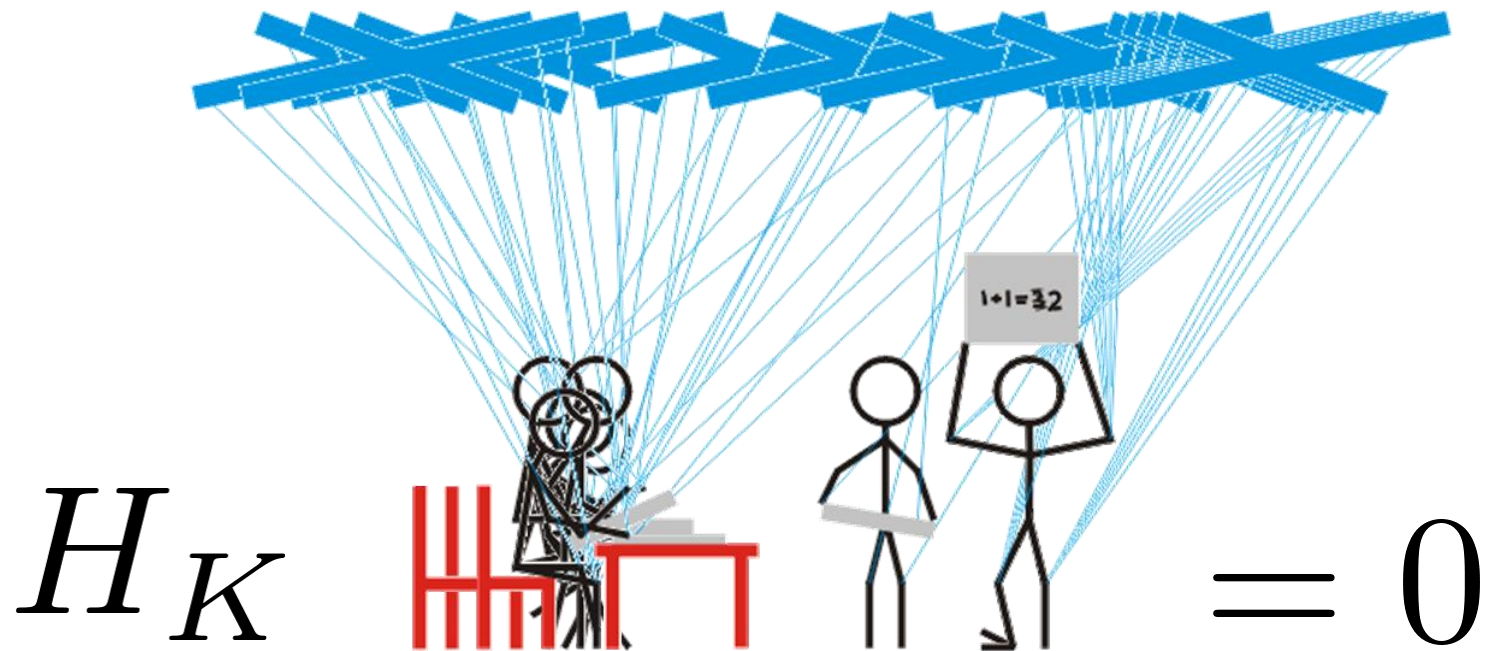
$I$	$A$	$B$
$A$	$AA$	$AB$
$B$	$BA$	$BB$



### 3 Implementing circuits locally: Feynman's computer



### 3 The history state: a ground state



$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

$\underbrace{\hspace{10em}}_{U_t \cdots U_1 |\varphi_0\rangle}$



### 3 The history state: a ground state

*k-local*

*c-o-n-d-i-t-i-o-n-s*

clock encoding  
state progression  
initialization

$$|\cdots\rangle|u\rangle \otimes |0\rangle$$

$$|\varphi_t\rangle \otimes |t\rangle$$

$$|\varphi_{t+1}\rangle \otimes |t+1\rangle$$

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

idling



most of the state has the result

### 3 A local Hamiltonian, $(N+n) \times 3 \times 3 \times (N+n)$

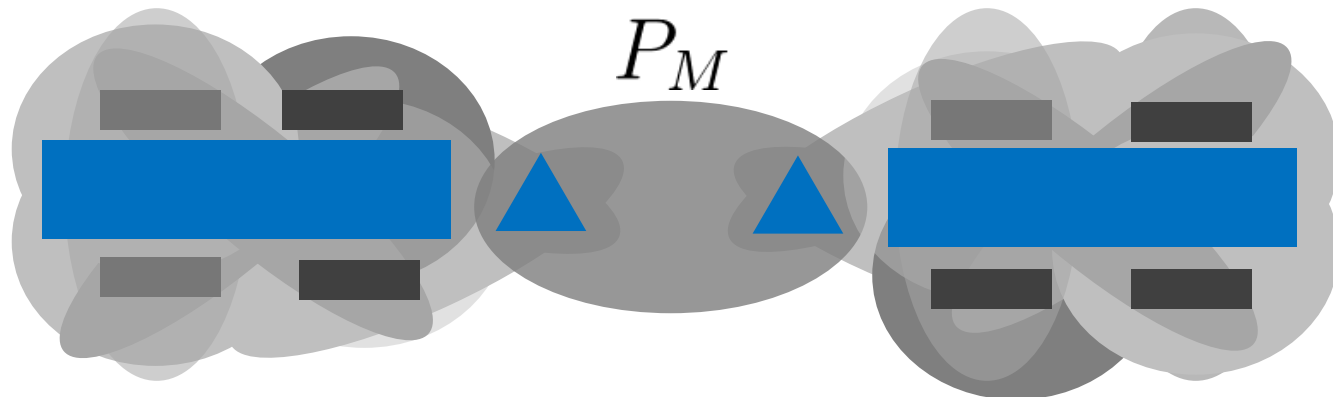
- frustrated, but still gapped

$O(1)$  norm terms

- a unique and still very entangled ground state

$$\approx |w\rangle \otimes$$

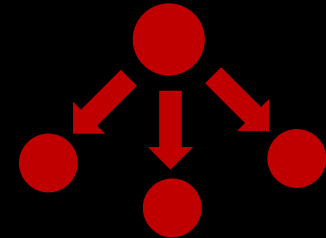
$I$	$A$	$B$
$A$	$AA$	$AB$
$B$	$BA$	$BB$



1

# q. expanders

maximally entangled states



2

# entanglement

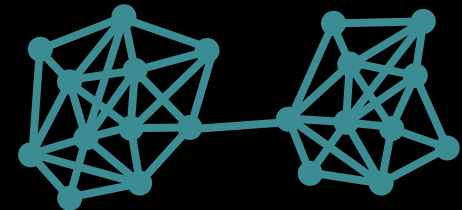
testing and communication



3

# area law

gaps, connections, correlations

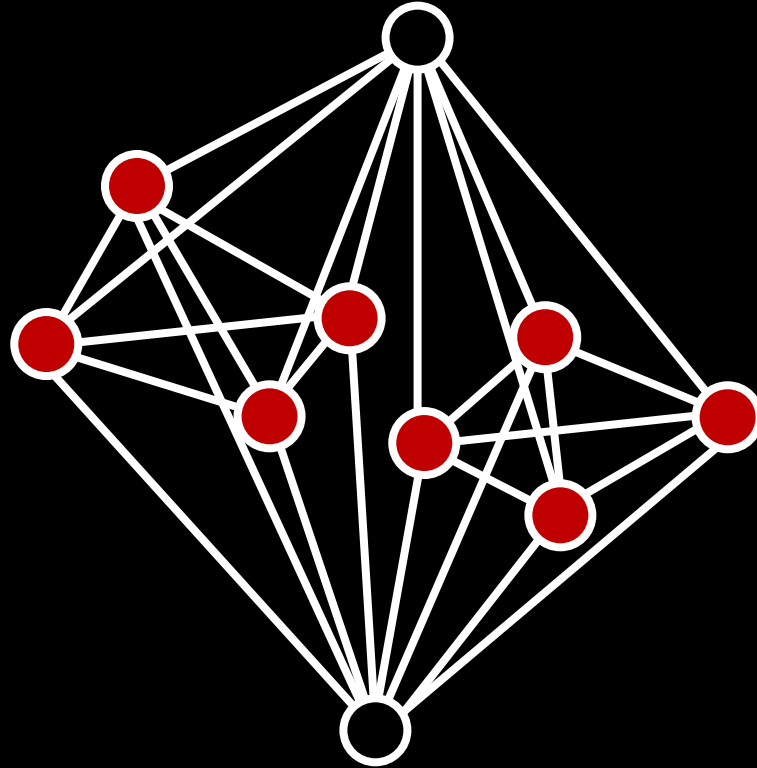


# Local tests of global entanglement and a counterexample to the generalized area law.









Daniel Nagaj



universität  
wien

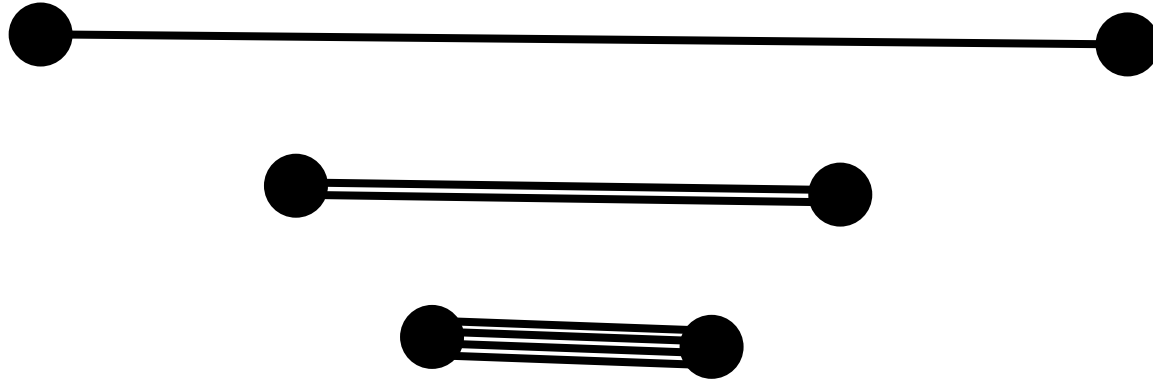


Yudong Cao

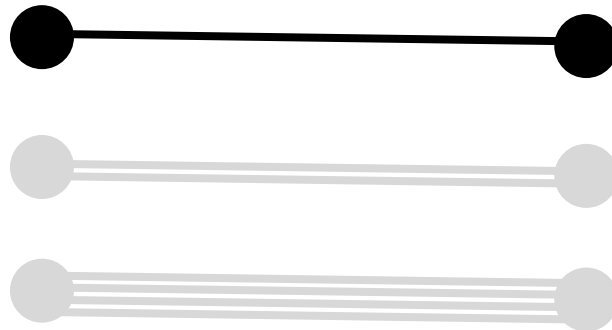
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## ■ interaction strength vs. distance



## ■ limited interaction strength





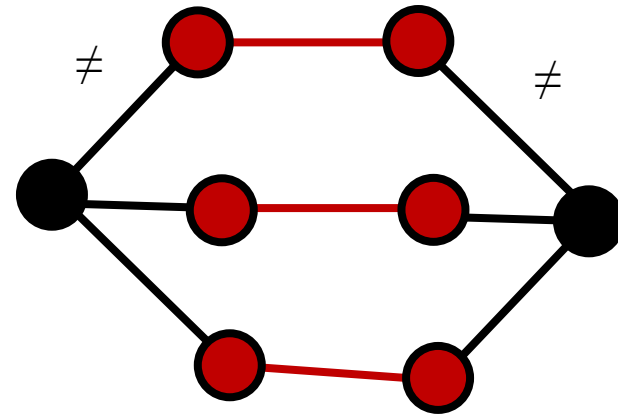
## ■ classical gadgets



00, 11, 22



00, 11, 22

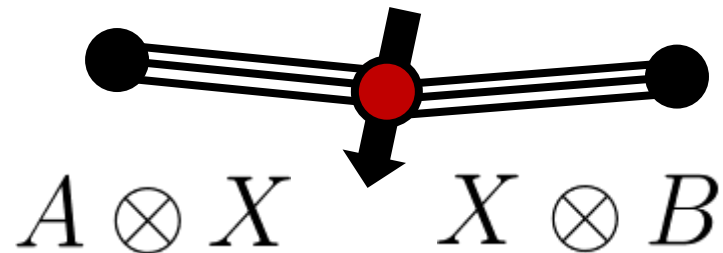


00, 11, 22

## ■ quantum gadgets



$A \otimes B$



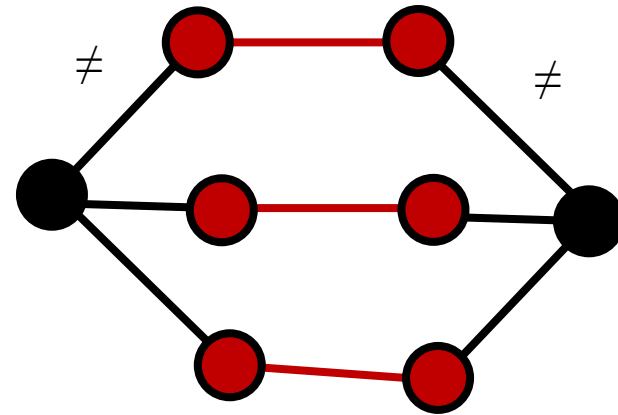
## ■ classical gadgets



00, 11, 22

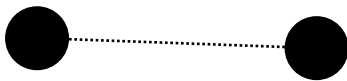


00, 11, 22

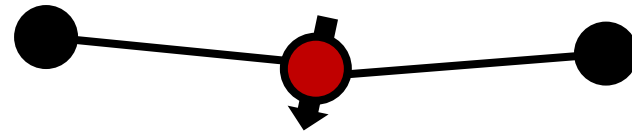


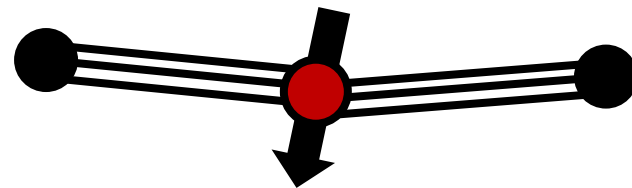
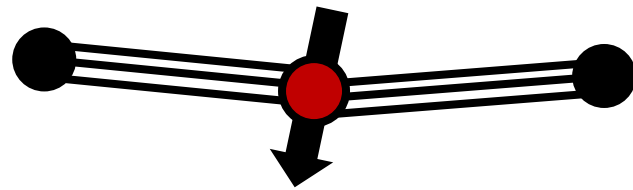
00, 11, 22

## ■ quantum gadgets

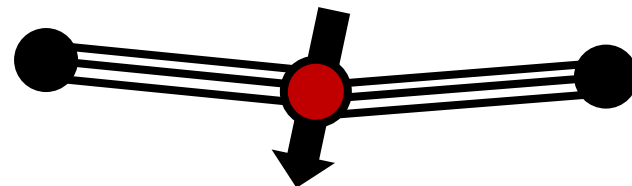
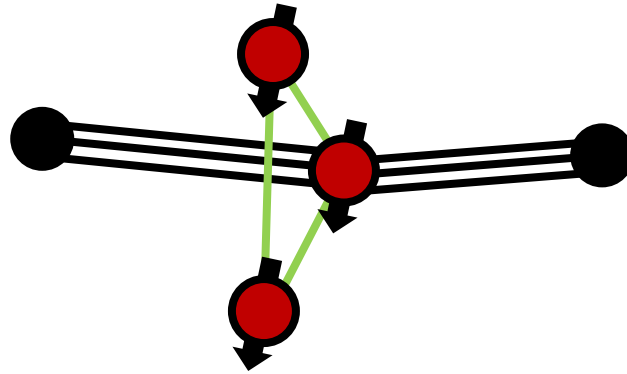


$A \otimes B$

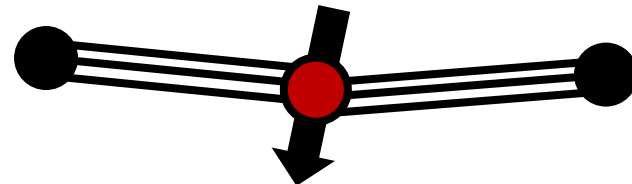
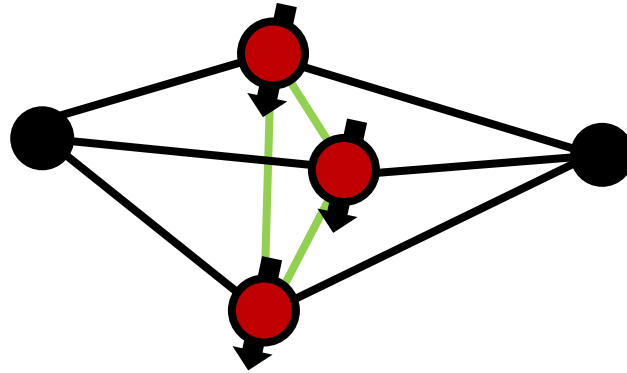




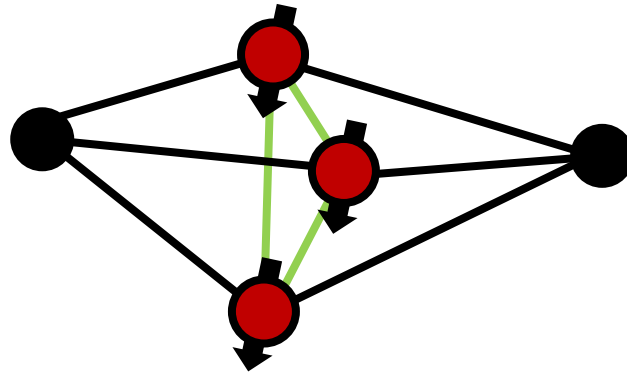
a “strong” field



“strong” interactions



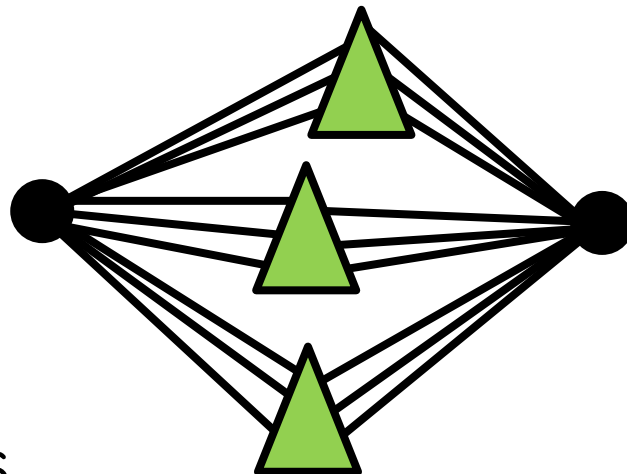
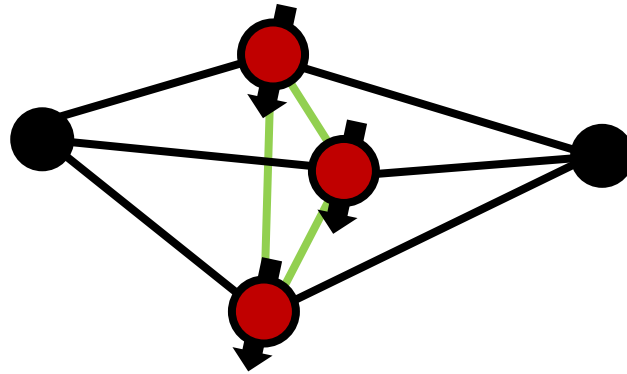
“strong” interactions



$$A \otimes B$$

one gadget

“strong” interactions



$$A \otimes B$$

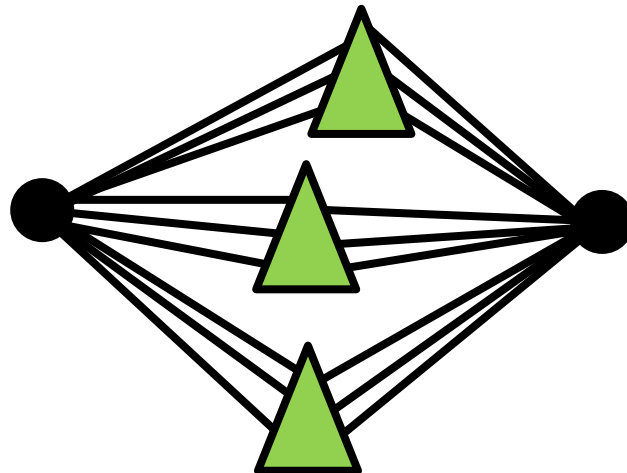
several gadgets

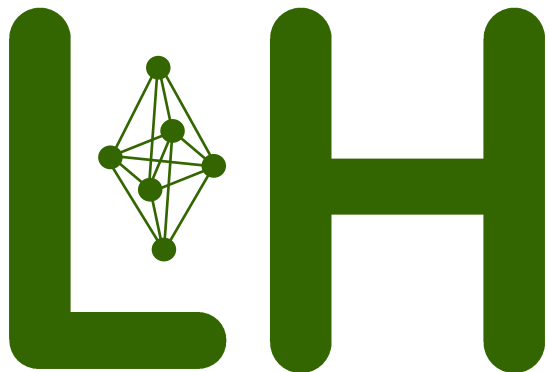


# **“strong” interactions**

weak components

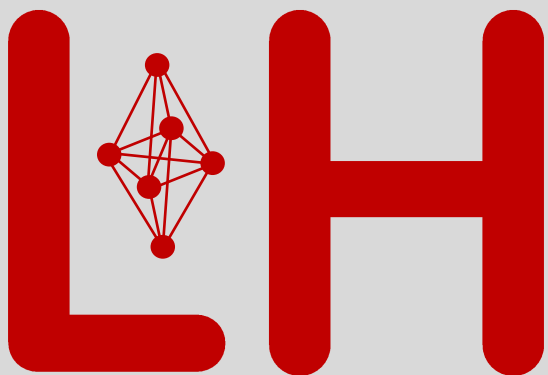
new parallel composition





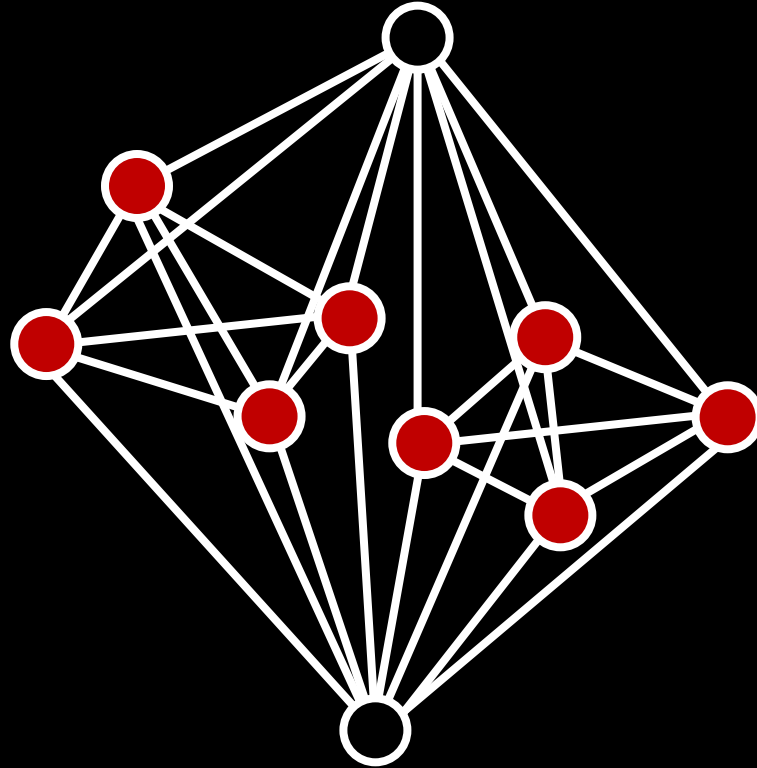
$O(1)$  terms? QMA-complete.

$1/\text{poly}$  gap? Constant gap.



High degree (poly).

Fractional gap? Worse.



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