many-body physics & complexity

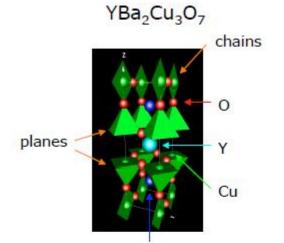
tough & interesting questions

Daniel Nagaj wien wien

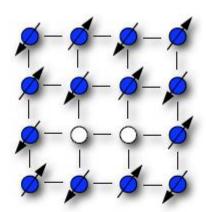




High Tc superconductors & Quantum Magnets

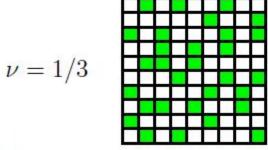


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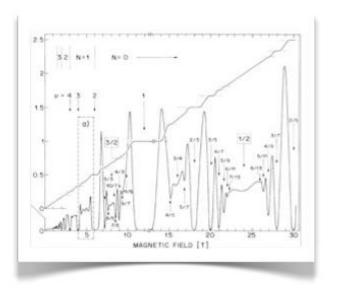




Fractional Quantum Hall Effect







[slide: A.Lauchli]

investigate & explain

construct something



Many-body systems & their Hamiltonians

MACROSCOPIC PROPERTIES

m-i-c-r-o-s-c-o-p-i-c r-u-l-e-s m-i-c-r-o-s-c-o-p-i-c r-u-l-e-

density, heat capacity, magnetization, correlations, spectrum ...

1

Many-body systems & their Hamiltonians

a model of interactions

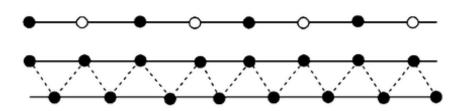
non-interacting = 1-body

$$E(s_1 \ldots s_n)$$

strings

$$\hat{H}|\psi_k\rangle = E_k|\psi_k\rangle$$

spins bosons fermions



a local Hamiltonian

$$\cdots + H_{a,b,c,d} + \cdots$$
 few-body

$$\cdots + H_{k,k+2} + \cdots$$
 short-range

long-range, many interactions: Aram

some lattice

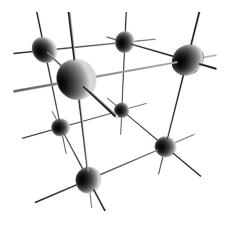
1D Ising, antiferromagnetic

$$E(s) = \sum_{j} s_{j} s_{j+1} - h \sum_{j} s_{j}$$
 O/1 strings

some lattice

3D Ising

$$E(s) = -J\sum_{\langle i,j\rangle} s_i s_j - h\sum_j s_j$$



[Melcom, de.wikipedia]

• lattice dimensionality: visit n^d points in n steps

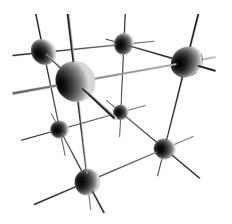
1

Geometry matters

some lattice

3D Ising, spin glass

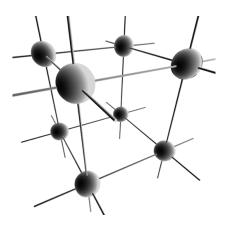
$$E(s) = -\sum_{\langle i,j\rangle} J_{ij} \, s_i s_j - h \sum_j s_j$$



[Melcom, de.wikipedia]

some lattice

$$E(s) = -J \sum_{\langle i,j \rangle} s_i s_j - h \sum_j s_j$$



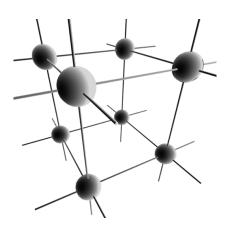
[Melcom, de.wikipedia]

1D Heisenberg (classical)

$$H = -\sum_{j} J_{j} \, \vec{s_{j}} \cdot \vec{s_{j+1}}$$
 3D vectors

some lattice

$$E(s) = -J\sum_{\langle i,j\rangle} s_i s_j - h\sum_j s_j$$



[Melcom, de.wikipedia]

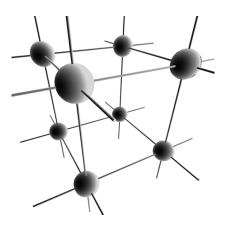
1D Heisenberg (quantum)

$$H = -\frac{1}{2} \sum_{j} \left(J_{x} \sigma_{j}^{x} \sigma_{j+1}^{x} + J_{y} \sigma_{j}^{y} \sigma_{j+1}^{y} + J_{x} \sigma_{j}^{z} \sigma_{j+1}^{z} \right) + h \sum_{j} \sigma_{j}^{z}$$

the Haldane conjecture about integer/half-integer spin

some lattice

$$E(s) = -J \sum_{\langle i,j \rangle} s_i s_j - h \sum_j s_j$$



[Melcom, de.wikipedia]

1D Heisenberg (quantum)

$$H = -\frac{1}{2} \sum_{j} \left(J_{x} \sigma_{j}^{x} \sigma_{j+1}^{x} + J_{y} \sigma_{j}^{y} \sigma_{j+1}^{y} + J_{x} \sigma_{j}^{z} \sigma_{j+1}^{z} \right) + h \sum_{j} \sigma_{j}^{z}$$

general graphs

general graphs Bose-Hubbard
$$H=t\sum_{\langle i,j\rangle}^{\text{hopping}}b_i^\dagger b_j+rac{U}{2}\sum_i^{\text{repulsion}}\hat{n}_i(\hat{n}_i-1)$$

a very complex graph [Childs Gosset Webb 13]

(fixed particle number)

a lifetime with a model



hubbard model

Web

Images

Maps

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Search tools

About 14,800,000 results (0.35 seconds)

Hubbard model - Wikipedia, the free encyclopedia

en.wikipedia.org/wiki/Hubbard_model *

The **Hubbard model** is an approximate model used, especially in solid state physics, to describe the transition between conducting and insulating systems.

Theory (Narrow energy band theory) - Example: 1D chain of hydrogen ...

Bose-Hubbard model - Wikipedia, the free encyclopedia

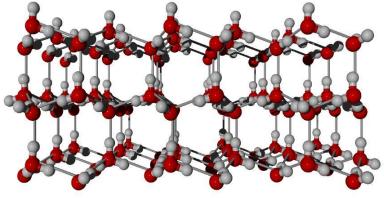
en.wikipedia.org/wiki/Bose-Hubbard model -

The Bose-Hubbard model gives an approximate description of the physics of interacting bosons on a lattice. It is closely related to the Hubbard model which ...

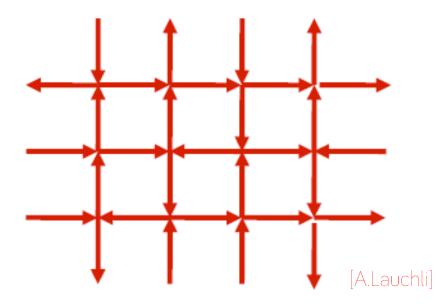
[PDF] The Hubbard Model for Dummies

www.physik.unizh.ch/lectures/electroncorrelations/**Hubbard**4Du.pdf \checkmark Jun 13, 2007 - Electron states for the (N=8) **Hubbard model** in terms of fermionic operators: c. 1\(\gamma\), \(\frac{1}{2}\), \(\frac{1}{

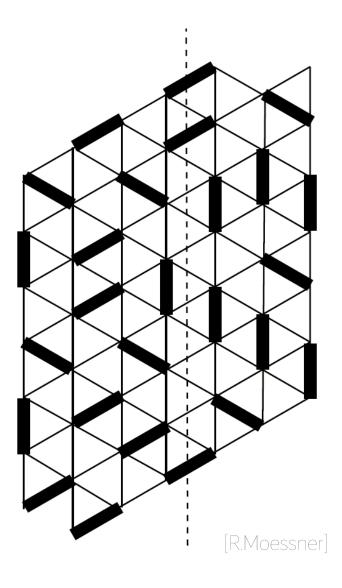
ice-type



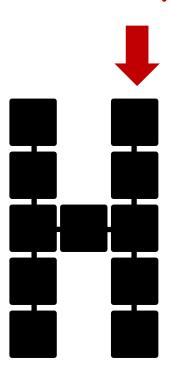
[Linkopings U.]

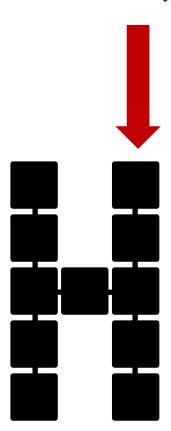


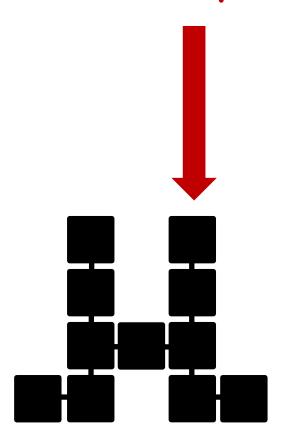
quantum dimers

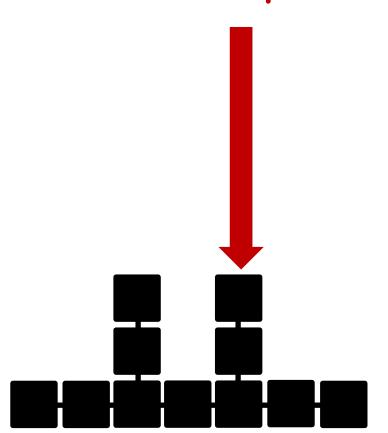


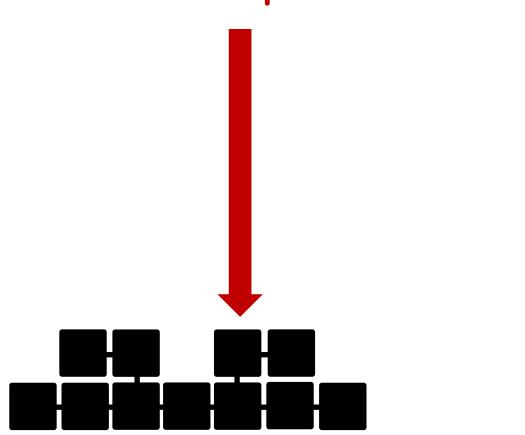
STATIC the questions





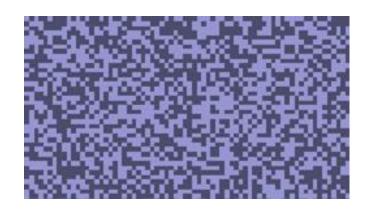




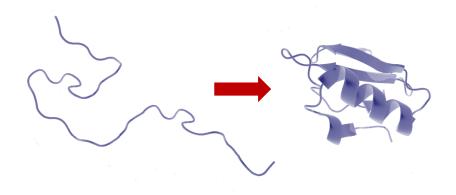


spin glasses

protein folding

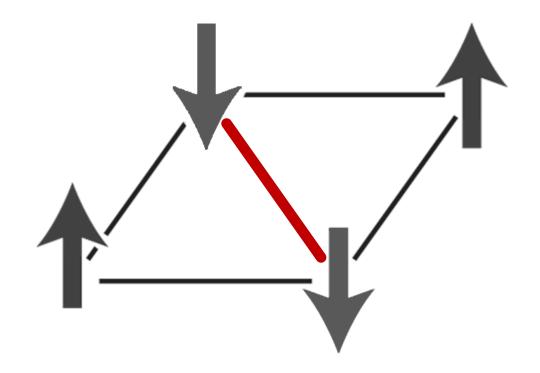






[wikipedia]

local Hamiltonians



a global ground state

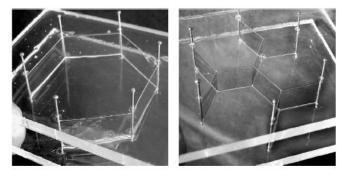
HARD?

finding it? describing it?



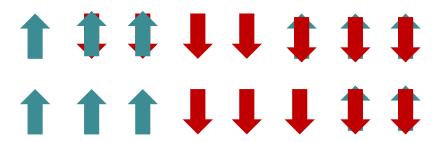
The ground state is cool

doesn't nature find it? can't we cool the system?



[Dutta+, arXiv:0806.1340]

simulated annealing:a nature-inspired heuristic



optimization, CSP & spin glasses

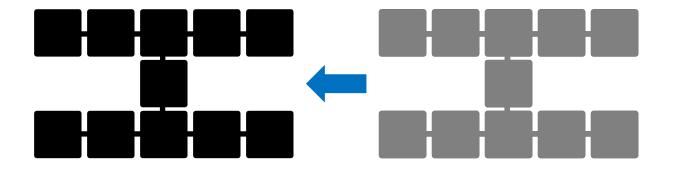
The questions we ask: static

- eigenstates?ground state(s)excitationsspectrum
- solvable? integrable?symmetriestransformationsnature of states
- bulk properties?
 phase transitions



[Oona Räisänen]

dynamical the questions



The state of the s



scattering

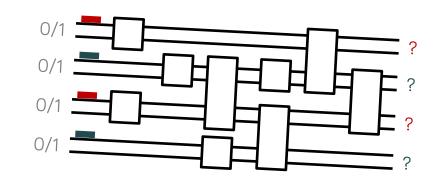
Universal computation by multi-particle quantum walk

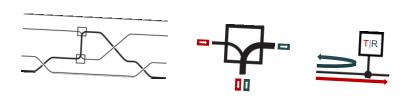
dual-rail encodingN wavepackets

$$a_j^{\dagger} a_k + a_k^{\dagger} a_j$$

CPHASE: interaction

$$a_j^{\dagger} a_k^{\dagger} a_j a_k$$





[Childs, Gosset, Webb, Science 339, 791 (2013)]

simulating dynamics

trouble: states become more entangled

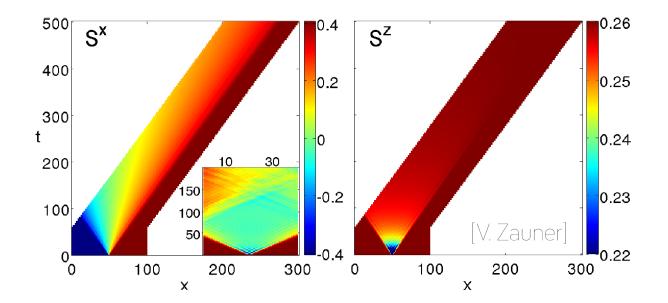
$$i\frac{d}{dt}|\psi(t)\rangle = H(t)|\psi(t)\rangle$$

a costly cut

- simulating dynamicstrouble: states become more entangled
- thermalization, quenching



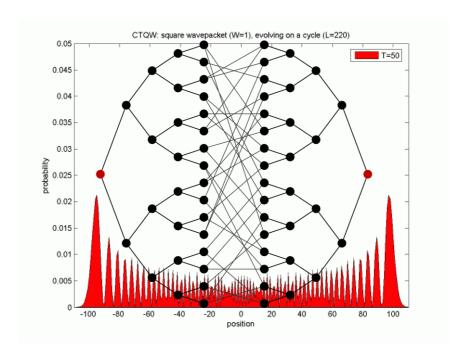
- simulating dynamics trouble: states become more entangled
- thermalization, quenching the behavior of correlations Lieb-Robinson bound $\|[A(t),B]\| \le c \|A\| \|B\| \exp(-\mu(dist(X,Y)-v|t|))$

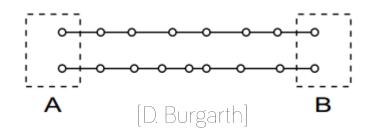


simulating dynamics

thermalization

utilizing the dynamics control (state preparation) transport (signals, energy) simulation algorithms!





the questions

three questions three questions three questions

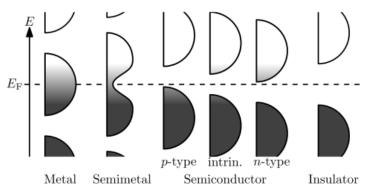
eigenstates & spectral properties symmetries & solvability parameters & transitions

The spectrum of a system: band gaps

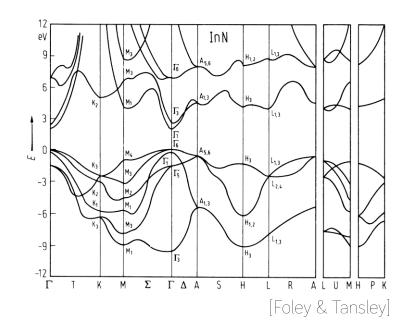
a lattice + electrons

orbitals... available levels filled with electrons

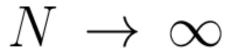
conductors (metals) semiconductors insulators



[wikimedia: Nanite]

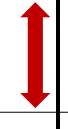


The (ground state) gap in a system



Are there states close to the ground state when we take the thermodynamic limit?

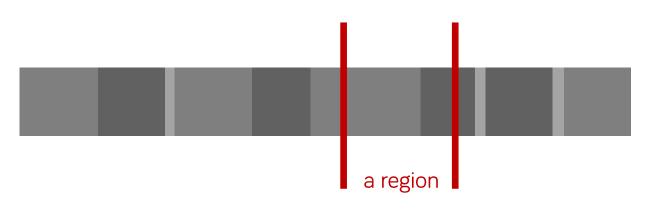
an inverse-poly gap?
$$\Delta = \frac{c}{N} \, \rightarrow \, 0$$



The (ground state) gap in a system

Do we have a gap?

■ YES? area law – 1D [Hastings, Arad Landau Vazirani] inspired numerics (MPS,...)



correlations fall off exponentially with distance

$$S(\rho) = -\text{Tr}(\rho \log_2 \rho)$$

entropy from
quantum correlations

The (ground state) gap in a system

Do we have a gap?

■ YES? area law inspired numerics (MPS, ...)

imaginary time evolution

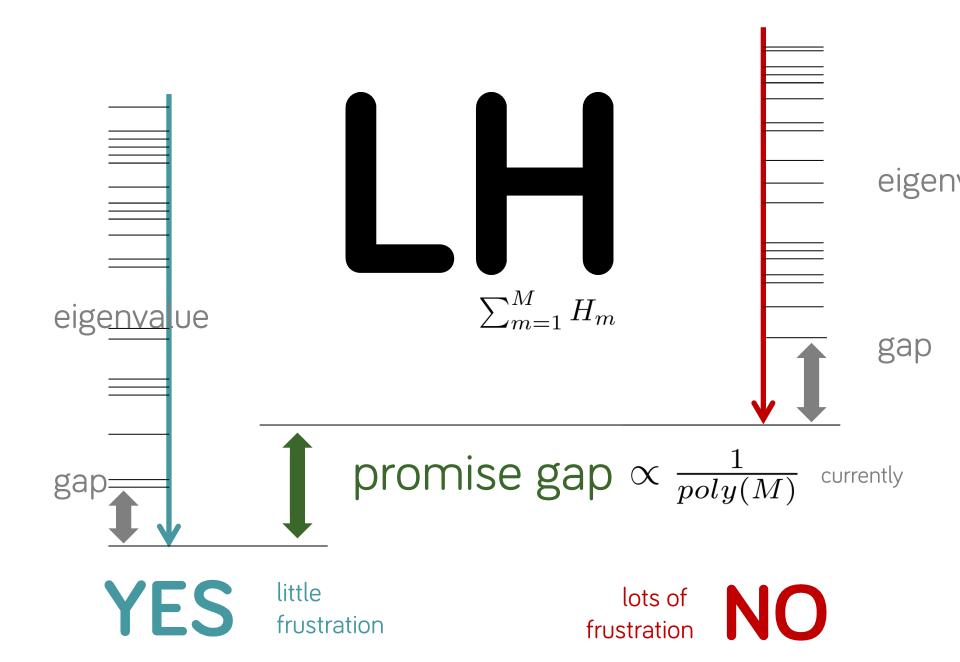
$$|\psi(t)\rangle = e^{-iHt}|\psi(0)\rangle$$

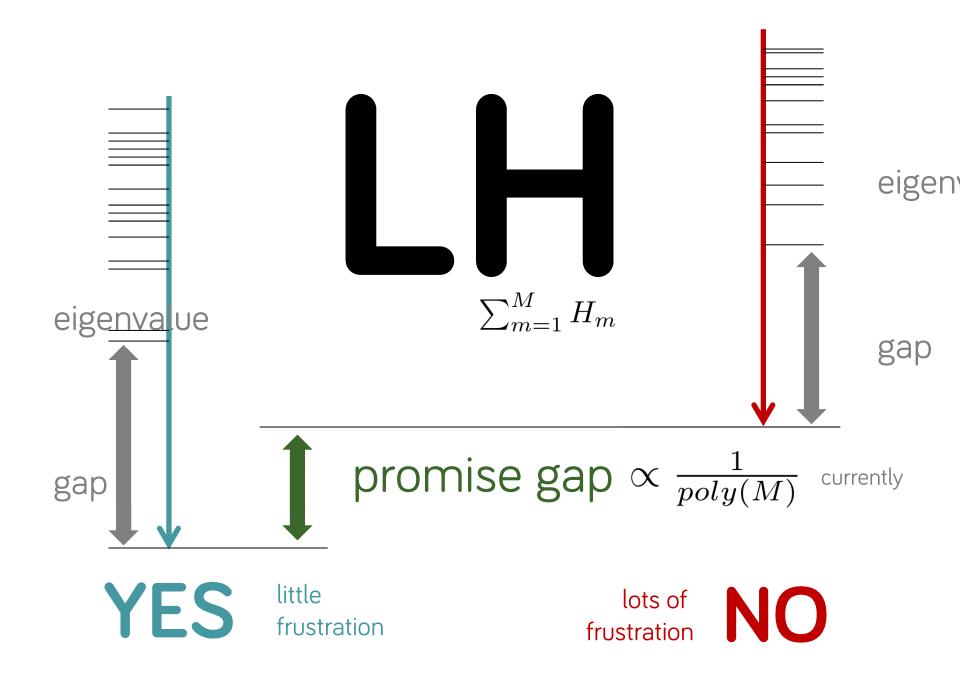
$$e^{-iH(-i\tau)}|\psi\rangle = e^{-H\tau}|\psi\rangle$$

$$e^{-H\tau} \approx e^{-H_1\tau}e^{-H_2\tau}$$

NO? what is nearby?
qPCP, no trivial states [Hastings]

deciding it: uncomputable [Cubitt Perez-Garcia Wolf]





1 "Solving" a system

exactly solvable

the free energy/partition function

$$E(s) = \sum_{j} s_j s_{j+1} - h \sum_{j} s_j$$

integrable systems

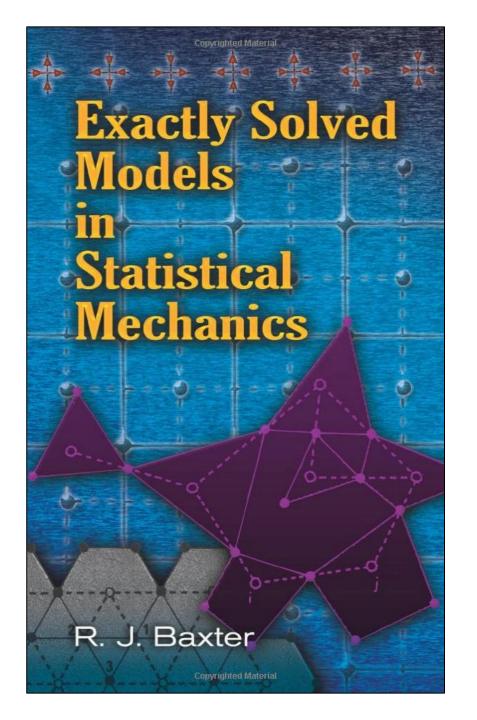
$$H = -J \sum_{n=1}^{N} \mathbf{S}_n \cdot \mathbf{S}_{n+1}$$

1D Heisenberg, conserved S_z

1- excitation waves

2-, 3-, ... algebraic Bethe ansatz

... new equations & lots of work left!



1 What can we exactly diagonalize ... today?

■ Spin S=1/2 models:

40 spins square lattice, 39 sites triangular, 42 sites Honeycomb lattice 48 sites kagome lattice 64 spins or more in elevated magnetization sectors up to ~500 billion basis states

Fractional quantum hall effect
different filling fractions ν, up to 16-20 electrons
up to 3.5 billion basis states

Hubbard models (~ Full CI in Quantum Chemistry)

20 sites square lattice at half filling, 21 sites triangular lattice

24 sites honeycomb lattice

up to 160 billion basis states

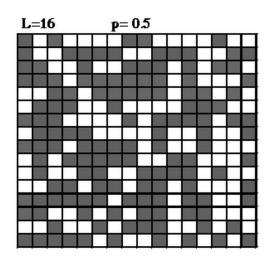
(low-lying states only)

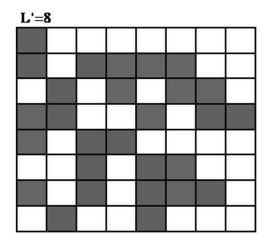
[A.Lauchli]

1 Trying to understand a system

- exact diagonalization
- Monte Carlo sampling
- small clusters
- mean-field averaging over neighbors
- series expansions
- renormalization group relevant lengthscales
- replica trick $\lim_{n \to 0} \frac{Z^n 1}{n} = \ln Z$

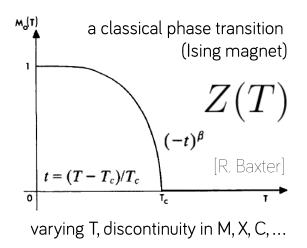
[D. Ceperley]





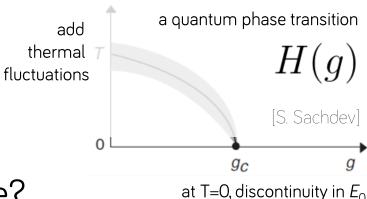
What's going on in a system

varying parameters =
 varying nature?
 thermal fluctuations
 response functions



phase transitions

order/disorder insulator/conductor power/exp correlation decay



is there one? what is its type?

scaling, critical exponents universality classes

what happens at T>0?

Phase transition in the quantum Ising chain

■ in a transverse field

$$H_I = -Jg \sum_{i} \hat{\sigma}_i^x - J \sum_{\langle ij \rangle} \hat{\sigma}_i^z \hat{\sigma}_j^z$$



$$g \ll 1$$

$$g = g_{c}$$

$$g = g_c$$

$$\sim \frac{1}{|x_i - x_j|^{d-2+\eta}}$$

$$g \gg 1$$

$$\langle 0|\hat{\sigma}_i^z\hat{\sigma}_j^z|0\rangle \sim N_0^2$$

 $\langle 0|\hat{\sigma}_i^z|0\rangle \sim N_0$

$$\sim \frac{1}{|x_i - x_j|^{d-2+\eta}}$$

$$\sim e^{-|x_i-x_j|/\xi}$$

frustration, states get hard to describe correlations grow, area law corrections critical scaling, gap closes (finite-size: avoided crossing)

1

Phase transitions in CS

the SAT/UNSAT transition in random SAT

solutions clustering
algorithms work/don't work
statistical physics analysis
algorithms (cavity belief su



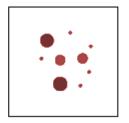
[Parisi, Mezard, Zecchina, Monasson]

rigorous analysis [Coja-Oghlan]







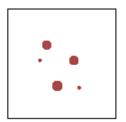


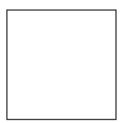
0.8

0.6

0.4

0.2





4000

3500

3000

2500

2000

1500

1000

500

[R. Zecchina]

[L. Zdeborová]

classifying the difficulties

classical

The classical Ising spin glass can be hard (NP-C)

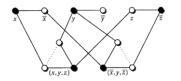
2D, varying couplings, external field | Barahona84|

$$E(s) = -\sum_{\langle i,j\rangle} J_{ij} s_i s_j - H \sum_i s_i$$
 Uniform H varying J

$$= -\sum_{E} J_{ij} - HN + 2\sum_{E^{\pm}} J_{ij} + 2HF_{\bullet}^{-}$$

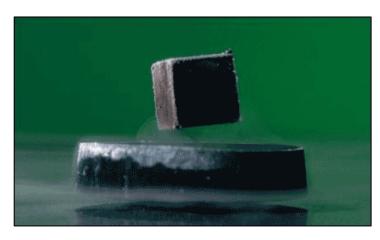
$$= const. - \sum_{E^{\pm}} w_{e}$$

MAX weighted CUT & finding E_{Ω}



[C. Moore, S. Mertens]

- OPTIMIZATION find the ground state
- UNDERSTANDING THE SYSTEM explain/predict expectation values



[northeastern.edu]

2 Partition function & counting

all the properties of the system derivatives... properties responses, excitations

$$Z = \sum_{s} e^{-\beta E(s)}$$
$$p(s) = \frac{1}{Z} e^{-\frac{E(s)}{kT}}$$

counting degeneracies

Z dominated by the ground states at low T Ising, planar graphs: computing determinants [Kasteleyn, Fisher] nonplanar: #P complete (Potts, perf. matchings, permanents)

sampling & approximation Markov Chain Monte Carlo [Jerrum Sinclair]

$$Z(T_n) = \frac{Z_n}{Z_{n-1}} \underbrace{Z_{n-1}}_{Z_{n-2}} \cdots \underbrace{Z_1}_{Z_0} Z_0$$

the quantumness

the quantum mess

Quantum Many-Body Hamiltonians

terms

non-commuting
$$H=-\sum_{\langle i,j\rangle}J_{ij}\sigma_i^z\sigma_j^z\,-\sum B_i\sigma_i^x$$
 terms

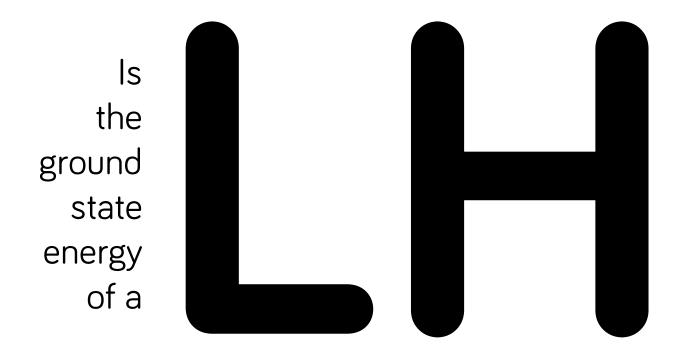
entanglement (vs. products) hard-to-find, hard-to describe ground states

Schmidt decomposition

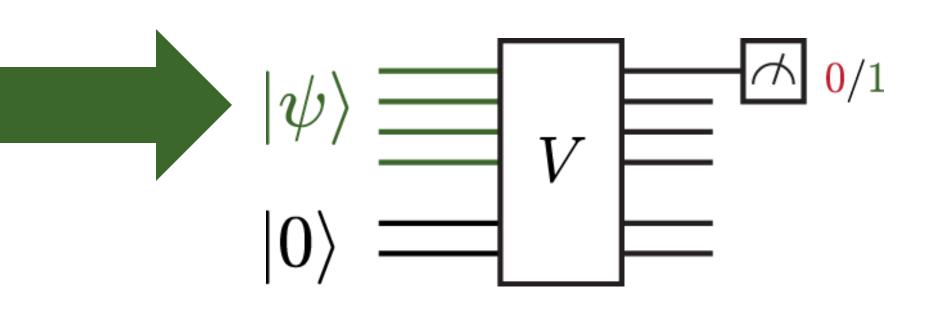
sition
$$\sum_k \lambda_k |a_k
angle \otimes |b_k
angle$$
 a costly cut

• the basic hard question: E_{\cap}

The Local Hamiltonian problem

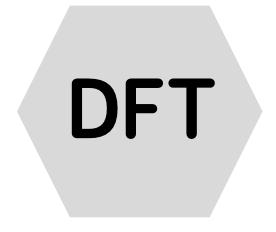


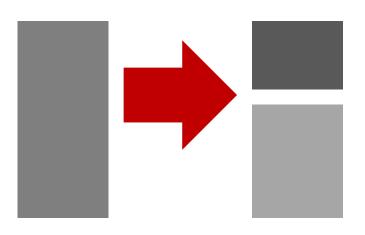




Accept a good proof with p > a. $\begin{bmatrix} a \\ b \end{bmatrix}$ Probability of accepting p < b.

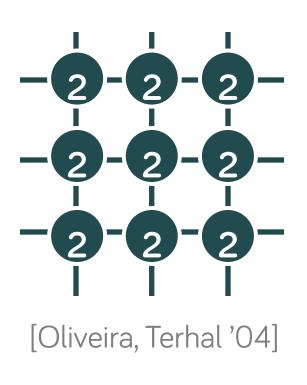




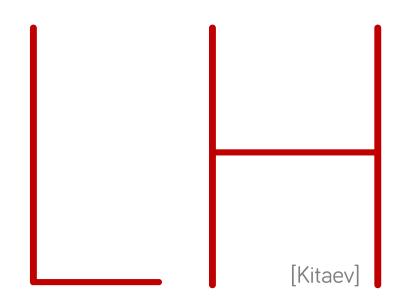




2-local Hamiltonian is QMA complete



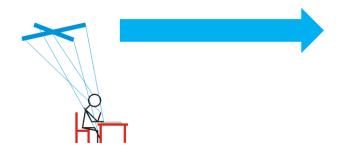
a global minimum



$$\sum H_{jk}$$
 -8-8-8-8-8-

[Hallgren, N, Narayanaswami '13]

clock/work registers

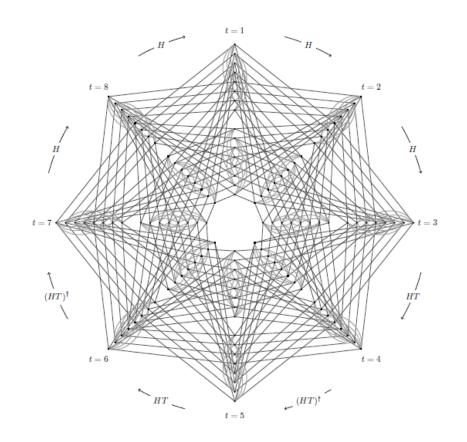


Can we get away from this?

QMA-complete problems that are more "natural"?

Problem 4 (Minimum Graph Eigenvalue). We are given a $D \times D$ symmetric 0-1 matrix A, specified by a classical deterministic circuit that takes as input a row $r \in [D]$ and computes the locations of the nonzero entries in that row. We are also given a real number a and a precision parameter $\epsilon = \frac{1}{N}$, where $N \in \mathbb{N}$ is specified in unary. We are promised that either the smallest eigenvalue of A is at most a (yes instance), or else it is at least $a + \epsilon$ (no instance), and asked to decide which is the case.

[Childs Gosset Webb 13]



Can we get away from this?

QMA-complete problems?

Universality & hardness?

$$H = t \sum_{\langle i,j \rangle} b_i^{\dagger} b_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1)$$

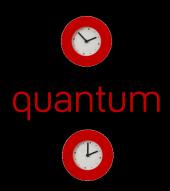
[Childs Gosset Webb 13]

many-body physics & complexity

tough & interesting questions



MODELS



What are they like?

ground state (energy)

QMA-complete problems

$$\sum_{k} H_k(t)$$

What are they good for?

control, transport, chemistry, making a (q) computer BQP universality local particle dimension



interaction geometry



time independence



translational invariance



promise and eigenvalue gaps



energyx time cost



many-body physics & complexity &

the problems & the tools

Daniel Nagaj wien wien